# NONTRIVIAL SOLUTIONS OF A SEMILINEAR ELLIPTIC PROBLEM VIA VARIATIONAL METHODS 

Zhi-Qing Han

Using variational methods, we investigate the existence of nontrivial solutions of a nonlinear elliptic boundary value problem at resonance under generalised Ahmad-Lazer-Paul conditions. Some new results are obtained and some results in the literature are improved.

## 1. Introduction

In this paper we consider the existence of nontrivial solutions of the following problem on a bounded domain $\Omega \subset \mathbb{R}^{n}$ with smooth boundary

$$
\begin{cases}\Delta u+\lambda_{k} u+g(x, u)=0 & x \in \Omega  \tag{1.1}\\ u=0, & x \in \partial \Omega\end{cases}
$$

where $g(x, t)$ is a Caratheodory function such that $g(x, 0)=0$ for almost everhwhere $x \in \Omega$ and $\lambda_{k}$ is the $k-t h(k \geqslant 2)$ eigenvalue for the elliptic linear operator $-\Delta$ with zero Dirichlet boundary condition.

It is well-known that the operator $-\Delta$ with zero Dirichlet boundary condition has discrete eigenvalues $(0<) \lambda_{1}<\lambda_{2}<\cdots<\lambda_{k}<\lambda_{k+1}<\cdots$ and each eigenspace, is finite dimensional. Denote the eigenspace corresponding to $\lambda_{i}$ by $E_{i}$ and suppose that $E_{k}=\operatorname{span}\left\{\phi_{1}, \phi_{2}, \ldots, \phi_{m}\right\}$.

We impose the following conditions.

$$
\text { (g1) }|g(x, t)| \leqslant C|t|^{\alpha}+b(x)
$$

where $b(x) \in L^{q}$ with $q=(2 N / N+2)$ for $N \geqslant 3, q=1$ for $N=1,2$ and $C>0,0 \leqslant \alpha<1$ are constants;

$$
(\mathrm{G} \pm) \frac{\int_{\Omega} G\left(x, \sum_{i=1}^{m} a_{i} \phi_{i}\right) d x}{\|a\|^{2 \alpha}} \rightarrow \pm \infty \text { as }\|a\|=\left(\sum_{i=1}^{m} a_{i}^{2}\right)^{1 / 2} \rightarrow \infty
$$

Received 7th August, 2003
Copyright Clearance Centre, Inc. Serial-fee code: 0004-9727/04 \$A2.00+0.00.
where $a=\left(a_{1}, a_{2}, \ldots, a_{m}\right) \in \mathbb{R}^{m}$ and $G(x, t)=\int_{0}^{t} g(x, s) d s$.
The conditions ( $\mathrm{G} \pm$ ) with $\alpha=0$ are proposed by Ahmad-Lazer-Paul ([1]) and are in some sense coercivity conditions on the function $G(x, t)$.

Much work has been done on the existence of solutions or multiple solutions to the problem (1.1) since the work in [12], either by topological or by variational methods; for example see $[6,10,14,15,16]$ and the references therein. To obtain (nontrivial) solutions, one of the difficulties is that $\left(g(x, t)+\lambda_{k} t\right) / t$ may approach some $\lambda_{i}$ as $t \rightarrow \infty$ (resonant at infinity) or as $t \rightarrow 0$ (resonant at 0 ), since when resonance at infinity occurs, it is difficult to obtain the priori estimates needed by the topological methods or obtain the Palais-Smale condition required by the variational methods. The Ahmad-Lazer-Paul conditions have been widely used in the literature to overcome the difficulty. For the strong resonance case, that is, $g(x, t) \rightarrow 0$ as $t \rightarrow \infty$, where Ahmad-Lazer-Paul conditions fail, $[3,4,5]$ develop some variational techniques to investigate the nontrivial solutions of (1.1). For the nonresonant or incompletely resonant case, there is also a lot of work in this respect; for example see $[2,11,14]$. But it seems that there is not too much work to deal with the middle case where $g(x, t)$ satisfies conditions like (g1). For some related results see [13]. In [9], we proposed conditions ( $\mathrm{G} \pm$ ) to investigate the existence of solutions and proved the following theorem.

Theorem 1.1. Suppose that condition pair (g1), ( $\mathrm{G}_{+}$) or (g1), (G-) holds. Then equation (1.1), where we do not assume that $g(x, 0)=0$ for almost everywhere $x \in \Omega$, has at least one solution in $H_{0}^{1}(\Omega)$.

The above theorem with $\alpha=0$ was proved in [1]; see also [15]. In this paper, we aim to investigate the nontrivial solutions to (1.1) under the conditions (g1) and ( $\mathrm{G}_{+}$) or (G-) and obtain the following results.

TheOrem 1.2 Suppose that conditions (g1) and ( $\mathrm{G}_{+}$) hold. If there exists $m \leqslant k$ such that

$$
\begin{equation*}
\limsup _{t \rightarrow 0} \frac{g(x, t)}{t}<\lambda_{m}-\lambda_{k} \tag{1.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\inf _{t \neq 0} \frac{g(x, t)}{t} \geqslant \lambda_{m-1}-\lambda_{k} \tag{1.3}
\end{equation*}
$$

uniformly for almost everywhere $x \in \Omega$, then equation (1.1) has at least one nontrivial solution in $H_{0}^{1}(\Omega)$.

Theorem 1.3. Suppose that conditions (g1) and (G-) hold. If there exists $m \geqslant k$ such that

$$
\begin{equation*}
\liminf _{t \rightarrow 0} \frac{g(x, t)}{t}>\lambda_{m}-\lambda_{k} \tag{1.4}
\end{equation*}
$$

and

$$
\begin{equation*}
\sup _{t \neq 0} \frac{g(x, t)}{t} \leqslant \lambda_{m+1}-\lambda_{k} \tag{1.5}
\end{equation*}
$$

uniformly for almost everywhere $x \in \Omega$, then equation (1.1) has at least one nontrivial solution in $H_{0}^{1}(\Omega)$.

It is natural to investigate the case $m>k$ in Theorem 1.2 and the case $m<k$ in Theorem 1.3. The corresponding results are interesting, since the coercivity condition ( $\mathrm{G}_{+}$) or condition ( $\mathrm{G}_{-}$) is not indispensable. In particular, Theorem 1.5 contains one of the main results in [14, Theorem 1] as a special case.

Theorem 1.4. Suppose that condition (g1) and the Palais-Smale condition for $J$ at any level $c<0$ hold. If there exists $m>k$ such that (1.4) and (1.5) hold, then equation (1.1) has at least one nontrivial solution in $H_{0}^{1}(\Omega)$.

Theorem 1.5. Suppose that condition (g1) and the Palais-Smale condition for $J$ at any level $c>0$ hold. If there exists $m<k$ such that such that (1.2) and (1.3) hold, then equation (1.1) has at least one nontrivial solution in $H_{0}^{1}(\Omega)$.

## 2. Proofs of the Theorems

In the following, the notations $\|\cdot\|$ and $\langle\cdot\rangle$ denote the norm in $H_{0}^{1}(\Omega)$ and the pairing between $H^{-1}(\Omega)$ and $H_{0}^{1}(\Omega)$. $C$ denotes a universal constant. For $u \in H_{0}^{1}(\Omega)$ and $p>0$, denote

$$
\|u\|_{p}=\left(\int_{\Omega}|u|^{p} d x\right)^{1 / p}
$$

and decompose $u$ as $u=\bar{u}+u^{0}+\widetilde{u}$, where $\bar{u} \in \sum_{i<k} E_{i}, u^{0} \in E_{k}$ and $\widetilde{u} \in \sum_{i>k} E_{i}$.
It is well-known that (weak) solutions of (1.1) in $H_{0}^{1}(\Omega)$ correspond to the critical points of the $C^{1}$ functional in $H_{0}^{1}(\Omega)$

$$
J(u)=\frac{1}{2} \int_{\Omega}|\nabla u|^{2} d x-\frac{1}{2} \lambda_{k} \int_{\Omega}|u|^{2} d x-\int_{\Omega} G(x, u) d x
$$

and

$$
\left\langle J^{\prime}(u), v\right\rangle=\int_{\Omega}\left(\nabla u \nabla v-\lambda_{k} u v-g(x, u) v\right) d x
$$

for $u, v \in H_{0}^{1}(\Omega)$ (see [15]).
If we want to get classical solutions, we need to impose more regularity assumptions on $g(x, t)$, for example that $g$ is locally Lipschitz in $\bar{\Omega} \times \mathbb{R}$; see [8] for more details.

Lemma 2.1. Under the condition pair (g1), (G+) or (g1), (G-), the functional J defined above satisfies the Palais-Smale condition.

Proof: We only prove the case where (g1) and ( $\mathrm{G}_{+}$) hold. The other case can be similarly proved.

Suppose that $\left\{u_{n}\right\} \subset H_{0}^{1}(\Omega)$ satisfies

$$
\begin{gather*}
J^{\prime}\left(u_{n}\right) \rightarrow 0 \quad \text { in } H^{-1}(\Omega), \text { as } n \rightarrow \infty ; \\
 \tag{2.1}\\
\left|J\left(u_{n}\right)\right| \leqslant C
\end{gather*}
$$

Hence

$$
\begin{aligned}
& \left\langle J^{\prime}\left(u_{n}\right),-\bar{u}_{n}\right\rangle \\
& \quad=\int_{\Omega}\left(-\left|\nabla \bar{u}_{n}\right|^{2}+\lambda_{k}|\bar{u}|^{2}+g\left(x, u_{n}\right) \bar{u}_{n}\right) d x \\
& \quad \geqslant\left(\lambda_{k}-\lambda_{k-1}\right) \int_{\Omega}\left|\bar{u}_{n}\right|^{2} d x-\int_{\Omega}\left|\bar{u}_{n}\right|\left(C\left|\bar{u}_{n}+u_{n}^{0}+\widetilde{u}_{n}\right|^{\alpha}+b\right) d x \\
& \quad \geqslant\left(\lambda_{k}-\lambda_{k-1}\right) \int_{\Omega}\left|\bar{u}_{n}\right|^{2} d x-\int_{\Omega}\left|\bar{u}_{n}\right| b d x-3^{\alpha} C \int_{\Omega}\left|\bar{u}_{n}\right|\left(\left|\bar{u}_{n}\right|^{\alpha}+\left|u_{n}^{0}\right|^{\alpha}+\left|\bar{u}_{n}\right|^{\alpha}\right) d x \\
& \quad \geqslant\left(\lambda_{k}-\lambda_{k-1}-\varepsilon\right) \int_{\Omega}\left|\bar{u}_{n}\right|^{2} d x-C \int_{\Omega}\left|\bar{u}_{n}\right|\left|u_{n}^{0}\right|^{\alpha} d x-C \int_{\Omega}\left|\bar{u}_{n}\right||\widetilde{u}|^{\alpha} d x-C(\varepsilon) \\
& \quad \geqslant\left(\lambda_{k}-\lambda_{k-1}-2 \varepsilon\right) \int_{\Omega}\left|\bar{u}_{n}\right|^{2} d x-C(\varepsilon) \int_{\Omega}\left|u_{n}^{0}\right|^{2 \alpha} d x-C(\varepsilon) \int_{\Omega}\left|\widetilde{u}_{n}\right|^{2 \alpha} d x-C(\varepsilon)
\end{aligned}
$$

where $C(\varepsilon)>0$ is a universal constant dependent on the arbitrary $\varepsilon>0$. Fixing $\varepsilon>0$ sufficiently small and noting that all norms in $\sum_{i<k} E_{i}$ are equivalent, we have

$$
\begin{equation*}
\left\|\bar{u}_{n}\right\|^{2} \leqslant C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha}+C\left\|\widetilde{u}_{n}\right\|_{2 \alpha}^{2 \alpha}+C . \tag{2.2}
\end{equation*}
$$

By a similar argument we can prove that

$$
\left\langle J^{\prime}\left(u_{n}\right), \widetilde{u}_{n}\right\rangle \geqslant\left(1-\frac{\lambda_{k}}{\lambda_{k+1}}-2 \varepsilon\right) \int_{\Omega}\left|\nabla \widetilde{u}_{n}\right|^{2} d x-C(\varepsilon) \int_{\Omega}\left|\bar{u}_{n}\right|^{2 \alpha} d x-C(\varepsilon) \int_{\Omega}\left|u_{n}^{0}\right|^{2 \alpha} d x-C(\varepsilon)
$$

Hence we have the inequality

$$
\begin{equation*}
\left\|\widetilde{u}_{n}\right\|^{2} \leqslant C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha}+C\left\|\bar{u}_{n}\right\|_{2 \alpha}^{2 \alpha}+C . \tag{2.3}
\end{equation*}
$$

Combining (2.2) and (2.3), we have

$$
\begin{align*}
\left\|\widetilde{u}_{n}\right\|^{2} & \leqslant C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha}+C\left\|\bar{u}_{n}\right\|^{2 \alpha}+C \\
& \leqslant C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha}+C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha^{2}}+C\|\tilde{u}\|_{2 \alpha}^{2 \alpha^{2}}+C . \tag{2.4}
\end{align*}
$$

By the Hölder inequality and the Sobolev inequality, in view of $2 \alpha^{2}<2 \alpha<2$, it follows immediately from (2.4) that

$$
\left\|\widetilde{u}_{n}\right\|^{2} \leqslant C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha}+C\|\widetilde{u}\|_{2 \alpha}^{2 \alpha^{2}}+C .
$$

Consequently,

$$
\begin{equation*}
\left\|\widetilde{u}_{n}\right\|^{2} \leqslant C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha}+C . \tag{2.5}
\end{equation*}
$$

By a similar argument to that in the proof (2.5), we obtain

$$
\begin{equation*}
\left\|\bar{u}_{n}\right\|^{2} \leqslant C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha}+C . \tag{2.6}
\end{equation*}
$$

Now we estimate $\int_{\Omega}\left(G\left(x, u_{n}\right)-G\left(x, u_{n}^{0}\right)\right) d x$.

$$
\begin{aligned}
\int_{\Omega}\left(G\left(x, u_{n}\right)-G\left(x, u_{n}^{0}\right)\right) d x= & \int_{\Omega} d x \int_{0}^{1} g\left(x, u_{n}^{0}+s\left(\widetilde{u}_{n}+\bar{u}_{n}\right)\right)\left(\widetilde{u}_{n}+\bar{u}_{n}\right) d s \\
\leqslant & \int_{\Omega} d x \int_{0}^{1}\left(\left|\widetilde{u}_{n}\right|+\left|\bar{u}_{n}\right|\right)\left(C\left|u_{n}^{0}+s\left(\widetilde{u}_{n}+\bar{u}_{n}\right)\right|^{\alpha}+b\right) d s \\
\leqslant & C \int_{\Omega}\left(\left|\widetilde{u}_{n}\right|\left|u_{n}^{0}\right|^{\alpha}+\left|\widetilde{u}_{n}\right|^{1+\alpha}+\left|\widetilde{u}_{n}\right|\left|\bar{u}_{n}\right|^{\alpha}+b\left|\widetilde{u}_{n}\right|\right) d x \\
& \quad+\int_{\Omega}\left(\left|\bar{u}_{n}\right|\left|u_{n}^{0}\right|^{\alpha}+\left|\bar{u}_{n}\right|\left|\widetilde{u}_{n}\right|^{\alpha}+\left|\bar{u}_{n}\right|^{1+\alpha}+b\left|\bar{u}_{n}\right|\right) d x .
\end{aligned}
$$

By (2.5) and (2.6) and a simple calculation, we can obtain

$$
\begin{equation*}
\int_{\Omega}\left(G\left(x, u_{n}\right)-G\left(x, u_{n}^{0}\right)\right) d x \leqslant C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha}+C \tag{2.7}
\end{equation*}
$$

Obviously, by (2.1) and the definition of J,

$$
-C \leqslant \frac{1}{2} \int_{\Omega}\left|\nabla \tilde{u}_{n}\right|^{2} d x-\int_{\Omega}\left(G\left(x, u_{n}\right)-G\left(x, u_{n}^{0}\right)\right) d x-\int_{\Omega} G\left(x, u_{n}^{0}\right) d x
$$

Moreover, by (2.3) and (2.7), we have

$$
-C \leqslant C\left\|u_{n}^{0}\right\|_{2 \alpha}^{2 \alpha}+C-\int_{\Omega} G\left(x, u_{n}^{0}\right) d x
$$

Write $u_{n}^{0}=\sum_{i=1}^{m} a_{i}^{n} \phi_{i}$. The above inequality is converted to

$$
-C \leqslant C\left(\sum_{i=1}^{m}\left(a_{i}^{n}\right)^{2}\right)^{\alpha}+C-\int_{\Omega} G\left(x, \sum_{i=1}^{m} a_{i}^{n} \phi_{i}\right) d x
$$

Hence $\left\{\sum_{i=1}^{m}\left(a_{i}^{n}\right)^{2}\right\}$ is bounded by the condition $\left(G_{+}\right)$. Furthermore, $\left\{u_{n}\right\}$ is bounded in $H_{0}^{1}(\Omega)$ by (2.5) and (2.6). A standard argument implies that J satisfies the Palais-Smale condition in $H_{0}^{1}(\Omega)$ (see [15]).

Proof of Theorem 1.2: We only need to prove the existence of nontrivial solutions of $J$ in $H_{0}^{1}(\Omega)$. Write $H_{0}^{1}(\Omega)=\sum_{i<m} E_{i} \oplus \sum_{i \geqslant m} E_{i}$. In order to use [15, Theorem 5.3], in view of Lemma 2.1, we need to verify the following conditions:
(i) there are $\rho, d>0$ such that $J \geqslant d$ on $\left\{u \in \sum_{i \geqslant m} E_{i} \mid\|u\|=\rho\right\}$;
(ii) there are $e \in \sum_{i \geqslant m} E_{i}$ with $\|e\|=1, R>\rho$, and $\varepsilon<d$ such that if $Q$ $=\left\{u \in \sum_{i<m} E_{i} \mid\|u\| \leqslant R\right\} \oplus\{t e: 0<t<R\}$, then $J \leqslant \varepsilon$ on $\partial Q$, where $\partial Q$ denotes the boundary of $Q$ in $\sum_{i<m} E_{i} \oplus \mathbb{R} e$.
By (g1) and (1.2), the condition (i) can be proved by the argument in the proof of [14, Theorem 1]. By (1.3), it is obvious that $J \leqslant 0$ on $\sum_{i \leqslant m-1} E_{i}$. Hence, in order to obtain (ii), we only need to prove

$$
\begin{equation*}
\lim _{\substack{\|u\| \rightarrow \infty \\ u \in \sum_{i \leqslant m} E_{i}}} J(u)=-\infty \tag{2.8}
\end{equation*}
$$

since then $e$ can be taken as any element in $E_{m}$ with $\|e\|=1, R$ any number sufficiently large and $\varepsilon<d$ any number sufficiently small.

In fact, for $u \in \sum_{i \leqslant m} E_{i}$, we have $u=\bar{u}+u^{0}$, since $m \leqslant k$. Suppose that $m=k$. Then

$$
\begin{aligned}
J(u) & =\frac{1}{2} \int_{\Omega}\left(|\nabla \bar{u}|^{2}-\lambda_{k}|\bar{u}|^{2}\right) d x-\int_{\Omega}\left(G(x, u)-G\left(x, u^{0}\right)\right) d x-\int_{\Omega} G\left(x, u^{0}\right) d x \\
& \leqslant \frac{1}{2}\left(\lambda_{k-1}-\lambda_{k}\right) \int_{\Omega}|\bar{u}|^{2} d x-\int_{\Omega} d x \int_{0}^{1} g\left(x, u^{0}+s \bar{u}\right) \bar{u} d s-\int_{\Omega} G\left(x, u^{0}\right) d x \\
& \leqslant \frac{1}{2}\left(\lambda_{k-1}-\lambda_{k}\right) \int_{\Omega}|\bar{u}|^{2} d x+C \int_{\Omega}|\bar{u}|\left(\left|u^{0}\right|^{\alpha}+|\bar{u}|^{\alpha}+b\right) d x-\int_{\Omega} G\left(x, u^{0}\right) d x \\
& \leqslant \frac{1}{2}\left(\lambda_{k-1}-\lambda_{k}+\varepsilon\right) \int_{\Omega}|\bar{u}|^{2} d x+C \int_{\Omega}|\bar{u}|\left|u^{0}\right|^{\alpha} d x-\int_{\Omega} G\left(x, u^{0}\right) d x+C(\varepsilon) \quad(\forall \varepsilon>0) \\
(2.9) & \leqslant \frac{1}{2}\left(\lambda_{k-1}-\lambda_{k}+2 \varepsilon\right) \int_{\Omega}|\bar{u}|^{2} d x+C(\varepsilon) \int_{\Omega}\left|u^{0}\right|^{2 \alpha} d x-\int_{\Omega} G\left(x, u^{0}\right) d x+C(\varepsilon) .
\end{aligned}
$$

Choosing $0<\varepsilon<\left(\lambda_{k}-\lambda_{k-1}\right) / 2$ and using the condition ( $\mathrm{G}_{+}$), we obtain (2.8).
If $m<k$, then $u=\bar{u}$. The proof of (2.8) is much easier. The theorem is proved. $]$
Proof of Theorem 1.3: Under the conditions of the theorem, we can prove
(i) there are $\rho, d>0$ such that $J \leqslant-d$ on $\left\{u \in \sum_{i \leqslant m} E_{i} \mid\|u\|=\rho\right\}$;
(ii) $J \geqslant 0$ on $\sum_{i \geqslant m+1} E_{i}$;
(iii) $J(u) \rightarrow+\infty$ as $u \in \sum_{i \geqslant m} E_{i}$ and $\|u\| \rightarrow \infty$.

Then, for $I=-J$, we use [15, Theorem 5.29] and obtain a positive(nonzero) critical value for $I$. This completes the proof.

Proofs of Theorems 1.4-1.5: Under the conditions of Theorem 1.4, we obtain a negative critical value for $J$ by obtaining (i)-(iii) in the proof of Theorem 1.3, where in the proof of (iii) we notice $m>k$ and use a similar (and simpler) argument to the proof of (2.9). The proof of Theorem 1.5 can be given as that of Theorem 1.2.

Lemma 2.2. Suppose that $g(x, t)$ satisfies (g1) and there exist $a(x) \in L^{1}(\Omega)$ $a(x) \geqslant 0$ with $\int_{\Omega} a(x) d x>0$ and $b(x) \in L^{1}(\Omega)$ such that
(g2) $a(x) \leqslant \liminf _{t \rightarrow \infty} \frac{g(x, t) t}{|t|^{1+\alpha}} \leqslant \underset{t \rightarrow \infty}{\limsup } \frac{g(x, t) t}{|t|^{1+\alpha}} \leqslant b(x)$
uniformly for almost everywhere $x \in \Omega$. Then the condition ( $\mathrm{G}_{+}$) holds.
Proof: It is easy to get an $\varepsilon>0$ such that

$$
\int_{\Omega} a(x)|v(x)|^{1+\alpha} d x \geqslant \varepsilon \int_{\Omega}|v(x)|^{1+\alpha} d x
$$

for all $v \in E_{k}, v \neq \theta$. For this $\varepsilon>0, \exists T>0$ such that $g(x, t) t \geqslant(a(x)-\varepsilon)|t|^{1+\alpha}$ for $|t|>T$. Hence

$$
\begin{align*}
G(x, t)-G(x, 0) & =\int_{0}^{1} g(x, t s) t d s \\
& =\int_{|t s|>T} \frac{1}{s} g(x, t s) t s d s+\int_{|t s| \leqslant T} \frac{1}{s} g(x, t s) t s d s \\
& \geqslant \int_{|t s|>T} \frac{1}{s}(a(x)-\varepsilon)|t s|^{1+\alpha} d s-|t| \int_{|t s| \leqslant T} h_{\varepsilon}(x) d s \\
& \geqslant \frac{1}{1+\alpha}(a(x)-\varepsilon)|t|^{1+\alpha}-h_{\varepsilon}(x)|t| \tag{2.10}
\end{align*}
$$

for some $h_{\varepsilon} \in L^{q}(\Omega)$.
For any sequence $\left\{u_{n}\right\} \in E_{k}$ with $\left\|u_{n}\right\| \rightarrow \infty$, set $v_{n}=u_{n} /\left\|u_{n}\right\|$. Without loss of generality, we assume that $v_{n} \rightarrow v$ in $C(\bar{\Omega})$ and $\left|v_{n}(x)\right| \leqslant C$ for almost everywhere $x \in \Omega$ and $n \geqslant 1$, where $v \neq \theta$. Therefore, by Fatou lemma and the inequality (2.10),

$$
\begin{aligned}
& \liminf _{n \rightarrow \infty}\left\|u_{n}\right\|^{-(\alpha+1)} \int_{\Omega} G\left(x, u_{n}\right) d x \\
& \quad \geqslant \liminf _{n \rightarrow \infty}\left\|u_{n}\right\|^{-(\alpha+1)} \int_{\Omega}\left(\frac{1}{1+\alpha}(a(x)-\varepsilon)\left|u_{n}\right|^{1+\alpha}-h_{\varepsilon}(x)\left|u_{n}\right|\right) d x \\
& \quad \geqslant \frac{1}{\alpha+1} \int_{\Omega}(a(x)-\varepsilon)|v(x)|^{1+\alpha} d x>0 .
\end{aligned}
$$

Consequently, by $\alpha+1>2 \alpha$, we have

$$
\lim _{n \rightarrow \infty}\left\|u_{n}\right\|^{-2 \alpha} \int_{\Omega} G\left(x, u_{n}\right) d x=+\infty .
$$

This completes the proof.
Lemma 2.3 Suppose that $g(x, t)$ satisfies (g1) and there exist $a(x) \in L^{1}(\Omega)$ and $b(x) \in L^{1}(\Omega), b(x) \leqslant 0$ with $\int_{\Omega} b(x) d x<0$ such that
(g3) $a(x) \leqslant \liminf _{t \rightarrow \infty} \frac{g(x, t) t}{|t|^{1+\alpha}} \leqslant \limsup _{t \rightarrow \infty} \frac{g(x, t) t}{|t|^{1+\alpha}} \leqslant b(x)$
uniformly for almost everywhere $x \in \Omega$. Then the condition (G-) holds.
Proof: The proof is similar to that of Lemma 2.2.

Corollary 2.1 Suppose that the assumptions in Theorem 1.2 hold with ( $\mathrm{G}_{+}$) replaced by (g2). Then (1.1) has at least one nontrivial solution in $H_{0}^{1}(\Omega)$.
REMARK 2.1. The above corollary is essentially proved in [16, Theorem 1]. Existence of nontrivial solutions of (1.1) under (g2) and some other conditions is also investigated in [13] by Morse theory where more regularity conditions on $g(x, t)$ and different conditions near $t=0$ are needed.

The condition (1.2) is a one-sided nonresonant condition at the origin with respect to the eigenvalue $\lambda_{m}$. Using the ideas in [7, 14], we can relax it. First, we make some preparations.

A measurable subset $E$ of $\mathbb{R}$ is said to have positive density at $+0(-0)$ if

$$
\liminf _{r \rightarrow+0} \frac{\operatorname{meas}(E \cap[0, r])}{\operatorname{meas}([0, r])}>0 \quad\left(\liminf _{r \rightarrow-0} \frac{\operatorname{meas}(E \cap[r, 0])}{\operatorname{meas}([r, 0])}>0\right)
$$

We say that a measurable subset $A$ of a measurable set $B$ is a full subset of $B$ if $B \backslash A$ has measurable zero. For $A \subset \Omega$ and $r>0$, write

$$
E(A, r)=\bigcap_{x \in A}\left\{t \in \mathbb{R} \backslash\{0\}: \frac{G(x, t)}{2 t^{2}} \leqslant r\right\} .
$$

Now we present the following improvement of Theorem 1.2. Other theorems in this paper can be similarly improved.

Theorem 2.1. Suppose that conditions (g1) and ( $\mathrm{G}_{+}$) hold. Assume there exists $m \leqslant k$ such that

$$
\limsup _{t \rightarrow 0} \frac{g(x, t)}{t} \leqslant \lambda_{m}-\lambda_{k}
$$

and

$$
\inf _{t \neq 0} \frac{g(x, t)}{t} \geqslant \lambda_{m-1}-\lambda_{k}
$$

uniformly for almost everywhere $x \in \Omega$. If there exists a full subset $\Omega^{\prime}$ of $\Omega$ and $\eta>0$ such that $E\left(\Omega^{\prime}, \lambda_{m}-\lambda_{k}-\eta\right)$ has positive density at +0 or ( -0 ), then equation (1.1) has at least one nontrivial solution in $H_{0}^{1}(\Omega)$.

Proof: The proof can be given combining the arguments in the proof of Theorem 1.2 in this paper and those in the proof of [14, Theorem 2].

Remark 2.2: Finally, we point out that since there are no coercive conditions on $g(x, t)$ in Theorems 1.4 and 1.5 , by weakening the Palais-Smale condition to the (C) condition (for example see [3]), these theorems may be applied to investigate the strong resonance case for the problem (1.1). Further research may appear elsewhere.

## References

[1] S. Ahmad, A.C. Lazer and J.L. Paul, 'Elementary critical point theory and perturbations of elliptic boundary value problems at resonance', Indiana Univ. Math. J. 25 (1976), 933-944.
[2] H. Amann and E. Zehnder, 'Nontrivial solutions for a class of non-resonance problems and applications to nonlinear differential equations', Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4) 7 (1980), 539-603.
[3] P. Bartolo, V. Benci and D. Fortunato, 'Abstract critical point theorems and applications to nonlinear problems with "strong" resonance at infinity', Nonlinear Anal. 7 (1983), 981-1012.
[4] A. Capozzi, D. Lupo and S. Solimini, 'On the existence of a nontrivial solution to nonlinear problems at resonance', Nonlinear Anal. 13 (1989), 151-163.
[5] K.C. Chang and J.Q. Liu, 'A strong resonance problem', Chinese Ann. Math. Ser. B 11 (1990), 191-210.
[6] E.N. Dancer, 'On the Dirichlet problem for weakly non-linear elliptic partial differential equations', Proc. Roy. Soc. Edinburgh Sect. A 76 (1977), 283-300.
[7] D.G. De Figueiredo and J.-P Gossez, 'Nonresonance below the first eigenvalue for a semilinear elliptic problem', Math. Ann. 281 (1988), 589-610.
[8] D. Gilbarg and N.S. Trudinger, Elliptic partial diffferential equations of second order (Springer-Verlag, Berlin, Heidelberg, New York, Tokyo, 1983).
[9] Z.Q. Han, 'An improvement on the Ahmad-Lazer-Paul's theorem', J. Math. (Wuhan) 16 (1996), 512-518.
[10] Z.Q. Han, 'Solvability of elliptic boundary value problems without standard Landes-man-Lazer conditions', Acta Math. Sin. Engl. Ser. 16 (2000), 349-360.
[11] N. Hirano, 'Existence of nontrivial solutions of semilinear elliptic equations', Nonlinear Anal. 13 (1989), 695-705.
[12] E.M. Landesman and A. Lazer, 'Nonlinear perturbations of a linear elliptic boundary value problems at resonance', J. Math. Mech. 19 (1970), 609-623.
[13] S.J. Li and W.M. Zou, 'The computations of the critical groups with an application to elliptic resonant problems at a higher eigenvalue', J. Math. Anal. Appl. 235 (1999), 237-259.
[14] N. Mizoguchi, 'Asymptotically linear elliptic equations without nonresonance conditions', J. Differential Equations 113 (1994), 150-165.
[15] P. Rabinowitz, Minimax methods in critical point theory with applications to differential equations, CBMS Regional Conference Series in Mathematics 65 (American Mathematical Society, Providence, RI, 1986)).
[16] C.L. Tang and Q.J. Gao, 'Elliptic resonant problems at higher eigenvalues with an unbounded nonlinear term', J. Differential Equations 146 (1998), 56-66.

Department of Applied Mathematics
Dalian University of Technology
Dalian 116023
Liaoning
Peoples Republic of China
e-mail: hanzhiq@dlut.edu.cn

