## ON THE GAUSS MAP OF RULED SURFACES

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(Received 25 June, 1991)

Let  $M^2$  be a (connected) surface in Euclidean 3-space  $E^3$ , and let  $G: M^2 \rightarrow S^2(1) \subset E^3$ be its Gauss map. Then, according to a theorem of E. A. Ruh and J. Vilms [3],  $M^2$  is a surface of constant mean curvature if and only if, as a map from  $M^2$  to  $S^2(1)$ , G is harmonic, or equivalently, if and only if

$$\Delta G = \|dG\|^2 G \tag{1.1}$$

where  $\Delta$  is the Laplace operator on  $M^2$  corresponding to the induced metric on  $M^2$  from  $E^3$  and where G is seen as a map from  $M^2$  to  $E^3$ . A special case of (1.1) is given by

$$\Delta G = \lambda G, \qquad (\lambda \in \mathbb{R}) \tag{1.2}$$

i.e., the case where the Gauss map  $G: M^2 \to E^3$  is an eigenfunction of the Laplacian  $\Delta$  on  $M^2$ .

On the other hand, F. Dillen, J. Pas and L. Verstraelen [2] recently proved that among the surfaces of revolution in  $E^3$ , the only ones whose Gauss map satisfies the condition

$$\Delta G = \Lambda G, \qquad (\Lambda \in \mathbb{R}^{3 \times 3}) \tag{1.3}$$

are the planes, the spheres and the circular cylinders.

We observe that from the surfaces of revolution in  $E^3$  which satisfy (1.3) the planes and the circular cylinders are ruled surfaces. On the other hand, for the helicoid  $X(s,t) = (t \cos s, t \sin s, \alpha s), \alpha \neq 0$  the Guass map is given by

$$G = \frac{1}{\sqrt{t^2 + \alpha^2}} (-\alpha \sin s, \alpha \cos s, -t).$$

Then it is easy to show that the Laplacian  $\Delta G$  of the Gauss map G can be expressed as follows

$$\Delta G = \frac{2\alpha^2}{\left(t^2 - \alpha^2\right)^{5/2}} \left(-\alpha \sin s, \, \alpha \cos s, \, -t\right)$$

which clearly doesn't satisfy condition (1.3).

A question which arises now is: What are the ruled surfaces satisfying condition (1.3)? In particular, we will prove the following:

THEOREM. Among the ruled surfaces in  $E^3$ , the only ones whose Gauss map satisfies (1.3) are the planes and the circular cylinders.

We first study cylindrical surfaces  $M^2$ . Let  $X(s, t) = \alpha(s) + t\beta$  be the position vector of  $M^2$  in  $E^3$  where  $\alpha(s)$  is the plane curve  $\alpha = (\alpha_1, \alpha_2, 0)$  parameterized by arc-length and  $\beta$  is the constant vector  $\beta = (0, 0, 1)$ . We have the following lemma.

LEMMA. The only cylindrical surfaces whose Gauss map satisfies (1.3) are the planes and the circular cylinders.

† This work was done while the first author was a visiting scholar at Michigan State University.

Glasgow Math. J. 34 (1992) 355-359.

*Proof.* The Gauss map of  $M^2$  is  $G = \alpha' \times \beta = (\alpha'_2, -\alpha'_1, 0)$  and the Laplacian of G is  $\Delta G = (-\alpha''_2, \alpha''_1, 0)$ . Thus from the condition (1.3) we have

(i) 
$$-\alpha'''_{2} = \lambda_{11}\alpha'_{2} - \lambda_{12}\alpha'_{1}$$
  
(ii)  $\alpha'''_{1} = \lambda_{21}\alpha'_{2} - \lambda_{22}\alpha'_{1}$   
(iii)  $0 = \lambda_{31}\alpha'_{2} - \lambda_{32}\alpha'_{1}$ 

where  $\Lambda = [\lambda_{ii}]$  is a constant matrix. Since  $|\alpha'| = 1$  we can put

$$\alpha'_1 = \cos \theta, \qquad \alpha'_2 = \sin \theta \tag{2.2}$$

where  $\theta = \theta(s)$ . Then from (2.1)(i), (ii) we obtain

$$\theta'' \cos \theta - \theta'^2 \sin \theta = -\lambda_{11} \sin \theta + \lambda_{12} \cos \theta$$
$$\theta'' \sin \theta + \theta'^2 \cos \theta = -\lambda_{21} \sin \theta + \lambda_{22} \cos \theta$$

which give

$$\theta'^{2} = -(\lambda_{12} + \lambda_{21})\sin\theta\cos\theta + \lambda_{11}\sin^{2}\theta + \lambda_{22}\cos^{2}\theta$$
(2.3)

$$\theta'' = (\lambda_{22} - \lambda_{11})\sin\theta\cos\theta + \lambda_{12}\cos^2\theta - \lambda_{21}\sin^2\theta.$$
(2.4)

Taking the derivative of (2.3) and using (2.4) we obtain

$$\theta'[4(\lambda_{22} - \lambda_{11})\sin\theta\cos\theta + (3\lambda_{12} + \lambda_{21})\cos^2\theta - (\lambda_{12} + 3\lambda_{21})\sin^2\theta] = 0$$

If  $\theta' = 0$ , the Gauss map G is constant and hence  $M^2$  is a plane. So, suppose  $\theta' \neq 0$ . Since  $\sin^2 \theta$ ,  $\cos^2 \theta$  and  $\sin \theta \cos \theta$  are linearly independent functions of  $\theta = \theta(s)$ , we obtain from (2.5).

$$\lambda_{11} = \lambda_{22}, \, 3\lambda_{12} + \lambda_{21} = 0, \, \lambda_{12} + 3\lambda_{21} = 0.$$

Thus  $\lambda_{12} = \lambda_{21} = 0$ . Substitution into (2.3) then gives  $\theta'^2 = \frac{1}{r^2}$ , where  $\frac{1}{r^2} = \lambda_{11} = \lambda_{22} = \text{const.}$ 

Now from (2.1)(i) and (ii) we conclude that the curve  $\alpha$  is the circle

$$\alpha = (r\sin(rs + c) + d_1, -r\cos(rs + c) + d_2, 0)$$
(2.6)

where  $c, d_1$ , and  $d_2$  are constants. Also from (2.1)(iii) we obtain  $\lambda_{31} = \lambda_{32} = 0$ .

REMARK. The matrix  $\Lambda = [\lambda_{ij}]$  in the condition (1.3) when  $M^2$  is the circular cylinder on the circle (2.6) is given by

$$\Lambda = \begin{bmatrix} \frac{1}{r^2} & 0 & \lambda_{13} \\ 0 & \frac{1}{r^2} & \lambda_{23} \\ 0 & 0 & \lambda_{33} \end{bmatrix}$$

where  $\lambda_{i3}$ , i = 1, 2, 3 are arbitrary constants.

Proof of the theorem. We suppose that  $M^2$  is a non-cylindrical ruled surface in  $E^3$ . The surface  $M^2$  can be expressed in terms of a directrix curve  $\alpha(s)$  and a unit vector field  $\beta(s)$  pointing along the rulings as

$$X(s,t) = \alpha(s) + t\beta(s).$$

Moreover, we can choose the parameter s to be arc length along the spherical curve  $\beta(s)$ . Thus for the curves  $\alpha$ ,  $\beta$  we have

$$\langle \beta, \beta \rangle = 1, \langle \beta', \beta' \rangle = 1, \langle \alpha', \beta \rangle = 0.$$
 (2.7)

If we define a function q by

$$q = \|\alpha' + t\beta'\|^2 = t^2 + 2ut + v$$
(2.8)

where  $u = \langle \alpha', \beta' \rangle$  and  $v = \langle \alpha', \alpha' \rangle$ , then the Gauss map of the surface is given by

$$G = q^{-1/2}((\alpha' + t\beta') \times \beta).$$

It is easy to show that the Laplacian  $\Delta$  of M can be expressed as (see [1])

$$\Delta = -\frac{\partial^2}{\partial t^2} - \frac{1}{q} \frac{\partial^2}{\partial s^2} + \frac{1}{2} \frac{\partial q}{\partial s} \frac{1}{q^2} \frac{\partial}{\partial s} - \frac{1}{2} \frac{\partial q}{\partial t} \frac{1}{q} \frac{\partial}{\partial t}.$$
 (2.10)

For convenience we put

$$G = (G_1, G_2, G_3) = q^{-1/2} (A_1 + tB_1, A_2 + tB_2, A_2 + tB_3)$$
(2.11)

where

$$(A_1, A_2, A_3) = \alpha' \times \beta$$
  

$$(B_1, B_2, B_3) = \beta' \times \beta.$$
(2.12)

We now compute the Laplacian of the functions  $G_i$ . We have

$$\begin{aligned} \frac{\partial G_i}{\partial t} &= q^{-3/2} [B_i q - (A_i + tB_i)(t+u)] = q^{-3/2} C_i \\ \frac{\partial^2 G_i}{\partial t^2} &= q^{-5/2} [(B_i u - A_i)q - 3(B_i q - (A_i + tB_i)(t+u))(t+u)] = q^{-5/2} D_i \\ \frac{\partial G_i}{\partial s} &= \frac{1}{2} q^{-3/2} [2(A_i' + tB_i')q - (A_i + tB_i)(2u't+v')] = \frac{1}{2} q^{-3/2} E_i \\ \frac{\partial^2 G_i}{\partial s^2} &= \frac{1}{2} q^{-5/2} \{ [2(A_i'' + tB_i')q + (A_i' + tB_i')(2u't+v') - (A_i + tB_i)(2u''t+v'')] q \\ &\quad -\frac{3}{2} [2(A_i' + tB_i')q - (A_i + tB_i)(2u't+v')] (2u't+v') \} \\ &= \frac{1}{2} q^{-5/2} F_i. \end{aligned}$$

Thus, from the above relations and (2.10),

$$\Delta G_i = -q^{-5/2} D_i - \frac{1}{2} q^{-7/2} F_i + \frac{1}{4} q^{-7/2} (2u't + v') E_i - q^{-5/2} (t+u) C_i.$$

Now if we put  $\Lambda = [\lambda_{ij}]$  from (1.3) and (2.11) we have

$$-4qD_i - 2F_i + (2u't + v')E_i - 4q(t+u)C_i = 4\sum_{j=1}^3 \lambda_{ij}(A_j + tB_j)q^3, \qquad i = 1, 2, 3.$$
(2.13)

We consider the powers of t in equation (2.13). From the coefficient of  $t^7$  we have

$$\sum_{j=1}^{3} \lambda_{ij} B_j = 0, \qquad i = 1, 2, 3.$$
(2.14)

Considering the coefficients of the other powers of t and using (2.14) we obtain for any i = 1, 2, 3

$$\sum_{i=1}^{3} \lambda_{ij} A_j = 0$$
 (2.15)

$$B_i'' = 0$$
 (2.16)

$$A_i - B_i u - 3B'_i u' + A''_i - B_i u'' = 0$$
(2.17)

$$-8A_{i}u + 4B_{i}u^{2} + 4B_{i}v - 8A_{i}'u + 6A_{i}'u' + 12B_{i}'uu' + 3B_{i}'v' + 2A_{i}u'' + B_{i}v'' + 4B_{i}uu'' - 8B_{i}u'^{2} = 0 \quad (2.18)$$

$$12B_{i}uv - 12A_{i}u^{2} - 4A_{i}'v + 3A_{i}'v' + A_{i}v'' - 8A_{i}'u^{2} + 12A_{i}'uu' + 6B_{i}'u'v + 6B_{i}'uv' + 4A_{i}uu'' + 2B_{i}uv'' + 2B_{i}u''v$$

$$-8A_{i}u'^{2}-8B_{i}u'v'=0 \quad (2.19)$$

$$4B_{i}u^{2}v + 4B_{i}v^{2} - 8A_{i}u^{3} - 8A_{i}''uv + 6A_{i}'uv' + 2A_{i}uv'' + 6A_{i}'u'v + 3B_{i}'vv' + 2A_{i}u''v + B_{i}vv'' - 8A_{i}u'v' - 2B_{i}v'^{2} = 0 \quad (2.20)$$

$$2B_{i}uv^{2} + 2A_{i}v^{2} - 4A_{i}u^{2}v - 2A_{i}'v^{2} + 3A_{i}'vv' + A_{i}vv'' - 2A_{i}v'^{2} = 0.$$
(2.21)

We remark that det  $\Lambda = 0$ , for if we assume det  $\Lambda \neq 0$ , then from (2.14),  $B_i = 0$ , i = 1, 2, 3. Thus, from (2.12) we have  $\beta' \times \beta = 0$ , contradicting (2.7).

From (2.16) and (2.12) we have  $\beta' \times \beta = cs + d$ , where c and d are constant vectors. So  $1 = ||\beta' \times \beta||^2 = \langle c, c \rangle s^2 + 2\langle c, d \rangle s + \langle d, d \rangle$ , from which we conclude that  $\langle c, c \rangle = 0$ ,  $\langle d, d \rangle = 1$ , or equivalently  $\beta' \times \beta = d$ , where d is a constant unit vector. Since  $\beta$  is a spherical curve, this implies that  $\beta$  is a great circle. Let  $\beta = (\sin \theta \cos \varphi, \sin \theta \sin \varphi, \cos \theta)$ , where  $\theta = \theta(s)$  and  $\varphi = \text{const.}$  From  $\beta' \times \beta = d$  we conclude that  $\theta'^2 = 1$  and so

$$\beta' \times \beta = (B_1, B_2, B_3) = (\sin \varphi, -\cos \varphi, 0).$$

Now, from (2.17) we have  $A_i - B_i u + (A_i - B_i u)'' = 0$ . If we put  $A_i - B_i u = w_i$ , i = 1, 2, 3, then  $w_i + w_i'' = 0$  and  $(\alpha' - u\beta') \times \beta = w$  where  $w = (w_1, w_2, w_3)$ . So  $||\alpha' - u\beta'||^2 = \langle w, w \rangle$ , or

$$v = u^2 + w^2 \tag{2.22}$$

where  $w^2 = \langle w, w \rangle$ .

Since  $A_3 = w_3$  and  $w_3'' = -w_3$ , we have from (2.18)

$$3w_3'u' + w_3u'' = 0. \tag{2.23}$$

By using (2.23), from (2.19) we find that

$$-4w_3u^2 + 4w_3v + 3w'_3v' + w_3v'' - 8w_3u'^2 = 0.$$
 (2.24)

Using (2.22), (2.23) and (2.24), equations (2.20) and (2.21) can be written as

$$w_3(w^2)'u' = 0 \tag{2.25}$$

$$4w_3u'^2v - w_3v'^2 = 0. (2.26)$$

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Now, using the equations (2.22)–(2.26) we will prove that w = 0. Suppose, for the moment, that  $w_3 \neq 0$ . From (2.25) we have  $(w^2)'u' = 0$ . If u' = 0, from (2.26) we have v' = 0 and hence from (2.24)  $u^2 = v$ . Thus (2.22) implies w = 0, a contradiction. Thus  $u' \neq 0$  and so  $(w^2)' = 0$ . From (2.22) v' = 2uu' and from (2.26)  $v = u^2$ . Again (2.22) implies w = 0, a contradiction. So, we have  $w_3 = 0$ . This means that the vector w lies in the xy plane. But  $w = \alpha' \times \beta - u\beta' \times \beta$  and the vector  $\beta' \times \beta$  lies in the xy plane. So  $\alpha' \times \beta$  lies in the xy plane. This means that the vectors  $\alpha' \times \beta$  and  $\beta' \times \beta$  are parallel. If we put  $\alpha' \times \beta = \mu\beta' \times \beta$ , then  $(\alpha' - \mu\beta') \times \beta = 0$  or  $\alpha' = \mu\beta'$ . So  $\mu = \langle \alpha', \beta' \rangle = u$  and  $\alpha' = u\beta'$ , namely w = 0.

Now we conclude that  $q = (t + u)^2$  and the Gauss map is constant, which means that  $M^2$  is a plane.

## REFERENCES

1. B. Y. Chen, F. Dillen, L. Verstraelen and L. Vrancken, Ruled surfaces of finite type, Bull. Austral. Math. Soc. 42 (1990), 447-453.

2. F. Dillen, J. Pas and L. Verstraelen, On the Gauss map of surfaces of revolution, Bull. Inst. Math. Acad. Sinica, 18 (1990), 239-246.

3. E. A. Ruh and J. Vilms, The tension field of the Gauss map, Trans. Amer. Math. Soc. 149 (1970), 569-573.

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