On the deformation of the tangent *m*-plane of a V_n^m

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1. Schouten and van Kampen $(1)^1$ have studied the deformation of a V_n^m . Applying the methods of that paper to the tangent vectors $B_i^{\lambda}(\lambda, \mu, \nu, \ldots = 1, 2, \ldots, n; i, j, k, \ldots = 1, 2, \ldots, m)$, which exist by hypothesis at all points of a certain region $V_{m'}(m' > m)$ of V_n , we shall have

(1)
$$\begin{cases} d^{\lambda} B_{i}^{\lambda} = v^{\mu} dt \partial_{\mu} B_{i}^{\lambda} \\ d^{\lambda} B_{i}^{\lambda} = -\Gamma_{\mu\nu}^{\lambda} B_{i}^{\mu} v^{\nu} dt, \qquad \Gamma_{\mu\nu}^{\lambda} = \begin{pmatrix} \lambda \\ \mu\nu \end{pmatrix} \\ \frac{3}{d} B_{i}^{\lambda} = B_{i}^{\mu} \partial_{\mu} v^{\lambda} dt \end{cases}$$

whence we define the differentials

(2)
$$d = d = \delta = D dt$$
 (r, s = 1, 2, 3 with r $\pm s$).

In the application of the $D^{(r,s)}$ to B_i^{λ} the lower index is treated as an ordinal index only. We shall not be concerned with any extension of the $D^{(r,s)}$ to indices other than those of the general V_n (see 1, equ. 3.24). We here consider the application of $D^{(r,s)}$ to the *m*-vector $B_1^{[\lambda_1} \ldots B_m^{\lambda_m]}$ determining the *m*-plane tangent to the facet V_n^m at any point. This gives

(3)
$$\overset{(r,s)}{D}(B_1^{[\lambda_1}\dots B_m^{\lambda_m]}) = (\overset{(r,s)}{D}B_1^{[\lambda_1]})B_2^{\lambda_2}\dots B_m^{\lambda_m]} + B_1^{[\lambda_1}(\overset{(r,s)}{D}B_2^{\lambda_2})\dots B_m^{\lambda_m]} + \dots + B_1^{[\lambda_1}\dots B_{m-1}^{\lambda_{m-1}}(\overset{(r,s)}{D}B_m^{\lambda_m]}).$$

Since the vector $D B_i^{\lambda}$ will not be perpendicular to the B_i^{λ} in general, it will have a component in the tangent *m*-plane, and a component perpendicular to it. Consequently from (3) the *m*-vector $D B_1^{\lambda_1} \dots B_m^{\lambda_m}$ will have two components respectively parallel and perpendicular to the tangent *m*-plane. We shall study those components in a few special cases.

¹ These numbers refer to the list of papers at the end.

The component of the vector $\stackrel{(r, s)}{D} B_1^{\lambda_1}$ in the tangent *m*-plane is $B_{\nu}^{\lambda_1} \stackrel{(r, s)}{D} B_1^{\nu}$, which we may write as

(4)
$$(B^{1}_{\nu}D^{(r,s)}B^{\nu}_{1})B^{\lambda_{1}}_{1} + (B^{2}_{\nu}D^{(r,s)}B^{\nu}_{1})B^{\lambda_{1}}_{2} + \ldots + (B^{m}_{\nu}D^{(r,s)}B^{\nu}_{1})B^{\lambda_{1}}_{m}$$

and on inserting this series in the right-hand side of (3), it is evident that only the first term contributes anything, giving therefore

$$(B^{1}_{\nu} D^{(r, s)} B^{\nu}_{1}) B^{[\lambda_{1}}_{1} \dots B^{\lambda_{m}]}_{m}$$

Treating the vector $D B_{2}^{\lambda_{2}}$ in the same way, giving an expression corresponding to (4) and inserting in (3), we shall have

$$(B_{\nu}^2 \overset{(r,s)}{D} B_2^{\nu}) B_1^{[\lambda_1} \ldots B_m^{\lambda_m]},$$

. .

so that, on treating all the other vectors in the same way, we conclude that the total component of (3) in the original *m*-plane can be written in the form of the single *m*-vector

(5)
$$(B_{\nu}^{i} D B_{i}^{\nu}) B_{1}^{[\lambda_{1}} \dots B_{m}^{\lambda_{m}]}.$$

Let us now put for brevity

(6)
$$P_{j^{j}}^{\lambda_{j}} = C_{v}^{\lambda_{j}} \stackrel{(r,s)}{D} B_{j}^{v};$$

then the component of (3) perpendicular to the tangent *m*-vector can be written as the sum of *m* such *m*-vectors: —

$$P_{1}^{[\lambda_{1}} B_{2}^{\lambda_{2}} \dots B_{m}^{\lambda_{m}]} + B_{1}^{[\lambda_{1}} P_{2}^{\lambda_{2}} \dots B_{m}^{\lambda_{m}]} + \dots + B_{1}^{[\lambda_{1}} \dots B_{m-1}^{\lambda_{m-1}} P_{m}^{\lambda_{m}]}$$

so that (3) becomes

(7)
$$\overset{(r,s)}{D} B_1^{[\lambda_1} \dots B_m^{\lambda_m]} = (B_\nu^i \overset{(r,s)}{D} B_i^\nu) B_1^{[\lambda_1} \dots B_m^{\lambda_m]} + \sum_{j=1}^m B_1^{[\lambda_1} \dots P_j^{\lambda_j} \dots B_m^{\lambda_m]}.$$

2. Let us now take the special case D = D, and $v^{\lambda} = B_{\kappa}^{\lambda} u^{\kappa}$. Then

(8)
$$B_{\nu}^{i} \overset{(1,2)}{D} B_{i}^{\nu} = B_{\nu}^{i} D_{\kappa} B_{(i)}^{\nu} = \sum_{i=1}^{m} \left\{ \begin{array}{c} i \\ i \\ \kappa \end{array} \right\} u^{\kappa}$$

where the index (i) is ordinal.

If the metric coefficients of the V_n^m are $b_{ij} = a_{\lambda\mu} B_{ij}^{\lambda\mu}$, with $b = |b_{ij}|$, then it is a well-known fact that

(9)
$$\sum_{i=1}^{m} \left\{ \begin{matrix} i \\ i \\ \kappa \end{matrix} \right\} = \frac{\partial \log \sqrt{b}}{\partial \eta^{\kappa}},$$

E. T. DAVIES

where η has the meaning used in 1, so that $u^{\kappa} dt = d\eta^{\kappa}$ and

(10)
$$\sum_{i=1}^{m} \left\{ \begin{matrix} i \\ i \\ \kappa \end{matrix} \right\} = \frac{d \log \sqrt{b}}{dt}.$$

Also for this case

204

(11)
$$P_{j}^{\lambda_{j}} = C_{\nu^{j}}^{\lambda_{j}} D_{\kappa} B_{(j)}^{\nu} = H_{j\kappa}^{\lambda} u^{\lambda}$$

where $H_{j\kappa}^{\lambda_j}$ is the first tensor of Eulerian Curvature. Using (8), (10) and (11) we have

(12)
$$\begin{array}{c} D^{(1,2)} B_1^{[\lambda_1} \dots B_m^{\lambda_m]} = \frac{d}{dt} \log \sqrt{b} \cdot B_1^{[\lambda_1} \dots B_m^{\lambda_m]} \\ + u^{\kappa} \sum_{j=1}^m B_1^{[\lambda_1} \dots H_{j^{\kappa}}^{[\lambda_j]} \dots B_m^{\lambda_m]}. \end{array}$$

A sufficient condition for V_n^m to be geodesic is that $H_{j\kappa}^{\lambda_j} = 0$, in which case (12) becomes

(13)
$$D^{(1,2)} B_1^{[\lambda_1} \dots B_m^{\lambda_m]} - \frac{d}{dt} \log \sqrt{b} B_1^{[\lambda_1} \dots B_m^{\lambda_m]} = 0.$$

If, instead of a V_n^m , we have a single curve V_1 , and if the displacement $d\xi^{\lambda} = v^{\lambda} dt$ is along the tangent, then by taking t to be the parameter of the curve, the $B_1^{\lambda} = \frac{d\xi^{\lambda}}{dt}$ becomes the (non-unit) tangent vector, so that, if s is the length of arc reckoned from a certain point, then

$$\sqrt{b} = \sqrt{b_{11}} = \sqrt{\left(a_{\lambda\mu}\frac{d\xi^{\lambda}}{dt} \frac{d\xi^{\mu}}{dt}\right)} = \frac{ds}{dt},$$

and consequently

$$\frac{d \log \sqrt{b}}{dt} = \frac{d \log \frac{ds}{dt}}{dt} = \frac{\frac{d^2s}{dt^2}}{\frac{ds}{dt}},$$

19

and (13) becomes

(14)
$$\frac{d^2 \xi^{\mu}}{dt^2} + \left\{ \begin{array}{c} \mu \\ \rho \nu \end{array} \right\} \frac{d\xi^{\rho}}{dt} \quad \frac{d\xi^{\nu}}{dt} - \frac{\frac{d^2 s}{dt^2}}{\frac{ds}{dt}} \quad \frac{d\xi^{\mu}}{dt} = 0,$$

which is the well-known form for the equation of geodesic lines when t is not the arc length.

3. If
$$D = D$$
, we shall have $D B_i^{(3,2)} = D_i v^{\lambda}$, and consequently

$$B_{\nu}^{i} \overset{(3, 5)}{D} B_{i}^{\nu} = B_{\nu}^{i} D_{i} v^{\nu} = \frac{1}{2} b^{ij} B_{ij}^{\lambda\mu} D_{\lambda\mu} a_{\lambda\mu},$$

where D is the "Lie derivative" (1).

If d_{τ}^{m} is the element of volume in the local R_{m} , with (see 1, formula 3.50)

$$d^{m}_{\tau} = (d\eta)^{1} (d\eta)^{2} \dots (d\eta)^{m} \sqrt{b},$$

then the deformation of that element of volume is of amount

$$d\tau^{m}_{\frac{1}{2}}b^{ij}B^{\lambda\mu}_{ij}\delta_{L}a_{\lambda\mu} = d\tau^{m}B^{i}_{\nu}D_{i}v^{\nu}dt,$$

so that if $d^{m}_{\tau'}$ denotes the deformed volume element, we can write

(15)
$$\frac{d\tau'}{d\tau} = 1 + B^i_{\nu} D_i v^{\nu} dt,$$

 $B^i_{\nu} D_i v^{\nu}$ being therefore the "dilatation" at the point. If we are dealing with a curve, whose element of arc is ds and whose unit tangent vector is $i^{\lambda} = \frac{d\xi^{\lambda}}{ds}$, then if ds' is the element of arc of the deformed curve, the above equation becomes

(16)
$$\frac{ds'}{ds} = 1 + \phi dt, \text{ with } \phi = i_{\lambda} \frac{D}{Ds} v^{\lambda}$$

If $C_{\nu}^{\lambda_j} D B_j^{\nu} = 0$, (j = 1, 2, ..., m), then the deformed *m*-vector $B_1^{\lambda_1} \dots B_m^{\lambda_{m_1}}$ will be parallel to the original one. For a curve, the resulting equations

(17)
$$D^{(3,2)}_{i\lambda} = \phi i^{\lambda}$$

determine the parallel-tangent deformations (Hayden, 2).

4. We pass now to the component of the *m*-vector (3) perpendicular to the tangent plane. For this purpose we remark that if da is the angle between the two *m*-vectors v and v+wdt, and if w=w'+w'', where w' is parallel to v and w'' is perpendicular to it, then

(18)
$$\frac{\sin da}{dt} = \frac{|w''|}{|v|},$$

where |v| denotes the "measure" (3, Chapter 1) of the *m*-vector v.

Applying this to the case where v and w have the components

 $v^{\lambda_1 \dots \lambda_m} = B_1^{[\lambda_1} \dots B_m^{\lambda_m]}$ and $w^{\lambda_1 \dots \lambda_m} = D^{(r,s)} B_1^{[\lambda_1} \dots B_m^{\lambda_m]}$ we obtain on simplification

(19) $\frac{\sin d\alpha}{dt} = \sqrt{(b^{ij} a_{\lambda\mu} P_i^{\lambda} P_j^{\mu})}.$

If D = D, and if we are dealing with a curve, with $v^{\lambda} = \frac{d\xi^{\lambda}}{dt}$ as the tangent vector, the P_i^{λ} reduces to the left-hand side of equation (14), so that the invariant $\sin da/dt$ is the first curvature of the curve. For an isolated V_m in general, with D = D and $v^{\lambda} = B_i^{\lambda} u^i$, the right-hand side of (19) becomes the square root of the "forma angolare" of Bortolotti, which can therefore be regarded as the first curvature of the V_m in V_n . The particular form taken by (19) in that case has already been given (4, p. 294), together with its expression when D = D and $v^{\lambda} = C_p^{\lambda} w^p$.

- J. A. Schouten and E. R. van Kampen, "Beiträge zur Theorie der Deformation," Prace Mat. Fiz. Warszawa, 41 (1933), 1-19.
- H. A. Hayden, "Deformations of a curve in a Riemannian space," Proc. London Math. Soc. (2), 32 (1931), 321-336.
- E. Cartan, "Leçons sur la Géométrie des Espaces de Riemann," Paris, Gauthier-Villars, 1928.
- 4. E. T. Davies, "On the second and third fundamental forms of a sub-space," Journal London Math. Soc., 12 (1937), 290-295.

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206