# ON THE GAUSS MAP OF RULED SURFACES <br> by CHRISTOS BAIKOUSSIS $\dagger$ and DAVID E. BLAIR 

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Let $M^{2}$ be a (connected) surface in Euclidean 3-space $E^{3}$, and let $G: M^{2} \rightarrow S^{2}(1) \subset E^{3}$ be its Gauss map. Then, according to a theorem of E. A. Ruh and J. Vilms [3], $M^{2}$ is a surface of constant mean curvature if and only if, as a map from $M^{2}$ to $S^{2}(1), G$ is harmonic, or equivalently, if and only if

$$
\begin{equation*}
\Delta G=\|d G\|^{2} G \tag{1.1}
\end{equation*}
$$

where $\Delta$ is the Laplace operator on $M^{2}$ corresponding to the induced metric on $M^{2}$ from $E^{3}$ and where $G$ is seen as a map from $M^{2}$ to $E^{3}$. A special case of (1.1) is given by

$$
\begin{equation*}
\Delta G=\lambda G, \quad(\lambda \in \mathbb{R}) \tag{1.2}
\end{equation*}
$$

i.e., the case where the Gauss map $G: M^{2} \rightarrow E^{3}$ is an eigenfunction of the Laplacian $\Delta$ on $M^{2}$.

On the other hand, F. Dillen, J. Pas and L. Verstraelen [2] recently proved that among the surfaces of revolution in $E^{3}$, the only ones whose Gauss map satisfies the condition

$$
\begin{equation*}
\Delta G=\Lambda G, \quad\left(\Lambda \in \mathbb{R}^{3 \times 3}\right) \tag{1.3}
\end{equation*}
$$

are the planes, the spheres and the circular cylinders.
We observe that from the surfaces of revolution in $E^{3}$ which satisfy (1.3) the planes and the circular cylinders are ruled surfaces. On the other hand, for the helicoid $X(s, t)=(t \cos s, t \sin s, \alpha s), \alpha \neq 0$ the Guass map is given by

$$
G=\frac{1}{\sqrt{t^{2}+\alpha^{2}}}(-\alpha \sin s, \alpha \cos s,-t) .
$$

Then it is easy to show that the Laplacian $\Delta G$ of the Gauss map $G$ can be expressed as follows

$$
\Delta G=\frac{2 \alpha^{2}}{\left(t^{2}-\alpha^{2}\right)^{5 / 2}}(-\alpha \sin s, \alpha \cos s,-t)
$$

which clearly doesn't satisfy condition (1.3).
A question which arises now is: What are the ruled surfaces satisfying condition (1.3)?
In particular, we will prove the following:
Theorem. Among the ruled surfaces in $E^{3}$, the only ones whose Gauss map satisfies (1.3) are the planes and the circular cylinders.

We first study cylindrical surfaces $M^{2}$. Let $X(s, t)=\alpha(s)+t \beta$ be the position vector of $M^{2}$ in $E^{3}$ where $\alpha(s)$ is the plane curve $\alpha=\left(\alpha_{1}, \alpha_{2}, 0\right)$ parameterized by arc-length and $\beta$ is the constant vector $\beta=(0,0,1)$. We have the following lemma.

Lemma. The only cylindrical surfaces whose Gauss map satisfies (1.3) are the planes and the circular cylinders.
$\dagger$ This work was done while the first author was a visiting scholar at Michigan State University.
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Proof. The Gauss map of $M^{2}$ is $G=\alpha^{\prime} \times \beta=\left(\alpha_{2}^{\prime},-\alpha_{1}^{\prime}, 0\right)$ and the Laplacian of $G$ is $\Delta G=\left(-\alpha_{2}^{\prime \prime \prime}, \alpha_{1}^{\prime \prime \prime}, 0\right)$. Thus from the condition (1.3) we have
(i) $-\alpha_{2}^{\prime \prime \prime}=\lambda_{11} \alpha_{2}^{\prime}-\lambda_{12} \alpha_{1}^{\prime}$
(ii) $\alpha_{1}^{\prime \prime \prime}=\lambda_{21} \alpha_{2}^{\prime}-\lambda_{22} \alpha_{1}^{\prime}$
(iii) $0=\lambda_{31} \alpha_{2}^{\prime}-\lambda_{32} \alpha_{1}^{\prime}$
where $\Lambda=\left[\lambda_{i j}\right]$ is a constant matrix. Since $\left|\alpha^{\prime}\right|=1$ we can put

$$
\begin{equation*}
\alpha_{1}^{\prime}=\cos \theta, \quad \alpha_{2}^{\prime}=\sin \theta \tag{2.2}
\end{equation*}
$$

where $\theta=\theta(s)$. Then from (2.1)(i), (ii) we obtain

$$
\begin{aligned}
& \theta^{\prime \prime} \cos \theta-\theta^{\prime 2} \sin \theta=-\lambda_{11} \sin \theta+\lambda_{12} \cos \theta \\
& \theta^{\prime \prime} \sin \theta+\theta^{\prime 2} \cos \theta=-\lambda_{21} \sin \theta+\lambda_{22} \cos \theta
\end{aligned}
$$

which give

$$
\begin{align*}
\theta^{\prime 2} & =-\left(\lambda_{12}+\lambda_{21}\right) \sin \theta \cos \theta+\lambda_{11} \sin ^{2} \theta+\lambda_{22} \cos ^{2} \theta  \tag{2.3}\\
\theta^{\prime \prime} & =\left(\lambda_{22}-\lambda_{11}\right) \sin \theta \cos \theta+\lambda_{12} \cos ^{2} \theta-\lambda_{21} \sin ^{2} \theta \tag{2.4}
\end{align*}
$$

Taking the derivative of (2.3) and using (2.4) we obtain

$$
\theta^{\prime}\left[4\left(\lambda_{22}-\lambda_{11}\right) \sin \theta \cos \theta+\left(3 \lambda_{12}+\lambda_{21}\right) \cos ^{2} \theta-\left(\lambda_{12}+3 \lambda_{21}\right) \sin ^{2} \theta\right]=0
$$

If $\theta^{\prime}=0$, the Gauss map $G$ is constant and hence $M^{2}$ is a plane. So, suppose $\theta^{\prime} \neq 0$. Since $\sin ^{2} \theta, \cos ^{2} \theta$ and $\sin \theta \cos \theta$ are linearly independent functions of $\theta=\theta(s)$, we obtain from (2.5).

$$
\lambda_{11}=\lambda_{22}, 3 \lambda_{12}+\lambda_{21}=0, \lambda_{12}+3 \lambda_{21}=0 .
$$

Thus $\lambda_{12}=\lambda_{21}=0$. Substitution into (2.3) then gives $\theta^{\prime 2}=\frac{1}{r^{2}}$, where $\frac{1}{r^{2}}=\lambda_{11}=\lambda_{22}=$ const. Now from (2.1)(i) and (ii) we conclude that the curve $\alpha$ is the circle

$$
\begin{equation*}
\alpha=\left(r \sin (r s+c)+d_{1},-r \cos (r s+c)+d_{2}, 0\right) \tag{2.6}
\end{equation*}
$$

where $c, d_{1}$, and $d_{2}$ are constants. Also from (2.1)(iii) we obtain $\lambda_{31}=\lambda_{32}=0$.
Remark. The matrix $\Lambda=\left[\lambda_{i j}\right]$ in the condition (1.3) when $M^{2}$ is the circular cylinder on the circle (2.6) is given by

$$
\Lambda=\left[\begin{array}{ccc}
\frac{1}{r^{2}} & 0 & \lambda_{13} \\
0 & \frac{1}{r^{2}} & \lambda_{23} \\
0 & 0 & \lambda_{33}
\end{array}\right]
$$

where $\lambda_{i 3}, i=1,2,3$ are arbitrary constants.
Proof of the theorem. We suppose that $M^{2}$ is a non-cylindrical ruled surface in $E^{3}$. The surface $M^{2}$ can be expressed in terms of a directrix curve $\alpha(s)$ and a unit vector field $\beta(s)$ pointing along the rulings as

$$
X(s, t)=\alpha(s)+t \beta(s)
$$

Moreover, we can choose the parameter $s$ to be arc length along the spherical curve $\beta(s)$. Thus for the curves $\alpha, \beta$ we have

$$
\begin{equation*}
\langle\beta, \beta\rangle=1,\left\langle\beta^{\prime}, \beta^{\prime}\right\rangle=1,\left\langle\alpha^{\prime}, \beta\right\rangle=0 \tag{2.7}
\end{equation*}
$$

If we define a function $q$ by

$$
\begin{equation*}
q=\left\|\alpha^{\prime}+t \beta^{\prime}\right\|^{2}=t^{2}+2 u t+v \tag{2.8}
\end{equation*}
$$

where $u=\left\langle\alpha^{\prime}, \beta^{\prime}\right\rangle$ and $v=\left\langle\alpha^{\prime}, \alpha^{\prime}\right\rangle$, then the Gauss map of the surface is given by

$$
G=q^{-1 / 2}\left(\left(\alpha^{\prime}+t \beta^{\prime}\right) \times \beta\right)
$$

It is easy to shwo that the Laplacian $\Delta$ of $M$ can be expressed as (see [1])

$$
\begin{equation*}
\Delta=-\frac{\partial^{2}}{\partial t^{2}}-\frac{1}{q} \frac{\partial^{2}}{\partial s^{2}}+\frac{1}{2} \frac{\partial q}{\partial s} \frac{1}{q^{2}} \frac{\partial}{\partial s}-\frac{1}{2} \frac{\partial q}{\partial t} \frac{1}{q} \frac{\partial}{\partial t} . \tag{2.10}
\end{equation*}
$$

For convenience we put

$$
\begin{equation*}
G=\left(G_{1}, G_{2}, G_{3}\right)=q^{-1 / 2}\left(A_{1}+t B_{1}, A_{2}+t B_{2}, A_{2}+t B_{3}\right) \tag{2.11}
\end{equation*}
$$

where

$$
\begin{align*}
& \left(A_{1}, A_{2}, A_{3}\right)=\alpha^{\prime} \times \beta  \tag{2.12}\\
& \left(B_{1}, B_{2}, B_{3}\right)=\beta^{\prime} \times \beta .
\end{align*}
$$

We now compute the Laplacian of the functions $G_{i}$. We have

$$
\begin{aligned}
\frac{\partial G_{i}}{\partial t}= & q^{-3 / 2}\left[B_{i} q-\left(A_{i}+t B_{i}\right)(t+u)\right]=q^{-3 / 2} C_{i} \\
\frac{\partial^{2} G_{i}}{\partial t^{2}}= & q^{-5 / 2}\left[\left(B_{i} u-A_{i}\right) q-3\left(B_{i} q-\left(A_{i}+t B_{i}\right)(t+u)\right)(t+u)\right]=q^{-5 / 2} D_{i} \\
\frac{\partial G_{i}}{\partial s}= & \frac{1}{2} q^{-3 / 2}\left[2\left(A_{i}^{\prime}+t B_{i}^{\prime}\right) q-\left(A_{i}+t B_{i}\right)\left(2 u^{\prime} t+v^{\prime}\right)\right]=\frac{1}{2} q^{-3 / 2} E_{i} \\
\frac{\partial^{2} G_{i}}{\partial s^{2}}= & \frac{1}{2} q^{-5 / 2}\left\{\left[2\left(A_{i}^{\prime \prime}+t B_{i}^{\prime}\right) q+\left(A_{i}^{\prime}+t B_{i}^{\prime}\right)\left(2 u^{\prime} t+v^{\prime}\right)-\left(A_{i}+t B_{i}\right)\left(2 u^{\prime \prime} t+v^{\prime \prime}\right)\right] q\right. \\
& \left.-\frac{3}{2}\left[2\left(A_{i}^{\prime}+t B_{i}^{\prime}\right) q-\left(A_{i}+t B_{i}\right)\left(2 u^{\prime} t+v^{\prime}\right)\right]\left(2 u^{\prime} t+v^{\prime}\right)\right\} \\
= & \frac{1}{2} q^{-5 / 2} F_{i} .
\end{aligned}
$$

Thus, from the above relations and (2.10),

$$
\Delta G_{i}=-q^{-5 / 2} D_{i}-\frac{1}{2} q^{-7 / 2} F_{i}+\frac{1}{4} q^{-7 / 2}\left(2 u^{\prime} t+v^{\prime}\right) E_{i}-q^{-5 / 2}(t+u) C_{i}
$$

Now if we put $\Lambda=\left[\lambda_{i j}\right]$ from (1.3) and (2.11) we have

$$
\begin{equation*}
-4 q D_{i}-2 F_{i}+\left(2 u^{\prime} t+v^{\prime}\right) E_{i}-4 q(t+u) C_{i}=4 \sum_{j=1}^{3} \lambda_{i j}\left(A_{j}+t B_{j}\right) q^{3}, \quad i=1,2,3 \tag{2.13}
\end{equation*}
$$

We consider the powers of $t$ in equation (2.13). From the coefficient of $t^{7}$ we have

$$
\begin{equation*}
\sum_{j=1}^{3} \lambda_{i j} B_{j}=0, \quad i=1,2,3 \tag{2.14}
\end{equation*}
$$

Considering the coefficients of the other powers of $t$ and using (2.14) we obtain for any $i=1,2,3$

$$
\begin{gather*}
\sum_{j=1}^{3} \lambda_{i j} A_{j}=0  \tag{2.15}\\
B_{i}^{\prime \prime}=0  \tag{2.16}\\
-8 A_{i} u+4 B_{i} u^{2}+4 B_{i} v-8 A_{i}^{\prime \prime} u+6 A_{i}^{\prime} u^{\prime}+12 B_{i}^{\prime} u u^{\prime}+3 B_{i}^{\prime} v^{\prime}  \tag{2.17}\\
+2 A_{i} u^{\prime \prime}+B_{i} v^{\prime \prime}+4 B_{i} u u^{\prime \prime}-8 B_{i} u^{\prime 2}=0
\end{gather*}
$$

$$
\begin{align*}
& 12 B_{i} u v-12 A_{i} u^{2}-4 A_{i}^{\prime \prime} v+3 A_{i}^{\prime} v^{\prime}+A_{i} v^{\prime \prime}-8 A_{i}^{\prime \prime} u^{2}+12 A_{i}^{\prime} u u^{\prime} \\
& \\
& +6 B_{i}^{\prime} u^{\prime} v+6 B_{i}^{\prime} u v^{\prime}+4 A_{i} u u^{\prime \prime}+2 B_{i} u v^{\prime \prime}+2 B_{i} u^{\prime \prime} v  \tag{2.19}\\
& \\
& \quad-8 A_{i} u^{\prime 2}-8 B_{i} u^{\prime} v^{\prime}=0
\end{align*}
$$

$$
\begin{align*}
& 4 B_{i} u^{2} v+4 B_{i} v^{2}-8 A_{i} u^{3}-8 A_{i}^{\prime \prime} u v+6 A_{i}^{\prime} u v^{\prime}+2 A_{i} u v^{\prime \prime}+6 A_{i}^{\prime} u^{\prime} v \\
&+3 B_{i}^{\prime} v v^{\prime}+2 A_{i} u^{\prime \prime} v+B_{i} v v^{\prime \prime}-8 A_{i} u^{\prime} v^{\prime}-2 B_{i} v^{\prime 2}=0  \tag{2.20}\\
& 2 B_{i} u v^{2}+2 A_{i} v^{2}-4 A_{i} u^{2} v-2 A_{i}^{\prime \prime} v^{2}+3 A_{i}^{\prime} v v^{\prime}+A_{i} v v^{\prime \prime}-2 A_{i} v^{\prime 2}=0 \tag{2.21}
\end{align*}
$$

We remark that $\operatorname{det} \Lambda=0$, for if we assume $\operatorname{det} \Lambda \neq 0$, then from (2.14), $B_{i}=0$, $i=1,2,3$. Thus, from (2.12) we have $\beta^{\prime} \times \beta=0$, contradicting (2.7).

From (2.16) and (2.12) we have $\beta^{\prime} \times \beta=c s+d$, where $c$ and $d$ are constant vectors. So $1=\left\|\beta^{\prime} \times \beta\right\|^{2}=\langle c, c\rangle s^{2}+2\langle c, d\rangle s+\langle d, d\rangle$, from which we conclude that $\langle c, c\rangle=0$, $\langle d, d\rangle=1$, or equivalently $\beta^{\prime} \times \beta=d$, where $d$ is a constant unit vector. Since $\beta$ is a spherical curve, this implies that $\beta$ is a great circle. Let $\beta=$ $(\sin \theta \cos \varphi, \sin \theta \sin \varphi, \cos \theta)$, where $\theta=\theta(s)$ and $\varphi=$ const. From $\beta^{\prime} \times \beta=d$ we conclude that $\theta^{\prime 2}=1$ and so

$$
\beta^{\prime} \times \beta=\left(B_{1}, B_{2}, B_{3}\right)=(\sin \varphi,-\cos \varphi, 0) .
$$

Now, from (2.17) we have $A_{i}-B_{i} u+\left(A_{i}-B_{i} u\right)^{\prime \prime}=0$. If we put $A_{i}-B_{i} u=w_{i}, i=1,2,3$, then $w_{i}+w_{i}^{\prime \prime}=0$ and $\left(\alpha^{\prime}-u \beta^{\prime}\right) \times \beta=w$ where $w=\left(w_{1}, w_{2}, w_{3}\right)$. So $\left\|\alpha^{\prime}-u \beta^{\prime}\right\|^{2}=$ $\langle w, w\rangle$, or

$$
\begin{equation*}
v=u^{2}+w^{2} \tag{2.22}
\end{equation*}
$$

where $w^{2}=\langle w, w\rangle$.
Since $A_{3}=w_{3}$ and $w_{3}^{\prime \prime}=-w_{3}$, we have from (2.18)

$$
\begin{equation*}
3 w_{3}^{\prime} u^{\prime}+w_{3} u^{\prime \prime}=0 . \tag{2.23}
\end{equation*}
$$

By using (2.23), from (2.19) we find that

$$
\begin{equation*}
-4 w_{3} u^{2}+4 w_{3} v+3 w_{3}^{\prime} v^{\prime}+w_{3} v^{\prime \prime}-8 w_{3} u^{\prime 2}=0 . \tag{2.24}
\end{equation*}
$$

Using (2.22), (2.23) and (2.24), equations (2.20) and (2.21) can be written as

$$
\begin{gather*}
w_{3}\left(w^{2}\right)^{\prime} u^{\prime}=0  \tag{2.25}\\
4 w_{3} u^{\prime 2} v-w_{3} v^{\prime 2}=0 . \tag{2.26}
\end{gather*}
$$

Now, using the equations (2.22)-(2.26) we will prove that $w=0$. Suppose, for the moment, that $w_{3} \neq 0$. From (2.25) we have $\left(w^{2}\right)^{\prime} u^{\prime}=0$. If $u^{\prime}=0$, from (2.26) we have $v^{\prime}=0$ and hence from (2.24) $u^{2}=v$. Thus (2.22) implies $w=0$, a contradiction. Thus $u^{\prime} \neq 0$ and so $\left(w^{2}\right)^{\prime}=0$. From (2.22) $v^{\prime}=2 u u^{\prime}$ and from (2.26) $v=u^{2}$. Again (2.22) implies $w=0$, a contradiction. So, we have $w_{3}=0$. This means that the vector $w$ lies in the $x y$ plane. But $w=\alpha^{\prime} \times \beta-u \beta^{\prime} \times \beta$ and the vector $\beta^{\prime} \times \beta$ lies in the $x y$ plane. So $\alpha^{\prime} \times \beta$ lies in the $x y$ plane. This means that the vectors $\alpha^{\prime} \times \beta$ and $\beta^{\prime} \times \beta$ are parallel. If we put $\alpha^{\prime} \times \beta=\mu \beta^{\prime} \times \beta$, then $\left(\alpha^{\prime}-\mu \beta^{\prime}\right) \times \beta=0$ or $\alpha^{\prime}=\mu \beta^{\prime}$. So $\mu=\left\langle\alpha^{\prime}, \beta^{\prime}\right\rangle=u$ and $\alpha^{\prime}=u \beta^{\prime}$, namely $w=0$.

Now we conclude that $q=(t+u)^{2}$ and the Gauss map is constant, which means that $M^{2}$ is a plane.

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