

# van der Pol Expansions of L-Series

David Borwein and Jonathan Borwein

*Abstract.* We provide concise series representations for various L-series integrals. Different techniques are needed below and above the abscissa of absolute convergence of the underlying L-series.

## 1 Preliminaries

In [8] the following odd looking integral evaluation is obtained.

$$(1) \quad \int_0^\infty \frac{(3 - 2\sqrt{2} \cos(t \log 2)) |\zeta(1/2 + it)|^2}{t^2 + 1/4} dt = \pi \log 2.$$

This identity turns out, formally, to be a case of a rather pretty, and perhaps useful, class of L-series evaluations given in Theorem 1 and Corollary 1 (*cf.* [2]). In Theorem 3 we recover (1) entirely rigorously.

Given a *Dirichlet series*

$$\lambda(s) := \sum_{n=1}^{\infty} \frac{\lambda_n}{n^s}, s = \sigma + i\tau, \sigma = \Re s > 0,$$

we consider the integral

$$\iota_\lambda(\sigma) := \frac{1}{2} \int_{-\infty}^{\infty} \left| \frac{\lambda(s)}{s} \right|^2 d\tau$$

as a function of  $\lambda$ . Observe that when the coefficients  $\lambda_n$  are real,

$$\iota_\lambda(\sigma) = \int_0^\infty \left| \frac{\lambda(s)}{s} \right|^2 d\tau,$$

but that this is not necessarily so when the coefficients are complex. We refer to [3, 5, 7] for other, largely standard details.

---

Received by the editors January 25, 2005; revised August 18, 2005.

The research of the first author was supported by NSERC. The research of the second author was supported by NSERC, the Canada Foundation for Innovation and the Canada Research Chair Program.

AMS subject classification: 11M35, 11M41, 30B50.

Keywords: Dirichlet series integrals, Hurwitz zeta functions, Plancherel theorems, L-series.

©Canadian Mathematical Society 2007.

## 2 Integrals Involving $s$ with Large Real Part

It is convenient to recall [2] that for  $u, a > 0$ ,

$$(2) \quad \int_{-\infty}^{\infty} \frac{\cos(at)}{t^2 + u^2} dt = \frac{\pi}{u} e^{-au}.$$

**Theorem 1** For  $a(s) := \sum_{n=1}^{\infty} a_n n^{-s}$ ,  $b(s) := \sum_{n=1}^{\infty} b_n n^{-s}$ , and  $s = \sigma + i\tau$  with fixed  $\sigma = \Re(s) > 0$  such that both Dirichlet series are absolutely convergent, we have

$$(3) \quad \iota_{a,b}(\sigma) := \frac{1}{2} \int_{-\infty}^{\infty} \frac{a(s) \bar{b}(s)}{\sigma^2 + \tau^2} d\tau = \frac{\pi}{2\sigma} \sum_{n=1}^{\infty} \frac{A_n \bar{B}_n - A_{n-1} \bar{B}_{n-1}}{n^{2\sigma}},$$

where  $A_n := \sum_{k=1}^n a_k$ ,  $B_n := \sum_{k=1}^n b_k$ ,  $A_0 := B_0 := 0$ .

**Proof** Let  $a_N(s) := \sum_{n=1}^N a_n n^{-s}$ ,  $b_N(s) := \sum_{n=1}^N b_n n^{-s}$ . Then, in view of (2), we have

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{a_N(s) \bar{b}_N(s)}{\sigma^2 + \tau^2} d\tau &= \int_{-\infty}^{\infty} \frac{\sum_{N \geq n, m > 0} a_n \bar{b}_m n^{-\sigma+i\tau} m^{-\sigma-i\tau}}{\sigma^2 + \tau^2} d\tau \\ &= \sum_{N \geq n > m > 0} \frac{a_n \bar{b}_m}{(nm)^\sigma} \int_{-\infty}^{\infty} \frac{\cos(\tau \log(n/m))}{\sigma^2 + \tau^2} d\tau \\ &\quad + \sum_{N \geq n > m > 0} \frac{a_n \bar{b}_n}{(nm)^\sigma} \int_{-\infty}^{\infty} \frac{\cos(\tau \log(m/n))}{\sigma^2 + \tau^2} d\tau + \sum_{n=1}^N \frac{a_n \bar{b}_n}{n^{2\sigma}} \int_{-\infty}^{\infty} \frac{1}{\sigma^2 + \tau^2} d\tau \\ &\quad + i \sum_{N \geq m, n > 0} \frac{a_n \bar{b}_m}{(nm)^\sigma} \int_{-\infty}^{\infty} \frac{\sin(\tau \log(n/m))}{\sigma^2 + \tau^2} d\tau \\ &= \sum_{N \geq n > m > 0} \frac{a_n \bar{b}_m + a_m \bar{b}_n}{(nm)^\sigma} \int_{-\infty}^{\infty} \frac{\cos(\tau \log(n/m))}{\sigma^2 + \tau^2} d\tau + \frac{\pi}{\sigma} \sum_{n=1}^N \frac{a_n \bar{b}_n}{n^{2\sigma}} \\ &= \frac{\pi}{\sigma} \sum_{N \geq n > m > 0} \frac{a_n \bar{b}_m + a_m \bar{b}_n}{(nm)^\sigma (n/m)^\sigma} + \frac{\pi}{\sigma} \sum_{n=1}^N \frac{a_n \bar{b}_n}{n^{2\sigma}} \\ &= \frac{\pi}{\sigma} \sum_{n=1}^N \frac{a_n \bar{B}_{n-1} + A_{n-1} \bar{b}_n + a_n \bar{b}_n}{n^{2\sigma}} = \frac{\pi}{\sigma} \sum_{n=1}^N \frac{A_n \bar{B}_n - A_{n-1} \bar{B}_{n-1}}{n^{2\sigma}}. \end{aligned}$$

Note that the imaginary part in the above evaluation vanished because we integrated an odd function over the range  $-\infty < \tau < \infty$ .

Next, we observe that

$$|a_N(s) \bar{b}_N(s)| \leq \sum_{n=1}^{\infty} \frac{|a_n \bar{b}_n|}{n^{2\sigma}} = M < \infty,$$

where  $M$  is independent of  $\tau$ . Hence

$$\left| \frac{a_N(s)\bar{b}_N(s)}{\sigma^2 + \tau^2} \right| \leq \frac{M}{\sigma^2 + \tau^2},$$

and (3) follows by Lebesgue's theorem on dominated convergence on letting  $N \rightarrow \infty$ . ■

As an immediate consequence we have:

**Corollary 1** If  $\lambda(s) := \sum_{n=1}^{\infty} \lambda_n n^{-s}$  with  $s = \sigma + i\tau$  and fixed  $\sigma = \Re(s) > 0$  such that the Dirichlet series is absolutely convergent, then

$$(4) \quad \iota_\lambda(\sigma) := \frac{1}{2} \int_{-\infty}^{\infty} \left| \frac{\lambda(s)}{s} \right|^2 d\tau = \frac{\pi}{2\sigma} \sum_{n=1}^{\infty} \frac{|\Lambda_n|^2 - |\Lambda_{n-1}|^2}{n^{2\sigma}},$$

where  $\Lambda_n := \sum_{k=1}^n \lambda_k$ ,  $\Lambda_0 := 0$ .

If, in addition, all the coefficients  $\lambda_n$  are real, then

$$(5) \quad \iota_\lambda(\sigma) = \int_0^{\infty} \left| \frac{\lambda(s)}{s} \right|^2 d\tau = \frac{\pi}{2\sigma} \sum_{n=1}^{\infty} \frac{\Lambda_n^2 - \Lambda_{n-1}^2}{n^{2\sigma}}.$$

Note that, by Dirichlet's test [9], the final series in (5) is convergent for all  $\sigma > 0$  when  $\Lambda_n$  is bounded, but we cannot automatically guarantee that it is equal to the integral in this case, or even that the integral is finite. Simple continuation arguments will not work. Of course, similar difficulties arise with regard to (3) and (4). This in part motivates the first example and the following section.

It is, however, easy now to check that

$$\langle a, b \rangle_\sigma := \frac{1}{2} \int_{-\infty}^{\infty} \frac{a(s)\bar{b}(s)}{\sigma^2 + \tau^2} d\tau$$

defines an extended-value inner product on the space of Dirichlet series such that  $\langle \alpha, \alpha \rangle_\sigma = \iota_\alpha(\sigma)$ , which is typically finite for  $\sigma$  large enough.

In the sequel, we let  $L_\mu(s) := \sum_{n=1}^{\infty} (\frac{\mu}{n}) n^{-s}$  denote the primitive L-function corresponding to the Kronecker symbol  $(\frac{\mu}{n})$ , [3]. Below,  $\lfloor x \rfloor$  is the integer part and  $\lceil x \rceil$  is the truncation of  $x$ , so that

$$\lceil x \rceil = \begin{cases} \lfloor x \rfloor & \text{when } x \geq 0, \\ \lceil x \rceil & \text{when } x < 0. \end{cases}$$

As usual,  $\{x\} := x - \lfloor x \rfloor$  denotes the fractional part of  $x$ .

**Example 1** For the Riemann zeta function ( $\zeta = L_1$ ), and for  $\sigma > 1$ , Corollary 1 applies and yields

$$\frac{\sigma}{\pi} \iota_\zeta(\sigma) = \zeta(2\sigma - 1) - \frac{1}{2} \zeta(2\sigma),$$

as  $\lambda_n = 1$  and  $\Lambda_n = n$ . By contrast it is known [6, (69)] that on the critical line  $s = \frac{1}{2} + i\tau$

$$\frac{1}{2\pi} \iota_\zeta\left(\frac{1}{2}\right) = \log(\sqrt{2\pi}) - \frac{1}{2} \gamma.$$

More broadly, for  $0 < \sigma < 1$ , Crandall (recorded in [6]) has found that

$$(6) \quad \iota_\zeta(\sigma) = \pi \int_0^1 \sum_{n=0}^{\infty} \frac{\theta^2}{(n + \theta)^{1+2\sigma}} d\theta = \pi \int_0^1 \theta^2 \zeta(1 + 2\sigma, \theta) d\theta,$$

where

$$\zeta(s, a) := \sum_{n=0}^{\infty} (n + a)^{-s}$$

is the Hurwitz zeta function, which is easy to compute. This devolves from *van der Pol's representation*

$$(7) \quad \frac{\zeta(s)}{s} = - \int_{-\infty}^{\infty} e^{-\sigma\omega} (e^\omega - \lfloor e^\omega \rfloor) e^{-i\tau\omega} d\omega, \quad s = \sigma + i\tau \text{ with } 0 < \sigma < 1.$$

One way to obtain identity (7) is to note that

$$\begin{aligned} \int_{-\infty}^{\infty} e^{-s\omega} (e^\omega - \lfloor e^\omega \rfloor) d\omega &= \int_0^{\infty} t^{-s-1} (t - \lfloor t \rfloor) dt = \sum_{n=0}^{\infty} \int_n^{n+1} t^{-s-1} (t - \lfloor t \rfloor) dt \\ &= \int_0^1 \theta \zeta(1 + s, \theta) d\theta = \sum_{n=0}^{\infty} \int_0^1 \theta (\theta + n)^{-s-1} d\theta \\ &= -\frac{1}{s} \lim_{N \rightarrow \infty} \left( \sum_{n=1}^N \frac{1}{n^s} - \frac{N^{1-s}}{1-s} \right), \end{aligned}$$

on integrating once by parts. Now set

$$\sigma_N := \sum_{n=1}^N \frac{1}{n^s} - \frac{N^{1-s}}{1-s}$$

and observe that we have shown that  $\sigma_N$  converges to some number  $\sigma_\infty$ . Further,

$$\sigma_{2N} - 2^{1-s} \sigma_N = \sum_{n=1}^{2N} \frac{1}{n^s} - 2 \sum_{n=1}^N \frac{1}{(2n)^s} = \sum_{n=1}^{2N} \frac{(-1)^{n+1}}{n^s} \rightarrow (1 - 2^{1-s}) \zeta(s),$$

which implies that  $\sigma_\infty = \zeta(s)$ , as required.

Another way to establish (7) (kindly suggested by David Bradley) is to start with the standard identity

$$\frac{\zeta(s)}{s} = \frac{1}{s-1} - \int_1^\infty x^{s-1}(x - [x]) dx,$$

which can easily be derived under the restriction  $\Re s > 1$  and then extended by analytic continuation to the punctured half-plane  $\{s, \Re s > 0, s \neq 1\}$ . This, together with the identity

$$\frac{1}{s-1} = - \int_0^1 x^{s-1}(x - [x]) dx,$$

which is valid for  $\Re s < 1$ , yields (7) via the change of variable  $x = e^\omega$ .

We now write (7) as a Fourier transform:

$$\frac{\zeta(s)}{s} = -\mathcal{F}_\tau(e^{-\sigma\omega}\{e^\omega\})$$

and so obtain for  $0 < \sigma < 1$ ,

$$2\iota_\zeta(\sigma) = \int_{-\infty}^\infty \left| \frac{\zeta(s)}{s} \right|^2 d\tau = 2\pi \int_{-\infty}^\infty e^{-2\sigma\omega} |e^\omega - [e^\omega]|^2 d\omega,$$

from the  $\mathcal{L}_2$  Plancherel theorem [9].

It follows that

$$\begin{aligned} \int_{-\infty}^\infty e^{-2\sigma\omega} |e^\omega - [e^\omega]|^2 d\omega &= \int_0^\infty t^{-2\sigma-1} |t - [t]|^2 dt \\ &= \sum_{n=0}^\infty \int_n^{n+1} t^{-2\sigma-1} |t - [t]|^2 dt = \sum_{n=0}^\infty \int_0^1 \theta^2 (\theta + n)^{-2\sigma-1} d\theta \\ &= \int_0^1 \theta^2 \sum_{n=0}^\infty (\theta + n)^{-2\sigma-1} d\theta, \end{aligned}$$

as required to yield (6). Note that all terms are absolutely convergent, which legitimates the operations. We have also established, *inter alia*, that for  $0 < \sigma < 1$ ,

$$\frac{1}{2\pi} \int_{-\infty}^\infty \left| \int_0^1 \theta \zeta(1 + s, \theta) d\theta \right|^2 d\tau = \int_0^1 \theta^2 \zeta(1 + 2\sigma, \theta) d\theta.$$

Moreover, reversing the order of integration and summation above leads to

$$\iota_\zeta(\sigma) = \pi \sum_{n=0}^\infty \int_0^1 \frac{\theta^2}{(n + \theta)^{1+2\sigma}} d\theta = -\frac{\pi}{2\sigma} \left( \frac{2\zeta(2\sigma - 1)}{2\sigma - 1} + \zeta(2\sigma) \right)$$

which in the limit as  $\sigma \rightarrow \frac{1}{2}$  recaptures the evaluation quoted above. Recapitulating, we have

$$(8) \quad \frac{\sigma}{\pi} \iota_{\zeta}(\sigma) = \begin{cases} -\frac{1}{2}\zeta(2\sigma) - \frac{\zeta(2\sigma-1)}{2\sigma-1} & 0 < \sigma < 1, \\ -\frac{1}{2}\zeta(2\sigma) + \zeta(2\sigma - 1) & 1 < \sigma < \infty. \end{cases}$$

There are similar formulae for  $s \mapsto \zeta(s - k)$  with  $k$  integral. For instance, applying the result in (5) with  $\zeta_1 := t \mapsto \zeta(t + 1)$  yields

$$\frac{1}{\pi} \int_0^\infty \frac{|\zeta(3/2 + i\tau)|^2}{1/4 + \tau^2} d\tau = \frac{1}{\pi} \iota_{\zeta_1}\left(\frac{1}{2}\right) = 2\zeta(2, 1) + \zeta(3) = 3\zeta(3),$$

on using Euler’s result (see [4]) that  $\zeta(2, 1) := \sum_{n=2}^\infty \frac{1}{n^2} \sum_{k=1}^{n-1} \frac{1}{k} = \zeta(3)$ . ■

**Example 2** For the *alternating zeta function*,  $\alpha := s \mapsto (1 - 2^{1-s})\zeta(s)$ , we recover, as in [7], that

$$\frac{\sigma}{\pi} \iota_{\alpha}(\sigma) = \frac{1}{2}\alpha(2\sigma),$$

as  $\lambda_n = (-1)^{n+1}$ ,  $\Lambda_n = (1 - (-1)^n)/2$  and  $\Lambda_n^2 - \Lambda_{n-1}^2 = (-1)^{n+1}/2$ .

Set  $\sigma := 1/2$ . Then

$$\begin{aligned} \frac{|\alpha(1/2 + it)|^2}{1/4 + t^2} &= \frac{|1 - 2^{1/2-it}|^2 |\zeta(1/2 + it)|^2}{1/4 + t^2} \\ &= (3 - 2\sqrt{2} \cos(t \ln(2))) \frac{|\zeta(1/2 + it)|^2}{1/4 + t^2} \end{aligned}$$

is precisely the integrand in (1). Thus, since  $\alpha(2\sigma) = \log 2$ , we see that  $\iota_{\alpha}(1/2) = \pi \log 2$  is the evaluation in [7]. Note that to justify the exchange of sum and integral implicit in (5), we should have to analyse more carefully the integrand, since  $1/2$  is below the abscissa of absolute convergence of the series, and note that this would not have been legitimate in Example 1 because of (8). This motivates the approach in the section below on the Hurwitz zeta function. ■

**Example 3**

(a) For the *Catalan zeta function* ( $\beta = L_{-4}$ ), and for  $\sigma > 1$  :

$$\frac{\sigma}{\pi} \iota_{\beta}(\sigma) = \frac{1}{2}\beta(2\sigma),$$

as  $\lambda_{2n} = 0$ ,  $\lambda_{2n+1} = (-1)^n$  and again  $\Lambda_n^2 - \Lambda_{n-1}^2 = \lambda_n$ .

(b) For  $L_{-3}$ , the same pattern holds, in that  $\frac{\sigma}{\pi} \iota_{L_{-3}}(\sigma) = \frac{1}{2} L_{-3}(2\sigma)$ , but not for  $L_5, L_{-7}$ , and so on.

(c) In general the series  $L_{\pm d}$  does not lead to output which is again a primitive L-series modulo  $d$ . For example,

$$\frac{\sigma}{\pi} \iota_{L_5}(\sigma) = - \sum_{5|n} \frac{(-1)^{n \bmod 5}}{n^{2\sigma}}, \text{ and } \frac{\sigma}{\pi} \iota_{L_{-8}}(\sigma) = L_{-8}(2\sigma) - \frac{1}{2} L_{-4}(2\sigma).$$

These are not character sums, though always the coefficients repeat modulo  $d$ .

(d) Finally, let  $\vartheta := s \mapsto (1 + 2^{-s})^{-1}$ . Then again  $\frac{\sigma}{\pi} \iota_{\vartheta}(\sigma) = \frac{1}{2} \vartheta(2\sigma)$ . ■

For each of these examples, which are discussed further in [2], the defect of Theorem 1 and Corollary 1 is that, as we have seen, they only directly apply when  $\sigma$  is large enough. The van der Pol approach offers a nice alternative, especially in the critical strip.

### 3 van der Pol's Approach to Hurwitz zeta

Recall that

$$\zeta(s, a) := \sum_{n=0}^{\infty} \frac{1}{(n+a)^s}, \quad a > 0,$$

is initially defined for  $\Re s > 1$  and appropriately analytically continued for  $\Re s < 1$ . Thus,  $\zeta(s) = \zeta(s, 1)$ . We first work out a Fourier transform representation for  $\zeta(a, s)$ .

**Proposition 1** For  $1 \geq a > 0$  and  $1 > \Re s > 0$ , we have

$$(9) \quad \int_{-\infty}^{\infty} e^{-s\omega} (e^{\omega} - \lceil e^{\omega} - a \rceil) d\omega = -\frac{1}{s} \zeta(s, a + 1).$$

**Proof** (i) Observe that the formula uses the truncation of  $e^w - a$  and not the floor, and that the two only differ in the interval  $-\infty < w < \log a$ . We have

$$\begin{aligned} (10) \quad \int_{\log a}^{\infty} e^{-s\omega} \{e^{\omega} - a\} d\omega &= \int_a^{\infty} t^{-s-1} (t - a - \lfloor t - a \rfloor) dt \\ &= \sum_{n=0}^{\infty} \int_{n+a}^{n+1+a} t^{-s-1} (t - a - \lfloor t - a \rfloor) dt \\ &= \int_0^1 \theta \zeta(1 + s, \theta + a) d\theta = \sum_{n=0}^{\infty} \int_0^1 \theta (\theta + n + a)^{-s-1} d\theta \\ &= -\frac{1}{s} \lim_{N \rightarrow \infty} \left( \sum_{n=1}^N \frac{1}{(n+a)^s} - \frac{(N+a)^{1-s} - a^{1-s}}{1-s} \right) \\ &= -\frac{1}{s} \zeta(s, a + 1) - \frac{a^{1-s}}{s(1-s)}, \end{aligned}$$

this evidently being true for  $\Re s > 1$  and so for  $0 < \Re s < 1$  by analytic continuation: both sides of (10) clearly being analytic for  $\Re s > 0$ . Incidentally, we have shown that

(10) is valid for  $\Re s > 0$ ,  $a > 0$ . The additional restriction  $\Re s < 1$ ,  $a \leq 1$  is needed for the next part of the proof.

It follows from (10) that

$$\begin{aligned} & \int_{-\infty}^{\infty} e^{-s\omega} (e^{\omega} - [e^{\omega} - a]) d\omega \\ &= \int_{\log a}^{\infty} e^{-s\omega} (\{e^{\omega} - a\} + a) d\omega + \int_{-\infty}^{\log a} e^{(1-s)\omega} d\omega \\ &= -\frac{1}{s} \zeta(s, a+1) - \frac{a^{1-s}}{s(1-s)} + \frac{a^{1-s}}{s} + \frac{a^{1-s}}{1-s} = -\frac{1}{s} \zeta(s, a+1), \end{aligned}$$

and this establishes (9). ■

(ii) Another proof of equation (10) (suggested by the David Bradley) is to proceed from Theorem 12.2 in Apostol [1, p. 269] which states that for  $0 < a \leq 1$ ,  $\Re s > 0$ ,  $N = 0, 1, 2, \dots$ ,

$$(11) \quad \zeta(s, a) := \sum_{n=0}^N \frac{1}{(n+a)^s} - \frac{(N+a)^{1-s}}{s-1} - s \int_N^{\infty} \frac{\{x\}}{(x+a)^{s+1}} dx.$$

Putting  $N = 0$  in (11) and making the change of variable  $x = e^{\omega} - a$  gives equation (10).

**Proposition 2** For  $\Re s > 0$  and  $a \geq 0$ , we have

$$\sigma(s, a) := \lim_{N \rightarrow \infty} \left( \sum_{n=1}^N \frac{1}{(n+a)^s} - \frac{(N+a)^{1-s}}{1-s} \right) = \zeta(s, a+1).$$

Another form of the limit is

$$\sigma(s, a) = -s \int_0^1 t \zeta(1+s, a+t) dt,$$

and indeed  $\sigma(s, 0) = \zeta(s)$ .

**Proof** The case  $a > 0$  follows immediately from the proof of (10) in part (i) above, and the case  $a = 0$  was treated in the earlier proof of (7). We can also derive the case  $1 \geq a > 0$  by letting  $N \rightarrow \infty$  in (11) and observing that  $\zeta(s, a) - a^{-s} = \zeta(s, a+1)$ . ■

Most of what follows concerns Dirichlet series having coefficients which repeat modulo  $N$ .

**Theorem 2** Let  $\lambda(s) := \sum_{n=0}^{\infty} \lambda_n n^{-s}$ ,  $s = \sigma + i\tau$ , with coefficients  $\lambda_n$  repeating modulo  $N$  (i.e.,  $\lambda_{N+n} = \lambda_n$ ). Then for  $0 < \sigma < 1$ ,

$$(12) \quad - \int_{-\infty}^{\infty} e^{-\sigma\omega} \sum_{k=1}^N \lambda_k \left( \frac{e^{\omega}}{N} - \left\lceil \frac{e^{\omega} + N - k}{N} \right\rceil \right) e^{-i\tau\omega} d\omega = \frac{\lambda(s)}{s}.$$

**Proof** Observe first that

$$(13) \quad \lambda(s) = \sum_{m=0}^{\infty} \sum_{k=1}^N \frac{\lambda_k}{(mN + k)^s} = \frac{1}{N^s} \sum_{k=1}^N \lambda_k \zeta\left(s, \frac{k}{N}\right);$$

strictly this is true for  $\sigma > 1$  in the first place, and the equating of the extreme terms for  $\sigma > 0$  follows by analytic continuation.

It follows from (9) that for  $k = 1, 2, \dots, N$ ,

$$-\frac{1}{N^s} \int_{-\infty}^{\infty} e^{-\sigma\omega} \left( e^{\omega} - \left\lceil e^{\omega} - \frac{k}{N} \right\rceil \right) e^{-i\tau\omega} d\omega = \frac{1}{sN^s} \zeta\left(s, \frac{k}{N}\right) - \frac{1}{sk^s}.$$

We now change variables,  $w \mapsto \omega - \log N$ , to obtain

$$-\int_{-\infty}^{\infty} e^{-\sigma w} \left( \frac{e^w}{N} - \left\lceil \frac{e^w - k}{N} \right\rceil \right) e^{-i\tau w} dw = \frac{1}{sN^s} \zeta\left(s, \frac{k}{N}\right) - \frac{1}{sk^s},$$

so that

$$(14) \quad - \int_{-\infty}^{\infty} e^{-\sigma\omega} \left( \frac{e^{\omega}}{N} - \left\lceil \frac{e^{\omega} + N - k}{N} \right\rceil \right) e^{-i\tau\omega} d\omega = \frac{1}{sN^s} \zeta\left(s, \frac{k}{N}\right)$$

since

$$\left\lceil \frac{e^{\omega} + N - k}{N} \right\rceil - \left\lceil \frac{e^{\omega} - k}{N} \right\rceil = \begin{cases} 1 & \text{when } w \geq \log k, \\ 0 & \text{when } w < \log k. \end{cases}$$

Summing (14) for  $k = 1, 2, \dots, N$ , and applying (13) yields (12), as desired. ■

We are now in position to prove the following companion to Theorem 1, in which we use the following notation: given a Dirichlet series  $\lambda(s) := \sum_{n=1}^{\infty} \lambda_n n^{-s}$  with coefficients repeating modulo  $N$ , we define an associated *kernel*

$$W_N(\lambda, t) := \sum_{k=1}^N \lambda_k \left\lceil t + \frac{N - k}{N} \right\rceil.$$

**Theorem 3** Suppose that the coefficients of the two Dirichlet series

$$a(s) := \sum_{n=1}^{\infty} a_n n^{-s} \quad b(s) := \sum_{n=1}^{\infty} b_n n^{-s}$$

repeat modulo  $N$ , that  $A_n := \sum_{k=1}^n a_k, B_n := \sum_{k=1}^n b_k, A_0 := B_0 := 0$ , and that

$$\iota_{a,b}(\sigma) := \frac{1}{2} \int_{-\infty}^{\infty} \frac{a(s)\bar{b}(s)}{\sigma^2 + \tau^2} d\tau$$

with  $s = \sigma + i\tau$ . Suppose further that  $A_N = B_N = 0$ . Then for  $0 < \sigma < 1$ ,

$$(15) \quad \iota_{a,b}(\sigma) = \frac{\pi}{N^{2\sigma}} \int_0^1 \zeta(2\sigma + 1, t) W_N(a, t) W_N(\bar{b}, t) dt.$$

Moreover, for all  $\sigma > 0$ ,

$$(16) \quad \iota_{a,b}(\sigma) = \frac{\pi}{2\sigma} \sum_{n=1}^{\infty} \frac{A_n \bar{B}_n - A_{n-1} \bar{B}_{n-1}}{n^{2\sigma}}.$$

**Proof** Suppose first that  $1 > \sigma > 0$ . Since  $A_N = B_N = 0$ , it follows from Theorem 2 that

$$\frac{a(s)}{s} = \int_{-\infty}^{\infty} e^{-\sigma\omega} \sum_{k=1}^N a_k \left[ \frac{e^{\omega} + N - k}{N} \right] e^{-i\tau\omega} d\omega,$$

with a corresponding formula for  $b(s)/s$ . Hence, by the  $\mathcal{L}_2$  Plancherel theorem [9],

$$\begin{aligned} (17) \quad & \int_{-\infty}^{\infty} \frac{a(s)\bar{b}(s)}{\sigma^2 + \tau^2} d\tau \\ &= 2\pi \int_{-\infty}^{\infty} e^{-2\sigma\omega} \left( \sum_{k=1}^N a_k \left[ \frac{e^{\omega} + N - k}{N} \right] \right) \left( \sum_{k=1}^N \bar{b}_k \left[ \frac{e^{\omega} + N - k}{N} \right] \right) d\omega \\ &= 2\pi \sum_{n=0}^N \int_0^N \frac{1}{(Nn + u)^{2\sigma+1}} \left( \sum_{k=1}^N a_k \left[ \frac{u}{N} + \frac{N - k}{N} \right] \right) \\ & \quad \times \left( \sum_{k=1}^N \bar{b}_k \left[ \frac{u}{N} + \frac{N - k}{N} \right] \right) du \\ &= \frac{2\pi}{N^{2\sigma}} \int_0^1 \zeta(2\sigma + 1, t) \left( \sum_{k=1}^N a_k \left[ t + \frac{N - k}{N} \right] \right) \\ & \quad \times \left( \sum_{k=1}^N \bar{b}_k \left[ t + \frac{N - k}{N} \right] \right) dt \\ &= \frac{2\pi}{N^{2\sigma}} \int_0^1 \zeta(2\sigma + 1, t) W_N(a, t) W_N(\bar{b}, t) dt. \end{aligned}$$

This establishes (15). We now denote the characteristic function of the interval  $(k/N, (k + 1)/N)$  by  $\chi_k$  and observe that since  $A_N = \bar{B}_N = 0$ , we can use summation

by parts to re-express the kernels as follows:

$$W_N(a, t) = \sum_{k=1}^{N-1} A_k \chi_k(t), \quad W_N(\bar{b}, t) = \sum_{k=1}^{N-1} \bar{B}_k \chi_k(t).$$

Thus,

$$W_N(a, t)W_N(\bar{b}, t) = \sum_{k=1}^{N-1} A_k \bar{B}_k \chi_k(t).$$

Consequently, by (15),

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{a(s)\bar{b}(s)}{\sigma^2 + \tau^2} d\tau &= \frac{2\pi}{N^{2\sigma}} \sum_{k=1}^{N-1} A_k \bar{B}_k \int_{\frac{k}{N}}^{\frac{k+1}{N}} \zeta(2\sigma + 1, t) dt \\ &= \frac{\pi}{\sigma} \sum_{n=0}^{\infty} \sum_{k=1}^{N-1} A_k \bar{B}_k \left( \frac{1}{(Nn+k)^{2\sigma}} - \frac{1}{(Nn+k+1)^{2\sigma}} \right) \\ &= \frac{\pi}{\sigma} \sum_{n=0}^{\infty} \sum_{k=1}^N \frac{A_k \bar{B}_k - A_{k-1} \bar{B}_{k-1}}{(Nn+k)^{2\sigma}} = \frac{\pi}{\sigma} \sum_{m=1}^{\infty} \frac{A_m \bar{B}_m - A_{m-1} \bar{B}_{m-1}}{m^{2\sigma}}, \end{aligned}$$

and this shows that (16) holds when  $0 < \sigma < 1$ .

By Theorem 1, (16) also holds when  $\sigma > 1$ , since the Dirichlet series defining  $a(s)$  and  $b(s)$  are absolutely convergent in this range. Further, since  $A_n \bar{B}_n - A_{n-1} \bar{B}_{n-1}$  is bounded, by Dirichlet’s test the series in (16) is absolutely convergent and thus continuous as a function of  $\sigma$  for  $\sigma > 0$ . Finally, it is easy to show by means of partial summation that  $a(s)$  and  $b(s)$  are analytic in the disk  $\{s, |s - 1| < \frac{1}{4}\}$ , and hence bounded therein by a constant  $M$ , say. It follows by Lebesgue’s theorem on dominated convergence that  $\iota_{a,b}(\sigma)$  is continuous for  $\frac{3}{4} < \sigma < \frac{5}{4}$ , and hence that (16) holds for all  $\sigma > 0$ . ■

As an immediate consequence we have the following companion to Corollary 1.

**Corollary 2** *Suppose that the coefficients of the Dirichlet series  $\lambda(s) := \sum_{n=1}^{\infty} \lambda_n n^{-s}$  repeat modulo  $N$ , and that  $\Lambda_n := \sum_{k=1}^n \lambda_k, \Lambda_0 := 0$ . Suppose further that  $\Lambda_N = 0$ . Then for  $s = \sigma + i\tau$  with  $\sigma > 0$  we have*

$$\iota_{\lambda}(\sigma) := \frac{1}{2} \int_{-\infty}^{\infty} \left| \frac{\lambda(s)}{s} \right|^2 d\tau = \frac{\pi}{2\sigma} \sum_{n=1}^{\infty} \frac{|\Lambda_n|^2 - |\Lambda_{n-1}|^2}{n^{2\sigma}},$$

where  $\Lambda_n := \sum_{k=1}^n \lambda_k, \Lambda_0 := 0$ . If, in addition, all the coefficients  $\lambda_n$  are real, then

$$\iota_{\lambda}(2\sigma) = \int_0^{\infty} \left| \frac{\lambda(s)}{s} \right|^2 d\tau = \frac{\pi}{\sigma} \sum_{n=1}^{\infty} \frac{\Lambda_n^2 - \Lambda_{n-1}^2}{n^{2\sigma}}.$$

Note that for  $N = 2$ ,  $\lambda_1 = 1 = -\lambda_2$  we reobtain from Corollary 2, a rigorous form of the original evaluation [8].

**Example 4** Recall, as in [4], that

$$\zeta(\bar{u}, v) := \sum_{n=2}^{\infty} \frac{(-1)^n}{n^u} \sum_{k=1}^{n-1} \frac{1}{k^v},$$

while

$$\zeta(u, \bar{v}) := \sum_{n=2}^{\infty} \frac{1}{n^u} \sum_{k=1}^{n-1} \frac{(-1)^k}{k^v}, \quad \zeta(\bar{u}, \bar{v}) := \sum_{n=2}^{\infty} \frac{(-1)^n}{n^u} \sum_{k=1}^{n-1} \frac{(-1)^k}{k^v}.$$

The same approach as in Example 1 applied to (5) and (16) produces

$$\frac{1}{\pi} \int_0^{\infty} \frac{\alpha(3/2 + i\tau) \overline{\alpha(3/2 + i\tau)}}{1/4 + \tau^2} d\tau = 2\zeta(\bar{2}, \bar{1}) + \zeta(3) = 3\zeta(2) \log 2 - \frac{9}{4}\zeta(3),$$

and

$$\begin{aligned} \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{\alpha(3/2 + i\tau) \overline{\zeta(3/2 + i\tau)}}{1/4 + \tau^2} d\tau &= \zeta(\bar{2}, 1) + \zeta(2, \bar{1}) + \alpha(3) \\ &= \frac{9}{8}\zeta(2) \log 2 - \frac{3}{4}\zeta(3), \end{aligned}$$

as companions to

$$\frac{1}{\pi} \int_0^{\infty} \frac{\zeta(3/2 + i\tau) \overline{\zeta(3/2 + i\tau)}}{1/4 + \tau^2} d\tau = 3\zeta(3),$$

since, by techniques discussed in [4, 5],

$$\zeta(\bar{2}, 1) = \frac{1}{8}\zeta(3), \quad \zeta(2, \bar{1}) = \zeta(3) - \frac{3}{2}\zeta(2) \log 2, \quad \zeta(\bar{2}, \bar{1}) = \frac{3}{2}\zeta(2) \log 2 - \frac{13}{8}\zeta(3).$$

## 4 Final Remarks

Many other similar results obtain. For example:

$$\begin{aligned} \int_{-\infty}^{\infty} e^{-2\sigma\omega} \{\max(e^\omega - a, 0)\}^2 d\omega &= \sum_{n=0}^{\infty} \int_{n+a}^{n+1+a} t^{-2\sigma-1} |t-a|^2 dt \\ &= \sum_{n=0}^{\infty} \int_0^1 \theta^2 (\theta + n + a)^{-2\sigma-1} d\theta \\ &= \int_0^1 \theta^2 \zeta(1 + 2\sigma, \theta + a) d\theta. \end{aligned}$$

While (4) and (16) give an effective way of evaluating the integral, directly evaluating the integral numerically to high precision presents a greater challenge. This is largely because of the severe oscillations of the integrand. The issue appears to lie in estimating the integrand well and so is intrinsically non-trivial.

**Acknowledgements** We should like to thank David Bradley for the many useful suggestions he made during the completion of this research.

## References

- [1] T. M. Apostol, *Introduction to Analytic Number Theory*. Springer-Verlag, New York, 1976.
- [2] J. Borwein. *A class of Dirichlet series integrals*. Amer. Math. Monthly **114**(2007), 70–76.
- [3] J. Borwein and D. Bailey, *Mathematics by Experiment. Plausible Reasoning in the 21st Century*. AK Peters, Natick, MA, 2004.
- [4] J. M. Borwein and D. M. Bradley, *Thirty-two Goldbach variations*. Int. J. Number Theory **2**(2006), no. 1, 65–103.
- [5] J. Borwein, D. Bailey, and R. Girgensohn *Experimentation in Mathematics. Computational Paths to Discovery*. AK Peters, Natick, MA, 2004.
- [6] J. M. Borwein, D. M. Bailey, and R. E. Crandall, *Computational strategies for the Riemann zeta function*. J. Comput. Appl. Math. **121**(2000), no. 1-2, 247–296.
- [7] G. H. Hardy and E. M. Wright, *Introduction to the Theory of Numbers*. Oxford University Press, Oxford, 1985.
- [8] Solution to AMM Problem 10939, proposed by Ivić, with his solution. Amer. Math. Monthly **110**(2003), no. 9, 847–848.
- [9] K. R. Stromberg *Introduction to Classical Real Analysis*. Wadsworth, Belmont, CA, 1981.

*Department of Mathematics*  
*University of Western Ontario*  
*London, ON*  
*N6A 5B7*  
*e-mail: dborwein@uwo.ca*

*Faculty of Computer Science*  
*Dalhousie University*  
*Halifax, NS*  
*B3H 1W5*  
*e-mail: jborwein@cs.dal.ca*