

Double tail structure in escaping atmospheres of magnetised close-in planets

A. A. Vidotto¹, S. Carolan², G. Hazra^{1,3}, C. Villarreal D'Angelo⁴ and D. Kubyshkina⁵

¹Leiden Observatory, Leiden University, PO Box 9513, 2300 RA Leiden, The Netherlands

²School of Physics, Trinity College Dublin, College Green, D02 PN40 Dublin 2, Ireland

³Dept. of Astrophysics, University of Vienna, Türkenschanzstrasse 17, A-1180 Vienna, Austria

⁴Inst. Astronomía Téorica y Exp. (CONICET-UNC), Laprida 854, Córdoba, Argentina

⁵Space Research Inst., Austrian Academy of Sciences, Schmiedlstrasse 6, A-8042 Graz, Austria

Abstract. High-energy stellar irradiation can photoevaporate planetary atmospheres, which can be observed in spectroscopic transits of hydrogen lines. Here, we investigate the effect of planetary magnetic fields on the observational signatures of atmospheric escape in hot Jupiters.

Keywords. MHD, planets and satellites: atmospheres, magnetic fields, planet-star interactions

In this work, we use our newly developed 3D self-consistent radiative magnetohydrodynamic (MHD) simulations (Carolan et al. 2021; Hazra et al. 2022) to study the effects of planetary magnetic fields on the dynamics of escaping atmosphere. To investigate the resulting observational signature, we couple the results of our 3D models to Lyman- α transit calculations (Vidotto et al. 2018; Allan & Vidotto 2019). In our models, we account for high energy stellar photons ionising the atmospheric neutral hydrogen atoms, which affects both the heating deposition in the atmosphere of the planet and its ionisation state. In addition to photoionisation, we also include collisional ionisation and recombination and track the proton and neutral components of the flow. Additionally, we also consider Lyman- α cooling and collisional cooling. The reader will be able to find further details of the model and this work in Carolan et al. (2021).

In our model, the stellar wind is injected in the numerical grid through an external boundary (see Figure 1). Using the same stellar wind property, we vary the planet's dipole field strength from 0 to 10 G. In Figure 1, we show a typical structure of the escaping atmosphere in magnetised planets (> 3 G). Escape occurs through polar outflows, as opposed to the predominantly comet-like tail from non-magnetised models (see comparison in Figure 2). We find a small increase in evaporation rate with planetary field, though this should not affect the timescale of atmospheric loss (Figure 3a).

The double-tail structure has some key effects in Lyman- α transit signatures, as summarised in Figure 3b. When considering magnetic fields, we see an increase in line centre absorption due to an increase of the size of dead-zones. Additionally, we also see an increase in redshifted absorption with increase in planetary field. Most of the red shifted material exists around the night-side orbital plane, as some material falls from the comet-like tails back towards the planet. Finally, we also see that with increased magnetic field, there is an initial decrease in blueshifted absorption, as planetary material begins to be

 \odot The Author(s), 2023. Published by Cambridge University Press on behalf of International Astronomical Union. This is an Open Access article, distributed under the terms of the Creative Commons Attribution licence (https://creativecommons.org/licenses/by/4.0/), which permits unrestricted re-use, distribution, and reproduction in any medium, provided the original work is properly cited.



Figure 1. The polar outflows seen in magnetised planets lead to the formation of a double tail structure, above and below the orbital plane, and a dead-zone around the equator.



Figure 2. Side view of the planet orbit showing the total (ρ) and neutral Hydrogen (ρ_{neu}) density for unmagnetised (left panels) and magnetised (right panels) Hot Jupiters.



Figure 3. In spite of the small increase in escape rate with planetary field (red line, panel a), we find that $Ly-\alpha$ transit observations are strongly affected by planetary magnetism (panel b).

launched above and below the orbital plane, instead of being fully funnelled on to the orbital plane by the stellar wind, as seen in the 0-G model.

For further details of this work, please see Carolan et al. (2021).

References

Allan, A. & Vidotto, A. A. 2019, MNRAS, 490, 3760

- Carolan, S., Vidotto, A. A., Hazra, G., Villarreal D'Angelo, C., & Kubyshkina, D. 2021, MNRAS, 508, 6001
- Hazra, G., Vidotto, A. A., Carolan, S., Villarreal D'Angelo, C., & Manchester, W. 2022, MNRAS, 509, 5858

Vidotto, A. A., Lichtenegger, H., Fossati, L., et al. 2018, MNRAS, 481, 5296