# $A V$-Courant Algebroids and Generalized CR Structures 

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#### Abstract

We construct a generalization of Courant algebroids that are classified by the third cohomology group $H^{3}(A, V)$, where $A$ is a Lie Algebroid, and $V$ is an $A$-module. We see that both Courant algebroids and $\mathcal{E}^{1}(M)$ structures are examples of them. Finally we introduce generalized CR structures on a manifold, which are a generalization of generalized complex structures, and show that every CR structure and contact structure is an example of a generalized CR structure.


## 1 Introduction

Courant algebroids and the Dirac structures associated with them were first introduced by Courant and Weinstein (see $[6,7]$ ) to provide a unifying framework for studying such objects as Poisson and symplectic manifolds. Aïssa Wade later introduced the related $\mathcal{E}^{1}(M)$-Dirac structures in [27] to describe Jacobi structures.

In [13], Hitchin defined generalized complex structures that are further described by Gualtieri [12]. Generalized complex structures unify both symplectic and complex structures, interpolating between the two, and have appeared in the context of string theory [17]. In [14] Iglesias and Wade describe generalized contact structures, an odd-dimensional analog to generalized complex structures, using the language of $\mathcal{E}^{1}(M)$-Dirac structures.

In this paper, we shall define $A V$-Courant Algebroids, a generalization of Courant algebroids that also allows one to describe $\mathcal{E}^{1}(M)$-Dirac structures. We will show that these have a classification similar to Severa's classification of exact Courant algebroids in [24].

To be more explicit, let $M$ be a smooth manifold, $A \rightarrow M$ be a Lie algebroid with anchor map $a: A \rightarrow T M$, and $V \rightarrow M$ a vector bundle that is an $A$-module. If we endow $V$ with the structure of a trivial Lie algebroid (that is, trivial bracket and anchor), then it is well known that the extensions of $A$ by $V$ are a geometric realization of $H^{2}(A, V)$ (see [18]). In this paper, we introduce $A V$-Courant algebroids and describe how they are a geometric realization of $H^{3}(A, V)$.

We then go on to show how to simplify the structure of certain $A V$-Courant algebroids by pulling them back to certain principal bundles. Indeed, in the most interesting cases, the pullbacks will simply be exact Courant algebroids.

We then introduce $A V$-Dirac structures, a special class of subbundles of an $A V$ Courant algebroid which generalize Dirac structures. Finally, we will introduce a special class of $A V$-Dirac structures, called generalized CR structures, which allow

[^0]us to describe any complex, symplectic, CR or contact structure on a manifold, as well as many interpolations of those structures. We show that associated with every generalized CR structure is a Jacobi bundle, introduced by Charles-Michel Marle [21] and Kirillov [16].

It is important to note that there are other constructions related to $A V$-Courant algebroids. For instance, recently Z. Chen, Z. Liu, and Y.-H. Sheng introduced the notion of $E$-Courant algebroids [5] in order to unify the concepts of omni-Lie algebroids (introduced in [3], see also [4]) and generalized Courant algebroids or Courant-Jacobi algebroids (introduced in [23] and [10] respectively; they are equivalent concepts; see [23]). The key property that both $E$-Courant algebroids and $A V$ Courant algebroids share is that they replace the $\mathbb{R}$-valued bilinear form of Courant algebroids with one taking values in an arbitrary vector bundle ( $E$ or $V$ respectively). Nevertheless, while there is some overlap between $E$-Courant algebroids and $A V$-Courant algebroids in terms of examples, these constructions are not equivalent; indeed, $A V$-Courant algebroids are classified by $H^{3}(A, V)$, while there is no simple classification of $E$-Courant algebroids. Moreover, this paper is distinguished from [5] by having the definition of generalized CR manifolds as one of its main goals.

Meanwhile, generalized CRF structures, introduced and studied in great detail by Izu Vaisman in [26], and generalized CR structures describe similar objects. To summarize, a complex structure on a manifold $M$ is a subbundle $H \subset T M \otimes \mathbb{C}$ such that

$$
\begin{equation*}
H \oplus \bar{H}=T M \otimes \mathbb{C} \tag{1.1}
\end{equation*}
$$

and $[H, H] \subset H$. The definition of a $C R$ structure simply relaxes (1.1) to $H \cap \bar{H}=0$. On the other hand, the definition of a generalized complex structure replaces $T M$ with the standard Courant algebroid $T M=T^{*} M \oplus T M$ in the definition of a complex structure, and in addition, requires $H \subset \mathbb{T} M \otimes \mathbb{C}$ to be isotropic.

The definition of a generalized CRF structure parallels the definition of a generalized complex structure, but relaxes the requirement that $H \oplus \bar{H}=\mathbb{T} M \otimes \mathbb{C}$ to $H \cap \bar{H}=0$. Among numerous interesting examples of generalized CRF structures are normal contact structures and normalized CR structures (namely those CR structures $H \subset T M \otimes \mathbb{C}$ for which there is a splitting $T M \otimes \mathbb{C}=H \oplus \bar{H} \oplus Q_{c}$ and $\left.\left[H, Q_{c}\right] \subset H \oplus Q_{c}\right)$.

Generalized CR structures differ from generalized CRF structures in multiple ways. In particular, they replace the standard Courant algebroid with an $A V$-Courant $\operatorname{algebroid} \mathbb{A}$, and furthermore, they take a different approach to describe contact and CR structures, using only maximal isotropic subbundles but allowing $H \cap \bar{H}$ to contain "infinitesimal" elements.

## 2 AV-Courant Algebroids

Let $M$ be a smooth manifold, $A \rightarrow M$ a Lie algebroid, and $V \rightarrow M$ a vector bundle that is an $A$-module, that is, there is a $C^{\infty}(M)$-linear Lie algebra homomorphism

$$
\begin{equation*}
\mathcal{L} .: \Gamma(A) \rightarrow \operatorname{End}(\Gamma(V)) \tag{2.1}
\end{equation*}
$$

satisfying the Leibniz rule. (See [18] for more details.)
For any $A$-module $V$, the sections of $V \otimes \wedge^{*} A^{*}$ have the structure of a graded right $\wedge^{*} \Gamma\left(A^{*}\right)$-module, and there are several important derivations of its module structure that we shall use throughout this paper. The first is the interior product with a section $X \in \Gamma(A)$,

$$
\iota_{X}: \Gamma\left(V \otimes \wedge^{i} A^{*}\right) \rightarrow \Gamma\left(V \otimes \wedge^{i-1} A^{*}\right)
$$

a derivation of degree -1 .
The second is the Lie derivative, a derivation of degree 0 , defined to be the unique derivation of $V \otimes \wedge^{*} A^{*}$ whose restriction to $V$ is given by (2.1), and such that the graded commutator with $\iota$. satisfies $\left[\mathcal{L}_{X}, \iota_{Y}\right]=\iota_{[X, Y]}$. Finally, the differential $d$, a derivation of degree 1 , is defined inductively by the graded commutator $\mathcal{L}_{X}=\left[d, \iota_{X}\right]$ (for all $X \in \Gamma(A)$ ).

It is easy to check that $d^{2}=0$, and the cohomology groups of the complex $(\Gamma) V \otimes$ $\left.\left.\wedge^{\bullet} A^{*}\right), d\right)$ are denoted $H^{\bullet}(A, V)$.

### 2.1 Definition of $A V$-Courant Algebroids

Let $A$ be a Lie algebroid and $V$ an $A$-module.
Definition 2.1 ( $A V$-Courant Algebroid) Let $A$ be a vector bundle over $M$, with a $V$-valued symmetric bilinear form $\langle\cdot, \cdot\rangle$ on the fibres of $\mathbb{A}$, and a bracket $\llbracket \cdot, \cdot \rrbracket$ on sections of $\mathbb{A}$. Suppose further that there is a short exact sequence of bundle maps

$$
\begin{equation*}
0 \rightarrow V \otimes A^{*} \xrightarrow{j} \mathbb{A} \xrightarrow{\pi} A \rightarrow 0 \tag{2.2}
\end{equation*}
$$

such that for any $e \in \Gamma(\mathbb{A})$ and $\xi \in \Gamma\left(V \otimes A^{*}\right)$,

$$
\begin{equation*}
\langle e, j(\xi)\rangle=\iota_{\pi(e)} \xi \tag{2.3}
\end{equation*}
$$

The bundle $A$ with these structures is called an $A V$-Courant algebroid if, for $f \in$ $C^{\infty}(M)$ and $e, e_{i} \in \Gamma(\mathbb{A})$, the following axioms are satisfied:
$(\mathrm{AV}-1) \llbracket e_{1}, \llbracket e_{2}, e_{3} \rrbracket \rrbracket=\llbracket \llbracket e_{1}, e_{2} \rrbracket, e_{3} \rrbracket+\llbracket e_{2}, \llbracket e_{1}, e_{3} \rrbracket \rrbracket$
$(\mathrm{AV}-2) \pi\left(\llbracket e_{1}, e_{2} \rrbracket\right)=\left[\pi\left(e_{1}\right), \pi\left(e_{2}\right)\right]$
(AV-3) $\llbracket e, e \rrbracket=\frac{1}{2} D\langle e, e\rangle$, where $D=j \circ d$
$(\mathrm{AV}-4) \mathcal{L}_{\pi\left(e_{1}\right)}\left\langle e_{2}, e_{3}\right\rangle=\left\langle\llbracket e_{1}, e_{2} \rrbracket, e_{3}\right\rangle+\left\langle e_{2}, \llbracket e_{1}, e_{3} \rrbracket\right\rangle$
we will often refer to $\llbracket \cdot, \cdot \rrbracket$ as the Courant bracket.
Remark 2.2 Axioms (AV-1) and (AV-4) state that $\llbracket e, \cdot \rrbracket$ is a derivation of both the Courant bracket and the bilinear form, while Axiom (AV-2) describes the relation of the Courant bracket to the Lie algebroid bracket of $A$. One should interpret Axiom (AV-3) as saying that the failure of $\llbracket \cdot, \cdot \rrbracket$ to be skew symmetric is only an "infinitesimal" $D(\cdot)$.

Remark 2.3 The bracket is also derivation of $\mathbb{A}$ as a $C^{\infty}(M)$-module in the sense that

$$
\llbracket e_{1}, f e_{2} \rrbracket=f \llbracket e_{1}, e_{2} \rrbracket+a \circ \pi\left(e_{1}\right)(f) \cdot e_{2}
$$

for any $e_{1}, e_{2} \in \Gamma(\mathbb{A})$ and $f \in C^{\infty}(M)$. In fact if $e_{3} \in \Gamma(\mathbb{A})$,

$$
\begin{aligned}
\left\langle a \circ \pi\left(e_{1}\right)(f) \cdot e_{2}+\right. & \left.f \llbracket e_{1}, e_{2} \rrbracket-\llbracket e_{1}, f e_{2} \rrbracket, e_{3}\right\rangle \\
(\mathrm{by}(\mathrm{AV}-4 \rrbracket)== & \left\langle a \circ \pi\left(e_{1}\right)(f) \cdot e_{2}+f \llbracket e_{1}, e_{2} \rrbracket, e_{3}\right\rangle-\pi\left(e_{1}\right)\left\langle f e_{2}, e_{3}\right\rangle \\
& +\left\langle f e_{2}, \llbracket e_{1}, e_{3} \rrbracket\right\rangle \\
= & a \circ \pi\left(e_{1}\right)(f)\left\langle e_{2}, e_{3}\right\rangle-\pi\left(e_{1}\right)\left\langle f e_{2}, e_{3}\right\rangle \\
& +f\left(\left\langle\llbracket e_{1}, e_{2} \rrbracket, e_{3}\right\rangle+\left\langle e_{2}, \llbracket e_{1}, e_{3} \rrbracket\right\rangle\right) \\
(\operatorname{by}(\mathrm{AV}-4))= & a \circ \pi\left(e_{1}\right)(f)\left\langle e_{2}, e_{3}\right\rangle-\pi\left(e_{1}\right)\left\langle f e_{2}, e_{3}\right\rangle+f \pi\left(e_{1}\right)\left\langle e_{2}, e_{3}\right\rangle \\
= & 0,
\end{aligned}
$$

where the last equality follows from the fact that $V$ is an $A$ module. Since this holds for all $e_{3} \in \Gamma(\mathbb{A})$, and $\langle\cdot, \cdot\rangle$ is non-degenerate, the statement follows.

Remark 2.4 One notices that (2.3) and exactness of (2.2) implies that the map $\mathbb{A} \rightarrow V \otimes \mathbb{A}^{*}$, given by $e \rightarrow\langle e, \cdot\rangle$, is an injection. Consequently, if $V$ is a line bundle (as in all the known interesting examples), it follows that $\mathbb{A} \simeq V \otimes \mathbb{A}^{*}$, and $j$ must be the composition

$$
j: V \otimes A^{*} \xrightarrow{\mathrm{id} \otimes \pi^{*}} V \otimes \mathbb{A}^{*} \simeq \mathbb{A} .
$$

Remark 2.5 Any $\mathfrak{D E} E$-Courant algebroid (an $A V$-Courant algebroid with $V=E$ and $A=\mathfrak{D E}$, the gauge Lie algebroid of $E$ ) is an $E$-Courant algebroid. However, not every $E$-Courant algebroid is a $\mathfrak{D E E}$-Courant algebroid, since there is no requirement in the definition of $E$-Courant algebroids for the sequence (2.2) in Definition 2.1) of $A V$-Courant algebroids to be exact, and the map (2.2) $j: E \otimes(\mathfrak{D} E)^{*} \rightarrow \mathbb{A}$ is only defined on the first jet bundle $\mathfrak{J}^{1} E \subset E \otimes(D E)^{*}$.

One could imagine some generalization of both $A V$-Courant algebroids and $E$-Courant algebroids that ignores the requirement that (2.2) be exact in the above definition (and perhaps allows $j$ to be defined on a smaller domain).

Conversely, if $\mathbb{A}$ is an $A V$-Courant algebroid, then there is a natural Lie algebroid morphism $\phi: A \rightarrow \mathfrak{D V}$ resulting from the fact that $V$ is an $A$-module. Consequently, (A), $\langle\cdot, \cdot\rangle, \llbracket \cdot, \cdot \rrbracket, \phi \circ \pi$ ) is an $E$-Courant algebroid (with $E=V$ ). So an $A V$-Courant algebroid can be thought of as an $E$-Courant algebroid with some additional structure, such as a exact sequence (2.2) and a factorization of the anchor map through a Lie algebroid $A$. This additional structure allows for a more comprehensive understanding of $A V$-Courant algebroids, including a simple classification of $A V$-Courant algebroids by $H^{3}(A, V)$, and when $A$ is a transitive Lie algebroid, a means of understanding both $A V$-Courant algebroids and $A V$-Dirac structures by relating them to standard Courant algebroids and Dirac structures on principal bundles.

### 2.2 Splitting

We call $\phi: A \rightarrow \mathbb{A}$ an isotropic splitting, if it splits the exact sequence (2.2), and $\phi(A)$ is an isotropic subspace of $\mathbb{A}$ with respect to the inner product.

Remark 2.6 Such splittings exist. In fact we may choose a splitting $\lambda: A \rightarrow \mathbb{A}$, which is not necessarily isotropic.

Then we have a map $\gamma: A \rightarrow V \otimes A^{*}$ given by the composition

$$
\gamma: A \xrightarrow{\lambda} \mathbb{A} \xrightarrow{e \rightarrow\langle e, \cdot\rangle} V \otimes \mathbb{A}^{*} \xrightarrow{\text { id } \otimes \lambda^{*}} V \otimes A^{*} .
$$

We let $\phi=\lambda-\frac{1}{2} j \circ \gamma$. It is easy to check that $\phi$ is an isotropic splitting.
If $\phi: A \rightarrow \mathbb{A}$ is an isotropic splitting, then we have an isomorphism $\phi \oplus j: A \oplus$ $\left(V \otimes A^{*}\right) \rightarrow \mathbb{A}$.

Proposition 2.7 Let $\phi: A \rightarrow \mathbb{A}$ be an isotropic splitting. Then under the above isomorphism, the bracket on $A \oplus\left(V \otimes A^{*}\right)$ is given by

$$
\begin{equation*}
\llbracket X+\xi, Y+\eta \rrbracket_{\phi}=[X, Y]+\mathcal{L}_{X} \eta-\iota_{Y} d \xi+\iota_{X} \iota_{Y} H_{\phi}, \tag{2.4}
\end{equation*}
$$

where $X, Y \in \Gamma(A), \xi, \eta \in \Gamma\left(V \otimes A^{*}\right)$ and $H_{\phi} \in \Gamma\left(V \otimes \wedge^{3} A^{*}\right)$, with $d H_{\phi}=0$.
Furthermore, if $\psi: A \rightarrow \mathbb{A}$ is a different choice of isotropic splitting, then $\psi(X)=$ $\phi(X)+j\left(\iota_{X} \beta\right)$ and $H_{\psi}=H_{\phi}-d \beta$, where $\beta \in \Gamma\left(V \otimes \wedge^{2} A^{*}\right)$.

The proof is relegated to the appendix, since it is parallel to the proof for ordinary Courant algebroids (see [2,24]).

Theorem 2.8 Let A be a Lie algebroid, and let $V$ be an A-module. Then the isomorphism classes of $A V$-Courant algebroids are in bijective correspondence with $H^{3}(A, V)$.

Proof If $H \in \Gamma\left(V \otimes \wedge^{3} A^{*}\right)$, and $d H=0$, then let $\mathbb{A}=A \oplus\left(V \otimes A^{*}\right)$. We define $\langle\cdot, \cdot\rangle$ by

$$
\begin{equation*}
\langle X+\xi, Y+\eta\rangle=\iota_{X} \eta+\iota_{Y} \xi \tag{2.5}
\end{equation*}
$$

where $\xi, \eta \in \Gamma\left(V \otimes A^{*}\right)$ and $X, Y \in \Gamma(A)$. We define the bracket to be given by equation (2.4). It is not difficult to check that this satisfies the axioms of an $A V$ Courant algebroid.

Conversely, by the above proposition, every $A V$-Courant algebroid defines a unique element of $H^{3}(A, V)$.

## 3 Examples

Example 3.1 Let $M$ be a point, then a Lie algebroid $A$ is simply a Lie algebra, and an $A$-module $V$ is a finite dimensional representation of $A$ as a Lie algebra. $H^{i}(A, V)$ is simply the $V$-valued Lie algebra cohomology, and $H^{3}(A, V)$ classifies the $A V$-Courant algebroids over a point. Note that an $A V$-Courant algebroid over a point is a Lie algebra if and only if $V$ is a trivial $A$-representation.

Example 3.2 (Exact Courant Algebroids) If we let $A \simeq T M$ and $V=M \times \mathbb{R}$ be the trivial line bundle over $M$ with a trivial $T M$-module structure, then we may identify $T^{*} M$ with $V \otimes T^{*} M$ by the map $\alpha \rightarrow 1 \otimes \alpha$. It follows that the class of $T M \mathbb{R}$-Courant algebroids over $M$ corresponds to the class of exact Courant algebroids (see [6,7]) on M,


Theorem $[2.8$ then corresponds to Ševera's classification of exact Courant algebroids.
Example $3.3\left(\mathcal{E}^{1}(M)\right.$ Structures) The bundle $\mathcal{E}^{1}(M)$ was introduced by A. Wade in [27] and is uniquely associated with a given manifold $M$. Within the context of this paper, it is easiest to define $\mathcal{E}^{1}(M)$ by using the language of $A V$-Courant algebroids.

We let $A=T M \oplus L$, where $L \simeq \mathbb{R}$ is spanned by the abstract symbol $\frac{\partial}{\partial t}$. The bracket is given by

$$
\left[X \oplus f \frac{\partial}{\partial t}, Y \oplus g \frac{\partial}{\partial t}\right]_{A}=[X, Y]_{T M} \oplus(X(g)-Y(f)) \frac{\partial}{\partial t},
$$

where $X, Y \in \mathcal{X}(M)$ and $f, g \in C^{\infty}(M)$.
Let $V$ be the trivial line bundle spanned by the abstract symbol $e^{t}$, so that $\Gamma(V)=$ $\left\{e^{t} h \mid h \in C^{\infty}(M)\right\} . V$ has an $A$-module structure (as suggested by the choice of symbols) given by

$$
\left(X \oplus f \frac{\partial}{\partial t}\right)\left(e^{t} h\right)=e^{t}(X(h)+f h) .
$$

We let $\mathbb{A}:=(T M \oplus L) \oplus\left(V \otimes\left(T^{*} M \oplus L^{*}\right)\right)$, and define a bracket on sections by equation (2.4). It is clear that this data defines an $A V$-Courant algebroid on $M$. If we set $H=0$ in equation (2.4), then the pair ( $\mathbb{A}, \llbracket \cdot, \cdot \rrbracket)$ associated with $M$ is the $\mathcal{E}^{1}(M)$-Structure, as introduced by Wade in [27].

Example 3.4 (Equivariant $A V$-Courant Algebroids on Principal Bundles) Let $\nu: P \rightarrow M$ be a $G$-principal bundle. Suppose that $A$ is a Lie algebroid over $P$ and $V$ is an $A$-module, and that there is an $A V$-Courant algebroid on $P$,

$$
0 \rightarrow V \otimes A^{*} \rightarrow \mathbb{A} \rightarrow A \rightarrow 0
$$

If the action of $G$ on $P$ lifts to an action by bundle maps on $V, A$ and $A$, such that all the structures involved are $G$-equivariant, then the quotient,

$$
0 \rightarrow\left(V \otimes A^{*}\right) / G \rightarrow \mathbb{A} / G \rightarrow A / G \rightarrow 0
$$

is an $A / G V / G$-Courant algebroid.
Example 3.5 Let $\nu: P \rightarrow M$ be a $G$-principal bundle, and let $W$ be a $k$-dimensional vector space possessing a linear action of $G$. We regard $W$ as a trivial bundle over $P$, and we consider the bundle $\mathbb{T}:=T P \oplus\left(W \otimes T^{*} P\right)$, endowed with a $W$-valued symmetric bilinear form given by equation (2.5). We also define a bracket on sections of $T$ by equation (2.4), where $H \in \Omega^{3}(P, W)^{G}$ is closed, then

$$
0 \rightarrow W \otimes T^{*} P \xrightarrow{j} \mathbb{T} \xrightarrow{\pi} T P \rightarrow 0
$$

is an equivariant $T P W$-Courant algebroid on $P$ (where $j$ and $\pi$ are the obvious inclusion and projection). Thus (as in Example 3.4), we have an $A V$-Courant algebroid on $P / G$, where $A=T P / G$ is the Atiyah algebroid, and $V=P \times{ }_{G} W$.

Note, if $W$ is 1-dimensional, then the $T P W$-Courant algebroid given above is simply an exact Courant algebroid.

As it turns out, this is quite a general example. Indeed, if $A$ is a transitive Lie algebroid, then locally all $A V$-Courant algebroids result from such a construction (see Section5).

Remark 3.6 In the above example, one could replace $P \times W$ with any flat bundle.
Example 3.7 As a special case of Example 3.5, if we take $G=\mathbb{R}$, then $P=M \times \mathbb{R}$ is a $\mathbb{R}$-principal bundle where the action is translation. We let $W=P \times \mathbb{R}$ be the trivial line bundle over $P$ and let $\lambda \in \mathbb{R}$ act on $W$ by scaling by $e^{-\lambda}$.

To describe the $G$-action explicitly,

$$
\lambda \cdot((x, s), t)=\left((x, s+\lambda), e^{-\lambda} t\right)
$$

where $\lambda \in \mathbb{R},(x, s) \in M \times \mathbb{R}=P$, and $((x, s), t) \in P \times \mathbb{R}=W$.
The quotient of the $T P W$-Courant algebroid on $P$ with $H=0$ under this action is precisely the $\mathcal{E}^{1}(M)$-Structure on $M=P / \mathbb{R}$.

Example 3.8 If $A$ is a Lie algebroid over $M, V$ is an $A$-module, and $A$ is an $A V$ Courant algebroid on the manifold $M$, and if $F \subset M$ is a leaf of the singular foliation defined by $a(A)$, then $i^{*} A$ is an $i^{*} A i^{*} V$-Courant algebroid on $F$, where $i: F \rightarrow M$ is the inclusion.

Remark 3.9 At this point, in the most interesting examples of $A V$-Courant algebroids, $V$ is a line bundle. Nevertheless, as mentioned in Theorem 2.8 for any Lie algebroid $A \rightarrow M$, any $A$-module $V$ over $M$, and any element $\gamma \in H^{3}(A, V)$, there is an $A V$-Courant algebroid (unique up to isomorphism) classified by $\gamma$. It is not yet known if these examples are of any importance.

## $4 A V$-Dirac Structures

Definition 4.1 ( $A V$-Dirac Structure) Let $M$ be a manifold, let $A \rightarrow M$ be a Lie algebroid over $M$, let $V \rightarrow M$ be an $A$-module, and let $\mathbb{A}$ be an $A V$-Courant algebroid. Suppose that $L \subset \mathbb{A}$ is a subbundle, since $\mathbb{A}$ has a non-degenerate inner product, we can define $L^{\perp}=\{v \in \mathbb{A} \mid\langle v, u\rangle=0 \forall u \in L\}$.

We call $L$ an almost $A V$-Dirac structure if $L^{\perp}=L$. An $A V$-Dirac structure is an almost $A V$-Dirac structure, $L \subset \mathbb{A}$ that is involutive with respect to the bracket $\llbracket \cdot, \cdot \rrbracket$.

Remark 4.2 If $L \subset \mathbb{A}$ is an $A V$-Dirac structure, then $\llbracket e, e \rrbracket=\frac{1}{2} D\langle e, e\rangle=0$ for any section $e \in \Gamma(L)$, so $\llbracket \cdot, \cdot \rrbracket$ is skew-symmetric when restricted to $L$, and then by the other properties of the bracket, it follows that $a \circ \pi: L \rightarrow T M$ is a Lie algebroid, and $\pi: L \rightarrow A$ is a Lie algebroid morphism.

Example 4.3 (Invariant Dirac Structure on a Principal Bundle) Using the notation of Example 3.4 suppose that the $A / G V / G$-Courant algebroid $\mathbb{A} / G$ on $M$ is the quotient of a $A V$-Courant algebroid $\mathbb{A}$ on $P$. If $L \subset \mathbb{A}$ is an $A V$-Dirac structure which is $G$ invariant, then it is clear that $L / G \subset \mathbb{A} / G$ is an $A / G V / G$-Dirac structure (see Example 3.4).

Example $4.4\left(\mathcal{E}^{1}(M)\right.$-Dirac Structures) Using Example 3.3, we can describe $\mathcal{E}^{1}(M)$, the bundle introduced by Wade in [27], as an $A V$-Courant algebroid. In this context, the $\mathcal{E}^{1}(M)$-Dirac structures (also introduced by Wade in [27]) correspond directly to the $A V$-Dirac structures.

## 5 Transitive Lie Algebroids

### 5.1 Simplifying $A V$-Courant Algebroids

Suppose that $A$ is a Lie algebroid, $V$ is an $A$-module, and $\mathbb{A}$ is an $A V$-Courant algebroid over $M$ (where we use the notation given in the definition of $A V$-Courant algebroids). We will assume for the duration of this section that $M$ is connected, and we require that $A$ be a transitive Lie algebroid, namely the anchor map $a: A \rightarrow T M$ is surjective (see [18] for more details).

Since $\mathbb{A}$ may be quite complicated, we wish to examine whether this $A V$-Courant algebroid is the quotient of a much simpler $A^{\prime} V^{\prime}$-Courant algebroid on a principal bundle over $M$, where $A^{\prime}$ is a very simple Lie algebroid and $V^{\prime}$ is a very simple $A^{\prime}$-module. To be more explicit, we wish to examine whether $\mathbb{A}$ results from the construction in Example 3.5 For this to be true, it is clearly necessary that $A$ be the Atiyah algebroid of that principal bundle; namely, if $P$ is the principal bundle, then $A=T P / G$. The existence of such a principal bundle is equivalent to the integrability of $A$ as a Lie algebroid:

Proposition 5.1 Suppose that $A \rightarrow M$ is an integrable transitive Lie algebroid, that is to say, there exists a source-simply connected Lie groupoid

with Lie algebroid A (see [18] for more details). Then A is the Atiyah algebroid of a principal bundle.

Conversely, if $A$ is the Atiyah algebroid of a principal bundle, then $A$ is an integrable Lie algebroid.

Proof Suppose first that $A$ is integrable, then using the notation in the statement of the proposition, where $s: \Gamma \rightarrow M$ is the source map and $t: \Gamma \rightarrow M$ is the target map, let $x \in M$, let $P=\Gamma_{x}:=s^{-1}(x)$, and let $G=\Gamma_{x}^{x}:=s^{-1}(x) \cap t^{-1}(x)$.

Since $A$ is transitive, $t: P \rightarrow M$ is a surjective submersion. For clarity, we define $p:=\left.t\right|_{P}$. Furthermore, if $y \in M$, and $g \in \Gamma_{x}^{y}$, then $g: p^{-1}(x) \rightarrow p^{-1}(y)$ is a diffeomorphism, so $p: P \rightarrow M$ is a fibre bundle, with its fibre diffeomorphic to $G$. In
addition, $G$ has a right action on $P$, given by right multiplication in the Lie groupoid. If $p^{-1}(y)=\Gamma_{x}^{y}$ is a fibre, and $g \in \Gamma_{x}^{y}$, then the diffeomorphism $g: p^{-1}(x) \rightarrow p^{-1}(y)$ is given by left groupoid multiplication while the action of $G$ on $P$ is given by right groupoid multiplication, so it is clear that the two operations commute, from which it follows that $G$ preserves the fibres of $P$, acting transitively and freely on them. Thus $P$ is a principal $G$ bundle.

Since $A$ is the Lie algebroid of $\Gamma$, it can be identified with the right invariant vectorfields on $\Gamma$ tangent to the source fibres. However, since $A$ is transitive, any two source fibres are diffeomorphic by right multiplication by some element. Thus $A$ can be identified with the $G$ invariant vector fields on $P$.

Conversely, if $A$ is the Atiyah algebroid of some principal bundle, it obviously integrates to the gauge groupoid associated with that principal bundle (see [9] or Remark 5.5), and we may take $\Gamma$ to be the source-simply connected cover of the gauge groupoid.

We now examine whether $V$ is an associated vector bundle.
Proposition 5.2 Suppose that $A$ is an integrable transitive Lie algebroid, and $V \rightarrow M$ is an A-module. Then there exists a (possibly disconnected) Lie group $G$ and a simply connected principal $G$-bundle $P \rightarrow M$ such that $V$ is the quotient bundle of $P \times \mathbb{R}^{k}$, for some $G$ action on $\mathbb{R}^{k}$. In this setting, the standard action of $X(P)$ on $C^{\infty}\left(P, \mathbb{R}^{k}\right)$ induces the module structure on $V$.

Proof Using the notation and the Lie groupoid described in the previous proposition, we consider $\Gamma_{x} \times V_{x}$, where $V_{x}$ is the fibre of $V$ at $x$. We may assume that $\Gamma$ is source-simply connected, and, consequently, since $V$ is an $A$-module, by Lie's second theorem there exists a Lie groupoid morphism $\Gamma \rightarrow \mathbf{G L}(V){ }^{1}$ Thus $\Gamma$ acts on $V$, and we have a $\operatorname{map} \tilde{p}: \Gamma_{x} \times V_{x} \rightarrow V$ given by $(g, v) \rightarrow g v$. This is clearly a surjective submersion 2 Furthermore,

$$
\tilde{p}(g, v)=\tilde{p}\left(g^{\prime}, v^{\prime}\right) \Leftrightarrow g^{-1} g^{\prime} \in \Gamma_{x}^{x} \quad \text { and } \quad v=\left(g^{-1} g^{\prime}\right) v^{\prime}
$$

Thus, letting $G=\Gamma_{x}^{x}$ and $P=\Gamma_{x}$, we have $V \simeq\left(\Gamma_{x} \times V_{x}\right) / G \simeq\left(P \times V_{x}\right) / G$.
Furthermore, identifying $V_{x}$ with $\mathbb{R}^{k}$, if $X \in \mathcal{X}(P) \simeq \mathcal{X}\left(\Gamma_{x}\right)$, and $\sigma \in C^{\infty}\left(P, \mathbb{R}^{k}\right)$, then the standard action of $X$ on $\sigma$ is given by $X(\sigma)_{z}=\left.\frac{\partial}{\partial t}\right|_{t=0} \sigma\left(e^{t X} z\right)$ for any $z \in P \simeq$ $\Gamma_{x}$. If we suppose that $X$ and $\sigma$ are $G$ invariant, then

$$
\tilde{p}\left(\left.\frac{\partial}{\partial t}\right|_{t=0} \sigma\left(e^{t X}(z)\right)\right)=\left.\frac{\partial}{\partial t}\right|_{t=0}\left(e^{-t X} \tilde{p}(\sigma)\right)_{p(z)}=\left(\mathcal{L}_{X} \tilde{p}(\sigma)\right)_{p(z)}
$$

since we defined the action of $\Gamma$ on $V$ in terms of the $A$-module structure of $V$.
Proposition 5.3 Suppose that $A$ is an integrable Lie algebroid, and $V \rightarrow M$ is an A-module. Then A results from the construction given in Example 3.5 Namely, there

[^1]exist a Lie group $G$ and a principal $G$-bundle $P \rightarrow M$ such that $\mathbb{A}$ is the quotient of a $T P \mathbb{R}^{k}$-Courant algebroid Furthermore, if $L \subset \mathbb{A}$ is an $A V$-Dirac structure, then it is also the quotient of a corresponding $T P \mathbb{R}^{k}$-Dirac structure on $P$.

Consequently, if $V$ is a line-bundle, then $\mathbb{A}$ is simply the quotient of an exact Courant algebroid on $P$.

Proof We choose some isotropic splitting of $\mathbb{A}$, so that $\mathbb{A} \simeq A \oplus\left(V \otimes A^{*}\right)$. The bracket is given by equation (2.4), and the symmetric bilinear form by equation (2.5). Then we can use the previous propositions to lift the right-hand side to a principal bundle.

By the above propositions, there exist a (possibly disconnected) Lie group $G$ and a simply connected $G$-principal bundle, $\nu: P \rightarrow M$, such that $A \simeq T P / G$. In addition to this there is a $G$-action on $W:=\mathbb{R}^{\operatorname{dim}(V)}$, say $\lambda: G \rightarrow \mathbf{G L}(W)$, such that $V=P \times_{G}$ $W$. In this setting, $\Gamma\left(V \otimes \wedge^{i} A^{*}\right) \simeq \Omega^{i}(P, W)^{G}$, and $d: \Gamma\left(V \otimes \wedge^{i} A^{*}\right) \rightarrow \Gamma\left(V \otimes \wedge^{i+1} A^{*}\right)$ is the restriction of the exterior derivative $d$ to $\Omega^{*}(P, W)^{G}$.

Thus since $H \in \Gamma\left(V \otimes \wedge^{3} A^{*}\right) \simeq \Omega^{3}(P, W)^{G}$, it is clear that we may view $H$ as a $G$-invariant element of $\Omega^{3}(P, W)$ and define the $T P W$-Courant algebroid $W \otimes$ $T^{*} P \rightarrow \mathbb{T} \rightarrow T P$ in terms of it: Namely, $\mathbb{T} \simeq T P \oplus\left(W \otimes T^{*} P\right)$ endowed with a $W$-valued symmetric bilinear form given by equation (2.5), and the bracket given by equation (2.4). (See Example 3.5 for more details on this construction.)

It is clear that $\mathbb{A}$ is the quotient of this TP $W$-Courant algebroid.
Equivalently, it is easy to see that $T P=\nu^{*} A, W=\nu^{*} V$, and $\mathbb{T}=\nu^{*} \mathbb{A}$. The $W$-valued symmetric bilinear form on $\mathbb{T}$ is simply the pullback of the $V$-valued symmetric bilinear form on $\mathbb{A}$, and if $e_{1}, e_{2} \in \Gamma(\mathbb{A})$, then $\llbracket \nu^{*} e_{1}, \nu^{*} e_{2} \rrbracket=\nu^{*} \llbracket e_{1}, e_{2} \rrbracket$, and the bracket on $\mathbb{T}$ is then extended to arbitrary sections of $\mathbb{T}$ by Axiom (AV -3 ) and Remark 2.2

Next, let $\tilde{L}=\nu^{*}(L) \subset \mathbb{T}$. It is obvious that $L^{\perp}=L \Rightarrow \tilde{L}^{\perp}=\tilde{L}$, and, similarly, since $L$ is involutive, so is $\tilde{L}$.

Thus $\tilde{L} \subset \mathbb{T}$ is a $T P W$-Dirac structure, and $\tilde{L} / G=L$.
Example 5.4 If $A=T M$ and $V$ is a flat vector bundle over $M$, then following the proof of Proposition 5.3 we see that $G=\pi_{1}(M)$ is the fundamental group, and $P=\tilde{M}$ is the simply connected covering space of $M$ over which the pullback of $V$ is a trivial vector bundle.

Remark 5.5 The above propositions construct the principal bundle $P$ and the Lie group $G$. Suppose however, that we already have a Lie group $G^{\prime}$ and a connected $G^{\prime}$-principal bundle $\nu^{\prime}: P^{\prime} \rightarrow M$ such that $A \simeq T P^{\prime} / G^{\prime}$. It will not be difficult to see that $\mathbb{A}$ is the quotient of a $A V$-Courant algebroid on $P^{\prime}$.

Let $\mathcal{G}=\left(P^{\prime} \times P^{\prime}\right) / G^{\prime}$, where we take the quotient by the diagonal action. Then $\mathcal{G} \underset{t}{\stackrel{s}{\longrightarrow}} M$ is a Lie groupoid with Lie algebroid $A$, where the source map is $s:[u, v] \rightarrow$ $\nu^{\prime}(v)$, the target map is $t:[u, v] \rightarrow \nu^{\prime}(u)$, and the multiplication is $[u, v] \cdot[v, w]=$ $[u, w]]^{3}$ Hence by Lie's second theorem (see $[8,19,22]$ for more details), since $\Gamma$, the Lie groupoid used in the proof of Proposition 5.3, is source-simply connected, there

[^2]is a unique Lie groupoid morphism $\Phi: \Gamma \rightarrow \mathcal{G}$ that restricts to the identity map on the Lie algebroid $A$.

It follows that $\left.\Phi\right|_{P}: P \rightarrow P^{\prime}$ is a covering map ${ }^{4}$ and $\left.\Phi\right|_{G}: G \rightarrow G^{\prime}$ is a covering morphism of Lie groups $5^{5}$ It is easy to see that $H=\operatorname{ker}\left(\left.\Phi\right|_{G}\right) \simeq \pi\left(P^{\prime}\right)$ and $P^{\prime}=P / H$.

Thus, we may take the quotient of the $T P W$-Courant algebroid on $P$ (constructed in Proposition 5.3) by $H$, to form a $T P^{\prime} W / H$-Courant algebroid on $P^{\prime}$ whose quotient by $G^{\prime}$ is $A$. It is important to note that while $W$ is a trivial vector bundle, $W / H$ is a flat vector bundle.

Remark 5.6 Proposition5.3 was observed for $\mathcal{E}^{1}(M)$ structures in [15].
Corollary 5.7 Suppose that $V$ is an A-module, and $M$ is contractible, then $\mathbb{A}$ is the quotient of a $T P \mathbb{R}^{k}$-Courant algebroid $\mathbb{R}^{k} \otimes T^{*} P \rightarrow \mathbb{T} \rightarrow T P$ on some principal G-bundle, P. (See Example 3.4). Furthermore, if $L \subset \mathbb{A}$ is an $A V$-Dirac structure, then it is also the quotient of a $T P \mathbb{R}^{k}$-Dirac structure $\tilde{L} \subset \mathbb{T}$.

Proof Every transitive Lie algebroid is integrable over a contractible space; see [18] for details.

### 5.2 Contact Manifolds

Iglesias and Wade show how to describe contact manifolds as $\mathcal{E}^{1}(M)$-Dirac structures in [14]. Thus in light of Example 4.4, we can describe them as $A V$-Dirac structures. We will now describe this same construction from a more geometric perspective, similar to their description in [15].

To simplify things, we assume that $(M, \xi)$ is a co-oriented contact manifold, namely $\xi \subset T M$ can be given as the kernel of a nowhere vanishing 1-form $\alpha \in$ $\Omega^{1}(M)$, and we use the fact that there is a one-to-one correspondence between cooriented contact manifolds and symplectic cones (see [1]). Recall, as in [1], that a symplectic manifold $\left(N, \omega_{N}\right)$ is a symplectic cone if

- $N$ is a principal $\mathbb{R}$ bundle over some manifold $B$, called the base of the cone, and
- the action of $\mathbb{R}$ expands the symplectic form exponentially, namely $\rho_{\lambda}^{*} \omega_{N}=e^{\lambda} \omega_{N}$, where $\rho_{\lambda}$ denotes the diffeomorphism defined by $\lambda \in \mathbb{R}$.
In particular, let

$$
N=\left\{q \in T^{*} M \mid q \in T_{x}^{*} M \text { and } q=e^{\tau} \cdot \alpha_{x} \text { for some } x \in M \text { and } \tau \in \mathbb{R}\right\} \subset T^{*} M
$$

then $\mathbb{R}$ acts on $N$ by $\rho_{\lambda}(q)=e^{\lambda} \cdot q\left(\right.$ for any $\lambda \in \mathbb{R}$, and $\left.q \in N \subset T^{*} M\right)$. Furthermore let $\omega_{N} \in \Omega^{2}(N)$ be the restriction to $N$ of the canonical symplectic form on $T^{*} M$, then $\left(N, \omega_{N}\right)$ is a symplectic cone over the base $M$ if and only if $(M, \xi)$ is a co-oriented contact manifold.

Since $\omega_{N}$ is expanded exponentially by the $\mathbb{R}$-action, we can simplify things by instead considering the $\mathbb{R}$-invariant section $1 \otimes \omega_{n}$ of $\Omega^{2}(N, W)$, where $W=N \times \mathbb{R}$

[^3]is the trivial line-bundle over $N$ on which $\lambda \in \mathbb{R}$ acts by scaling by $e^{-\lambda}$ (see Example 3.7).

Now $X \rightarrow 1 \otimes \iota_{X} \omega_{N}$ defines an isomorphism $T N \hookrightarrow W \otimes T^{*} N$. We let $L \subset$ $T N \oplus\left(W \otimes T^{*} N\right)$ be the graph of this morphism. It is easy to check that $L$ is a maximal isotropic subbundle of $T N \oplus\left(W \otimes T^{*} N\right)$, and since $\omega_{N}$ is closed, $L$ is a $T N W$-Dirac subbundle of the $T N W$-Courant algebroid on $N$ defined in Example 3.7. We note that $(M, \xi)$ is a contact manifold if and only if $L$ is the graph of an isomorphism, or simply $L \cap W \otimes T^{*} N=0$ and $L \cap T N=0$.

As described in Example 3.7, the quotient of the TN W-Courant algebroid on $N$ by the $\mathbb{R}$ action yields an $\mathcal{E}^{1}(M)$ bundle on $M$ or an $A V$-Courant algebroid, where $A=T N / \mathbb{R}$, and $V$ is the trivial line bundle on $M$.

Since $1 \otimes \omega_{N}$ is $\mathbb{R}$-invariant, it follows that its graph, $L$, is $\mathbb{R}$-invariant; consequently, $L$ defines an $\mathcal{E}^{1}(M)$-Dirac structure that we denote by $\tilde{L}_{\xi}$. It is perhaps important to note that $\tilde{L}_{\xi}$ is defined intrinsically. We may conclude the following(as shown in [14]).

Proposition $5.8(M, \xi)$ is a contact manifold if and only if $\tilde{L}_{\xi} \cap V \otimes A^{*}=0$ and $\tilde{L}_{\xi} \cap A=0$ (under the canonical splitting).

## 6 CR-structures and Courant Algebroids

Suppose $M$ is a smooth manifold; let $H \subset T M$ be a subbundle, and suppose $J \in$ $\Gamma(\operatorname{Hom}(H, H))$ is such that $J^{2}=-\mathrm{id}$. Then $(H, J)$ is called an almost CR structure. We let $H_{1,0} \subset \mathbb{C} \otimes H \subset \mathbb{C} \otimes T M$ denote the $+i$-eigenbundle of $J$. If $H_{1,0}$ is involutive, then it is called a CR-structure. It is possible to describe this as a Courant algebroid.

We consider the bundle $H^{*} \oplus H \simeq T^{*} M \oplus H / \operatorname{Ann}(H)$ and the bundle map $\mathbb{J}:=-J^{*} \oplus J \in \Gamma\left(\operatorname{Hom}\left(H^{*} \oplus H, H^{*} \oplus H^{*}\right)\right)$. It is clear that $\mathbb{J}^{2}=$-id. Let $L=\operatorname{ker}(J-i) \oplus \operatorname{Ann}(H) \subset \mathbb{C} \otimes\left(T M \oplus T^{*} M\right)$.

Proposition 6.1 L is involutive under the standard Courant bracket if and only if $J$ defines a CR structure.
Proof We notice that $L=H_{1,0} \oplus \operatorname{Ann}\left(H_{1,0}\right)$. Therefore, $L$ is involutive under the Courant bracket only if $\pi(L)=H_{1,0}$ is involutive, where $\pi: T M \oplus T^{*} M \rightarrow T M$ is the projection. Thus $J$ defines a CR structure.

Conversely, suppose that $H_{1,0}$ is involutive. Then if $I$ is the ideal generated by $\operatorname{Ann}\left(H_{1,0}\right)$ in $\Gamma\left(\mathbb{C} \otimes \wedge T^{*} M\right)$, then $I$ is closed under the differential: $d I \subset I$.

In particular, if we restrict our attention to a local neighborhood on $M$, and $\alpha_{i}$ is a local basis for $\operatorname{Ann}\left(H_{1,0}\right)$ and $\xi \in \Gamma\left(\operatorname{Ann}\left(H_{1,0}\right)\right)$, then $d \xi=\sum_{i} \beta_{i} \wedge \alpha_{i}$ for some $\beta_{i} \in \Omega^{1}(M, \mathbb{C})$. Thus, for any $X \in \Gamma\left(H_{1,0}\right)$, we have,

$$
\iota_{X} d \xi=\sum_{i} \beta_{i}(X) \alpha_{i} \in \Gamma\left(\operatorname{Ann}\left(H_{1,0}\right)\right)
$$

and

$$
\mathcal{L}_{X} \xi=d \iota_{X} \xi+\iota_{X} d \xi=\iota_{X} d \xi \in \Gamma\left(\operatorname{Ann}\left(H_{1,0}\right)\right)
$$

It follows that $L$ is involutive under the standard Courant bracket.

In the next section we shall generalize this construction.

## 7 Generalized CR structures

Suppose that $M$ is a manifold; $A$ is a Lie algebroid over $M ; V$ is an $A$-module of rank one over $M$, and $\mathbb{A}$ is an $A V$-Courant algebroid over $M$. Suppose further that $A$ has some distinguished subbundle $H \subset A$, and consider the bundle given by

$$
\mathbb{H I}=q\left(\pi^{-1}(H)\right), \text { where } q: \pi^{-1}(H) \rightarrow \pi^{-1}(H) / j(V \otimes \operatorname{Ann}(H)) .
$$

Then the pairing on $\mathbb{A}$ restricts non-degenerately to $\mathbb{H}$, and we have an exact sequence

$$
0 \rightarrow V \otimes H^{*} \xrightarrow{j} \mathbb{H} \xrightarrow{\pi} H \rightarrow 0
$$

Definition $7.1 J \in \Gamma(\operatorname{Hom}(H, H))$ is called a generalized CR structure if:
(i) J is orthogonal (preserves the pairing on $\mathbb{H I}$ );
(ii) $\mathrm{J}^{2}=-1$;
(iii) $L:=q^{-1}(\operatorname{ker}(J-i)) \subset \mathbb{C} \otimes \mathbb{A}$ is involutive.

Remark 7.2 We have that $L:=q^{-1}(\operatorname{ker}(\mathbb{J}-i)) \subset \mathbb{C} \otimes \mathbb{A}$ is a maximal isotropic subspace of $\mathbb{A}$, since $\operatorname{ker}(J-i)$ is a maximal isotropic subspace of $\mathbb{H}$. In particular, since we assume that $L$ is involutive, it is an $A V$-Dirac structure.

Remark 7.3 Here we have relaxed the requirement $L \cap \bar{L}=0$ in the definition of a generalized complex structure. While we have allowed $L \cap \bar{L}$ to be non-trivial, it must lie in $j(V \otimes \operatorname{Ann}(H)) \subset V \otimes A^{*}$. As pointed out in Remark 2.2, this can be interpreted as saying that $L \cap \bar{L}$ only fails to be trivial up to an "infinitesimal". On the other hand, we still require that $L$ be an $A V$-Dirac structure.

This is in contrast to the approach taken by generalized CRF structures, introduced by Izu Vaisman in [26], which requires $L \cap \bar{L}=0$, but does not require $L$ to be a Dirac structure.

It is well known that one can canonically associate a Poisson structure with every generalized complex structure. The analogue for generalized CR structures is to endow $V \otimes A^{*}$ with a non-trivial Lie algebroid structure, which we shall do in a canonical fashion following the corresponding argument given for generalized complex structures in [12].

We have an inclusion $i: H \rightarrow A$, and consequently, a map $J \circ j \circ\left(i d \otimes i^{*}\right): V \otimes$ $A^{*} \rightarrow \mathbb{H}$, which (abusing notation), we shall simply call J . We consider the family of subspaces of $A$ given by

$$
D_{t}:=e^{t J J}\left(V \otimes A^{*}\right)+V \otimes \operatorname{Ann}(H)=q^{-1}\left(e^{t J}\left(V \otimes H^{*}\right)\right)
$$

Since $e^{t J}=\cos (t)+\sin (t) \mathbb{J}: \mathbb{H} \rightarrow \mathbb{H}$ is orthogonal, and $j\left(V \otimes H^{*}\right)$ is a lagrangian subspace of $\mathbb{H}$, it follows that $D_{t}$ is lagrangian for each $t$.

The following proposition is a slight generalization of a result of Gualtieri [12].

Proposition 7.4 (Gualtieri) The family $D_{t}$ of almost $A V$-Dirac structures is integrable for all $t$.
Proof Let $\xi_{1}, \xi_{2} \in \Gamma\left(V \otimes A^{*}\right)$, then since $V \otimes A^{*} \subset L \oplus \bar{L}$, we may choose $X_{j} \in \Gamma(L)$ and $Y_{j} \in \Gamma(\bar{L})$, such that $\xi_{j}=X_{j}+Y_{j}$. It follows that $J \xi_{j}=i X_{j}-i Y_{j}+V \otimes \operatorname{Ann}(H)$. In fact, since $L \cap \bar{L}=V \otimes \operatorname{Ann}(H)$, by choosing $X_{j}$ and $Y_{j}$ appropriately, we may suppose that $i X_{j}-i Y_{j}$ is any given representative of $\mathrm{J} \circ i^{*}\left(\xi_{j}\right)$ in $\pi^{-1}(H)$. Abusing notation, we will use the term $J\left(\xi_{j}\right)$ and our particular choice of representative $i X_{j}-$ $i Y_{j}$ interchangeably. Then,

$$
\begin{aligned}
\llbracket \| \xi_{1}, J \xi_{2} \rrbracket-\llbracket \xi_{1}, \xi_{2} \rrbracket & =\llbracket i X_{1}-i Y_{1}, i X_{2}-i Y_{2} \rrbracket-\llbracket X_{1}+Y_{1}, X_{2}+Y_{2} \rrbracket \\
& =-2 \llbracket X_{1}, X_{2} \rrbracket-2 \llbracket Y_{1}, Y_{2} \rrbracket
\end{aligned}
$$

and

$$
\begin{aligned}
\llbracket \| \xi_{1}, \xi_{2} \rrbracket-\llbracket \xi_{1}, J \xi_{2} \rrbracket & =\llbracket i X_{1}-i Y_{1}, X_{2}+Y_{2} \rrbracket-\llbracket X_{1}+Y_{1}, i X_{2}-i Y_{2} \rrbracket \\
& =2 i \llbracket X_{1}, X_{2} \rrbracket-2 i \llbracket Y_{1}, Y_{2} \rrbracket .
\end{aligned}
$$

Thus, since $L$ and hence $\bar{L}$ are involutive, we have $\llbracket J \xi_{1}, J \xi_{2} \rrbracket-\llbracket \xi_{1}, \xi_{2} \rrbracket+V \otimes \operatorname{Ann}(H)=$ $J\left(\llbracket J \xi_{1}, \xi_{2} \rrbracket-\llbracket \xi_{1}, J \xi_{2} \rrbracket\right)+V \otimes \operatorname{Ann}(H)$.

We let $a=\cos (t)$ and $b=\sin (t)$, and we have,

$$
\begin{aligned}
& \llbracket(a+b \rrbracket) \xi_{1},(a+b \rrbracket) \xi_{2} \rrbracket \\
& \left.=a b\left(\llbracket \xi_{1}, \| \xi_{2} \rrbracket+\llbracket\right\rfloor \xi_{1}, \xi_{2} \rrbracket\right)+b^{2} \llbracket J \xi_{1}, \mathrm{~J} \xi_{2} \rrbracket \\
& =a b\left(\llbracket \xi_{1}, \| \xi_{2} \rrbracket+\llbracket J \xi_{1}, \xi_{2} \rrbracket\right)+b^{2}\left(\llbracket J \xi_{1}, \| \xi_{2} \rrbracket-\llbracket \xi_{1}, \xi_{2} \rrbracket\right) .
\end{aligned}
$$

So modulo $V \otimes \operatorname{Ann}(H)$, we see that
$\llbracket(a+b J) \xi_{1},(a+b J) \xi_{2} \rrbracket+V \otimes \operatorname{Ann}(H)=b(a+b J)\left(\llbracket \xi_{1}, J \xi_{2} \rrbracket+\llbracket J \xi_{1}, \xi_{2} \rrbracket\right)+V \otimes \operatorname{Ann}(H)$
Since $\llbracket \xi_{1}, J \xi_{2} \rrbracket+\llbracket J \xi_{1}, \xi_{2} \rrbracket \in V \otimes A^{*}$, it follows that $(\cos (t)+\sin (t) J)\left(V \otimes A^{*}\right)+$ $V \otimes \operatorname{Ann}(H)$ is involutive.

We next consider the map $P: V \otimes A^{*} \rightarrow H \xrightarrow{i} A$, which for $\xi, \eta \in V \otimes A^{*}$, is given by

$$
\langle P(\xi), \eta\rangle=\left\langle\left.\frac{\partial}{\partial t}\right|_{t=0} e^{t J J}(\xi), \eta\right\rangle=\langle\mathbb{J} \xi, \eta\rangle \quad\left(=\left\langle i \circ \pi \circ J \circ j \circ i^{*}(\xi), \eta\right\rangle\right)
$$

Clearly, since $\mathbb{J}$ is an orthogonal almost complex structure on $\mathbb{H}, P$ will be given by an element of $\Gamma\left(V^{*} \otimes \wedge^{2} A\right)$, which we will also denote by $P$. Adapting a proposition given in [12], we have the following.
Proposition 7.5 (Gualtieri) The bivector field $P=i \circ \pi \circ \mathrm{JJ} \circ j \circ i^{*}: V \otimes A^{*} \rightarrow A$ defines a Lie algebroid structure on $V \otimes A^{*}$. The bracket is given by

$$
\left.[\xi, \eta]=\iota_{P(\cdot, \xi)} d \eta-\iota_{P(\cdot, \eta)} d \xi+d(P(\xi, \eta))\right)
$$

where $\xi, \eta \in V \otimes A^{*}$, and the anchor map is given by $\xi \rightarrow a \circ P(\xi, \cdot) V \otimes A^{*} \rightarrow T M$, where $a: A \rightarrow T M$ is the anchor map of $A$. Furthermore, the map $\xi \rightarrow a \circ P(\xi, \cdot): V \otimes$ $A^{*} \rightarrow A$ is a Lie-algebroid morphism.

The proof is an adaptation of one found in [12].
Proof We choose a splitting of the $A V$-Courant algebroid and use the isomorphism and notation described in Proposition 2.7. Then if we choose $t$ sufficiently small, the $A V$-Dirac structures $D_{t}$ can be described as the graphs of $\beta_{t} \in \Gamma\left(V^{*} \otimes \wedge^{2} A\right)$.

In [24], it was shown that the integrability condition of a twisted Poisson structure $\beta$ over a 3-form background $\gamma$ is $[\beta, \beta]=\wedge^{3} \tilde{\beta}(\gamma)$, where $\tilde{\beta}: T^{*} M \rightarrow T M$ is given by $\tilde{\beta}(\xi)(\eta)=\beta(\xi, \eta)$. We would like to derive a similar equation for $\beta_{t}$, but we have not defined a bracket for sections of $V^{*} \otimes \Lambda^{2} A$. In order to define such a bracket, we first define a sheaf of rings over $M$.

We let $\mathcal{F}:=\left(S(V) \otimes S\left(V^{*}\right)\right) / I$, where $S(V)$ denotes the symmetric algebra generated by $V$, and $I$ is the ideal generated by $u \otimes f-f(u)$ for $f \in \Gamma\left(V^{*}\right)$ and $u \in \Gamma(V)$. Since $V$ is one dimensional, if $t \in \Gamma(V)$ is a local basis, then $\mathcal{F}$ is locally isomorphic to $C^{\infty}(M)\left[t, t^{-1}\right]$ as a ring. It is clear that it has a well-defined $\mathbb{Z}$ grading, which for a homogeneous $v \in \mathcal{F}$, we denote by $\tilde{v}$.
$\Gamma\left(S(V) \otimes S\left(V^{*}\right)\right)$ is a $\Gamma(A)$ module, where sections of $\Gamma(A)$ act as derivations, and it is easy to check that $\Gamma(I)$ is a sub-module. Thus it is clear that $\Gamma(A)$ acts on $\Gamma(\mathcal{F})$ by derivations satisfying the Leibniz rule with respect to the ring structure on $\mathcal{F}$.

We define a bracket on $\mathcal{F} \otimes \wedge^{*} A$, as follows (for $v, w \in \Gamma(\mathcal{F})$ and $P, Q \in \Gamma\left(\wedge^{*} A\right)$ ):

- $[X, v]=X v$ for any $X \in \Gamma(A)$, and $[v, w]=0$;
- $[P \wedge Q, v]=P \wedge[Q, v]+(-1)^{|Q|}[P, v] \wedge Q$;
- $[P, Q]$ is given by the Schouten-Nijenhuis bracket;
- $[v P, w Q]=(v[P, w]) Q-(-1)^{(|P|-1)(|Q|-1)}(w[Q, v]) P+v w[P, Q]$.

If we write $|v P|=i$ for $P \in \wedge^{i} A$, and $\operatorname{deg}(v P)=(\tilde{v},|v P|)$, then it is clear that our bracket satisfies the following identities (for homogeneous $a, b, c \in \Gamma\left(\mathcal{F} \otimes \wedge^{*} A\right)$ ):

- $\operatorname{deg}(a b)=\operatorname{deg}(a)+\operatorname{deg}(b)$ and $\operatorname{deg}([a, b])=\operatorname{deg}(a)+\operatorname{deg}(b)-(0,1)$;
- $(a b) c=a(b c)$ and $a b=(-1)^{|a||b|} b a$;
- $[a, b c]=[a, b] c+(-1)^{(|a|-1)|b|} b[a, c]$;
- $[a, b]=-(-1)^{(|a|-1)(|b|-1)}[b, a]$;
- $[a,[b, c]]=[[a, b], c]+(-1)^{(|a|-1)(|b|-1)}[b,[a, c]]$.

We next extend $d$ to a map $d: \mathcal{F} \otimes \wedge^{i} A^{*} \rightarrow \mathcal{F} \otimes \wedge^{i+1} A^{*}$ in the obvious way. We also have a natural $\mathcal{F}$-bilinear pairing on $\Gamma\left(\mathcal{F} \otimes \wedge^{*} A^{*}\right) \times \Gamma\left(\mathcal{F} \otimes \wedge^{*} A\right)$, which for $v_{i}, w_{j} \in \mathcal{F}, \alpha_{i} \in \Gamma\left(A^{*}\right)$, and $X_{j} \in \Gamma(A)$, is given by

$$
\left\langle\left(v_{1} \otimes \alpha_{1}\right) \cdots\left(v_{p} \otimes \alpha_{p}\right),\left(w_{1} \otimes X_{1}\right) \cdots\left(w_{q} \otimes X_{q}\right)\right\rangle= \begin{cases}0 & \text { if } p \neq q \\ \operatorname{det}\left(v_{i} w_{j} \otimes \alpha_{i}\left(X_{j}\right)\right) & \text { if } p=q\end{cases}
$$

We define a morphism $\iota: \mathcal{F} \otimes \wedge^{*} A \rightarrow \operatorname{End}\left(\mathcal{F} \otimes \wedge^{*} A^{*}\right)$ by $\langle\xi, P Q\rangle=\left\langle\iota_{P} \xi, Q\right\rangle$. For $P \in \mathcal{F} \otimes A, \iota_{P}$ is a derivation.

We also define a morphism $\breve{\iota}: \mathcal{F} \otimes \wedge^{*} A^{*} \rightarrow \operatorname{End}\left(\mathcal{F} \otimes \wedge^{*} A\right)$ by $\langle\xi \eta, P\rangle=\langle\xi, \breve{\iota}(\eta) P\rangle$. For $\alpha \in \mathcal{F} \otimes A^{*}, \breve{\iota}(\alpha)$ is a derivation on the right. Namely, $\breve{\iota}(\alpha)(P Q)=P \breve{\iota}(\alpha) Q+$ $(-1)^{|Q|}(\breve{\iota}(\alpha) P) Q$ (where $P, Q \in \mathcal{F} \otimes \wedge^{*} A$ are homogeneous).

Next, we notice that $\iota_{[P, Q]}=-\left[\left[\iota_{Q}, d\right], \iota_{P}\right]$. This is easy to check, following the argument given in [20]. Also following an argument in [20] one can verify that, for
$\eta \in \Gamma\left(\mathcal{F} \otimes A^{*}\right)$,

$$
\begin{aligned}
& \breve{\iota}(\eta)[P, Q]-[P, \breve{\iota}(\eta) Q]-(-1)^{|Q|-1}[\breve{\iota}(\eta) P, Q]= \\
& \quad(-1)^{|Q|-2}(\breve{\iota}(d \eta)(P Q)-P \breve{\iota}(d \eta) Q-(\breve{\iota}(d \eta) P) Q) .
\end{aligned}
$$

From this, we calculate, for any $\beta \in \Gamma\left(\mathcal{F} \otimes \wedge^{2} A\right)$ and $\xi, \eta \in \Gamma\left(\mathcal{F} \otimes A^{*}\right)$,

$$
\begin{aligned}
{[\breve{\iota}(\xi) \beta, \breve{\iota}(\eta) \beta]=\frac{1}{2} \breve{\iota}(\xi \eta)[\beta, \beta]+[\beta,\langle\eta \xi, \beta\rangle]+} & \frac{1}{2}\left(\breve{\iota}(\eta d \xi) \beta^{2}-\breve{\iota}(\xi d \eta) \beta^{2}\right) \\
& -\langle d \xi, \beta\rangle \breve{\iota}(\eta) \beta+\langle d \eta, \beta\rangle \breve{\iota}(\xi) \beta .
\end{aligned}
$$

Furthermore, it is not difficult to verify that $[\beta,\langle\eta \xi, \beta\rangle]=\breve{\iota}(d \beta(\eta, \xi)) \beta$, while

$$
\frac{1}{2}\left(\breve{\iota}(\eta d \xi) \beta^{2}-\breve{\iota}(\xi d \eta) \beta^{2}\right)-\langle d \xi, \beta\rangle \breve{\iota}(\eta) \beta+\langle d \eta, \beta\rangle \breve{\iota}(\xi) \beta=\breve{\iota}\left(\iota_{\breve{\iota}(\xi) \beta} d \eta-\iota_{\breve{\iota}(\eta) \beta} d \xi\right) \beta .
$$

Thus, we have, for $\beta \in \Gamma\left(V^{*} \otimes \wedge^{2} A\right)$,

$$
\begin{aligned}
\llbracket- & \breve{\iota}(\xi) \beta+\xi,-\breve{\iota}(\eta) \beta+\eta \rrbracket_{\phi} \\
= & {[\breve{\iota}(\xi) \beta, \breve{\iota}(\eta) \beta]-\iota_{\breve{\iota}(\xi) \beta} d \eta+\iota_{\breve{\iota}(\eta) \beta} d \xi+d(\beta(\xi, \eta))+\iota_{\breve{\iota}(\xi) \beta} \iota_{\breve{\iota}(\eta) \beta} H } \\
= & \breve{\iota}\left(\iota_{\breve{\iota}(\xi) \beta} d \eta-\iota_{\breve{\iota}(\eta) \beta} d \xi-d(\beta(\xi, \eta))\right) \beta-\iota_{\breve{\iota}(\xi) \beta} d \eta+\iota_{\breve{\iota}(\eta) \beta} d \xi+d(\beta(\xi, \eta)) \\
& +\frac{1}{2} \breve{\iota}(\xi \eta)[\beta, \beta]+\iota_{\breve{\iota}(\xi) \beta} \iota_{\breve{\iota}(\eta) \beta} H .
\end{aligned}
$$

It follows that $\beta_{t}$ defines an $A V$-Dirac structure under our chosen splitting if and only if $\frac{1}{2} \breve{\iota}(\eta \xi)\left[\beta_{t}, \beta_{t}\right]=\breve{\iota}\left(\iota_{\breve{l}(\xi) \beta_{t}} \iota_{\breve{\iota}(\eta) \beta_{t}} H\right) \beta_{t}$. To rewrite this, we let $\tilde{\beta}: \mathcal{F} \otimes A^{*} \rightarrow \mathcal{F} \otimes A$ be the map $\alpha \rightarrow-\breve{\iota}(\alpha) \beta$. The condition is then $\left[\beta_{t}, \beta_{t}\right]=2 \wedge^{3} \tilde{\beta}_{t}(H)$. We differentiate both sides by $t$ and evaluate at 0 . Since we have $P=\left.\frac{\partial}{\partial t}\right|_{0} \beta_{t}$ and $\beta_{0}=0$, the cubic term vanishes, and we see that the condition is $[P, P]=0$. The result follows immediately from this.

We also have a bracket $\{\cdot, \cdot\}$ on $\Gamma(V)$, which for $v, w \in \Gamma(V)$ is given by

$$
\begin{equation*}
\{v, w\}=P(d v, d w) \tag{7.1}
\end{equation*}
$$

It satisfies the following properties (for $f \in C^{\infty}(M)$ ):

- $\{\cdot, \cdot\}$ is bilinear;
- $\{v, w\}=-\{w, v\}$;
- $\{v, f w\}=f\{v, w\}+(a \circ P(d v)(f)) w ;$
- $\{u,\{v, w\}\}=\{\{u, v\}, w\}+\{v,\{u, w\}\}$ (for any $u, v, w \in \Gamma(V)$ ).

Since $V$ is a line-bundle, this is quite similar to a Poisson structure. In particular, if $U \subset M$ is an open set on which $\sigma \in \Gamma\left(\left.V\right|_{U}\right)$ is a local basis such that $P(\sigma)=0$, then we have a morphism

$$
\rho: C^{\infty}(U) \xrightarrow{f \rightarrow f \sigma} \Gamma\left(\left.V\right|_{U}\right)
$$

which allows us to define a Poisson structure on $U$, by

$$
\{f, g\}=\rho^{-1}\{\rho(f), \rho(g)\}
$$

In particular, if in some neighborhood $U \subset M, V$ admits a non-zero $A$-parallel section $\sigma \in \Gamma\left(\left.V\right|_{U}\right)$, then $P(\sigma)=0$, and thus $U$ is endowed with a Poisson structure. In fact, the Poisson structure associated with $U$ in this way is unique up to a constant multiple. Furthermore, if it exists at one point on a leaf of $A$, then it exists for any neighborhood of any point in that leaf.

Remark 7.6 (Poisson Structure on a Leaf of $A$ ) Suppose that $F \subset M$ is a connected leaf of the foliation given by $A$, then $a:\left.A\right|_{F} \rightarrow T F$ is a Lie algebroid, and we have an exact sequence of Lie algebroids given by $0 \rightarrow L=\left.\operatorname{ker}(a) \rightarrow A\right|_{F} \rightarrow T F \rightarrow 0$, where $L$ is actually a bundle of Lie algebras. The following are equivalent:

- $V$ admits an $\left.A\right|_{F}$-parallel section for any neighborhood $U \subset F$;
- $L$ acts trivially on $\left.V\right|_{F}$;
- $L_{x}$ acts trivially on $V_{x}$, for some point $x \in F{ }^{6}$

Note that, up to a constant multiple, there is a unique $A$-parallel section of $\left.V\right|_{F}$. Thus, if $\sigma \in \Gamma\left(\left.V\right|_{F}\right)$ is a non-zero $A$-parallel section, we can associate a Poisson structure with $F$, unique up to a constant multiple.

Remark 7.7 (Jacobi Bundle) A Jacobi bundle, introduced by Marle in [21] and Kirillov in [16], is a line bundle $P \rightarrow M$ over a manifold $M$, together with a bilinear map $\{\cdot, \cdot\}: \Gamma(P) \times \Gamma(P) \rightarrow \Gamma(P)$ on the sections of $P$ and a map $\Gamma(P) \xrightarrow{s \rightarrow X_{s}} \Gamma(T M)$ such that

- $\{\cdot, \cdot\}$ is bilinear;
- $\{v, w\}=-\{w, v\}$ (for any $v, w \in \Gamma(P))$;
- $\{v, f w\}=f\{v, w\}+\left(X_{v}(f)\right) w$ (for any $f \in C^{\infty}(M)$ and $\left.v, w \in \Gamma(P)\right)$;
- $\{u,\{v, w\}\}=\{\{u, v\}, w\}+\{v,\{u, w\}\}$ (for any $u, v, w \in \Gamma(P)$ ).

It follows that $V$ together with the bracket (7.1) is a Jacobi bundle canonically associated with the generalized CR structure.

Suppose for some $U \subset M$ there is a choice of a local basis $\sigma \in \Gamma\left(\left.V\right|_{U}\right)$. We may consider the isomorphism

$$
\rho: C^{\infty}(U) \xrightarrow{f \rightarrow f \sigma} \Gamma\left(\left.V\right|_{U}\right)
$$

which allows us to define a bracket on $C^{\infty}(U)$ by $[f, g]_{\sigma}=\rho^{-1}\{\rho(f), \rho(g)\}$. One notices that this bracket endows $C^{\infty}(U)$ with a Lie algebra structure that is local in the sense that the linear operator

$$
D_{f}: C^{\infty}(U) \xrightarrow{g \rightarrow[f, g]_{\sigma}} C^{\infty}(U)
$$

[^4]is local for all $f \in C^{\infty}(U)$. It is an important result (see [11, 16, 25]) that for any local Lie algebra structure, there exists unique $\Lambda \in \Gamma\left(\wedge^{2} T M\right)$, and $E \in \Gamma(T M)$ with $[\Lambda, \Lambda]=-2 \Lambda \wedge E$ and $[\Lambda, E]=0$ such that
$$
[f, g]_{\sigma}=\{f, g\}_{\Lambda}+f \mathcal{L}_{X} g-g \mathcal{L}_{X} f
$$
where $\{f, g\}_{\Lambda}=\breve{\iota}_{d f} \breve{\iota}_{d g} \Lambda$.
The triple $(U, \Lambda, E)$ is then called a Jacobi structure. Note however the dependence of $\Lambda$ and $E$ on $\sigma$; this is unlike the local Poisson structure that (if it exists) is unique up to a constant multiple.

Example 7.8 (CR Structures) As described in Section6, a CR-structure on a manifold $M$ can be described by a generalized CR structure. In this case, $V$ can be taken to be the trivial bundle, and $A$ can be taken to be TM. It follows from the above discussion that there is a Poisson structure $P \in \Gamma\left(\wedge^{2} T M\right)$ associated with the CR structure.

If $L \subset \mathbb{C} \otimes T M$ is the CR-structure, and $H=\operatorname{Re}(L \oplus \bar{L}) \subset T M$, then $P\left(T^{*} M\right) \subset H$. So the symplectic foliation associated with $P$ is everywhere tangent to $H$.

Example 7.9 (Quotients of Generalized Complex Structures) If the procedures described in Examples 4.3 and 3.4 are applied to a generalized complex structure, then one obtains a generalized CR structure.

Example 7.10 (Contact Structures and Generalized Contact Structures) Suppose that $M$ is a contact manifold, then there is a canonical way to associate a generalized CR structure with $M$. In particular, if $N=M \times \mathbb{R}$ is its symplectization, then $N$ admits a generalized complex structure corresponding to its symplectic structure. $\mathbb{R}$ acts on $N$, and the quotient is a generalized CR structure on $M$ (in the sense of Examples 3.7 and 4.4).

This procedure is also described in $[14,15]$, where they describe it as a generalized contact structure. In fact any generalized contact structure results from the quotient of generalized complex structure, and as such can also be described as a generalized CR structure.

Since the Lie algebroid $A$ and the vector bundle $V$ describe an $\mathcal{E}^{1}(M)$ structure, as given in Example 3.3, it can be checked that $V$ does not admit parallel sections, and thus, in general, $P \in \Gamma\left(V^{*} \otimes \wedge^{2} A\right)$ does not describe a Poisson structure, but rather a Jacobi structure. When the generalized contact structure is simply a contact structure, then $P$ corresponds to a Jacobi structure describing the contact structure.

To be more explicit, we let $M$ be a contact manifold with contact distribution $\xi \subset T M$, and $N=M \times \mathbb{R}$ its symplectization, where we let $t: M \times \mathbb{R} \rightarrow \mathbb{R}$ be the projection to the second factor, and $\omega \in \Omega^{2}(N)$ denote the corresponding symplectic form. (That is, $\omega=e^{t}(d \eta+d t \wedge \eta)$, where $\eta \in \operatorname{Ann}(\xi)$ is nowhere vanishing.) We note that $\mathcal{L}_{\frac{\partial}{\partial t}} \omega=\omega$.

Since $N$ is a symplectic manifold, we can associate a canonical generalized complex structure $\mathcal{J}: T N \oplus T^{*} N \rightarrow T N \oplus T^{*} N$ with it on the standard Courant algebroid

$$
0 \rightarrow T^{*} N \rightarrow T N \oplus T^{*} N \rightarrow T N \rightarrow 0
$$

(see [12] for details).
The Poisson bivector $\pi \in \Gamma\left(\wedge^{2} T N\right)$ associated with this generalized complex structure has the property that $\mathcal{L}_{\partial / \partial t} \pi=-\pi$ (since it is the Poisson bivector corresponding to $\omega$ ). It follows that we can write $\pi=e^{-t}(\Lambda+\partial / \partial t \wedge E)$ for $E \in \Gamma(M)$, and $\Lambda \in \Gamma\left(\wedge^{2} M\right)$. Then $[\pi, \pi]=0$ implies that

$$
\begin{aligned}
0 & =[\pi, \pi]=\left[e^{-t}\left(\Lambda+\frac{\partial}{\partial t} \wedge E\right), e^{-t}\left(\Lambda+\frac{\partial}{\partial t} \wedge E\right)\right] \\
& =e^{-2 t}[\Lambda, \Lambda]-2 e^{-2 t} \Lambda \wedge E+2 e^{-2 t} \frac{\partial}{\partial t} \wedge[\Lambda, E]
\end{aligned}
$$

From this it follows that $[\Lambda, \Lambda]=-2 \Lambda \wedge E$ and $[\Lambda, E]=0$, which are the defining conditions for a Jacobi structure $(\Lambda, E)$ on $M$.

Now, we consider the $T M \oplus \mathbb{R}-\mathbb{R}$ Courant algebroid structure on $M$, given by taking the quotient by the $G=\mathbb{R}$ action on $N=M \times \mathbb{R}$,

$$
0 \rightarrow T^{*} N / G \rightarrow\left(T N \oplus T^{*} N\right) / G \rightarrow T N / G \rightarrow 0
$$

and the generalized CR structure on $M$ given by quotient homomorphism

$$
\mathrm{J}:=\mathcal{J} / G:\left(T N \oplus T^{*} N\right) / G \rightarrow\left(T N \oplus T^{*} N\right) / G
$$

They define an $A V$-Courant algebroid, where $A=T N / G$, and the bundle $V \rightarrow$ $M$ is trivial, with $\Gamma(V) \simeq C^{\infty}(N)^{G}$ (this is in fact an $\mathcal{E}^{1}(M)$ structure; see [14]). Abusing notation, we denote by $e^{t} \in \Gamma(V)$ the section associated with the $G$-invariant function $e^{t} \in C^{\infty}(N)$.

Then the bivector $P \in \Gamma\left(V^{*} \otimes \wedge^{2} A\right)$ associated with the generalized CR structure on $M$ is simply $e^{-t}\left(\Lambda+\frac{\partial}{\partial t} \wedge E\right)$, and it defines a Jacobi structure on $M$, with bivector field $\Lambda$ and vector field $E$. Since $\Lambda^{n} \wedge E \neq 0$ (where $\operatorname{dim}(M)=2 n+1$ ), this Jacobi structure corresponds to a contact structure. In fact, the contact distribution is given by $\operatorname{span}\left\{\breve{\iota}_{\alpha} \Lambda \mid \alpha \in T * M\right\}$, and if $\theta \in \Omega^{1}(M)$ satisfies $\breve{\iota}_{\theta} \Lambda=0$ and $\breve{\iota}_{\theta} E=1$, then $\theta$ is a contact form. It is not difficult to see that this is the original contact structure, $\xi$, defined on $M$. (In fact, if $\omega=e^{t}(d \eta+d t \wedge \eta$ ) is the symplectic form on $N$ (where $\eta \in \operatorname{Ann}(\xi)$ is nowhere vanishing), then $E$ is a reeb vector field for $\eta$ and $\theta=\eta$.)

We must note that, if instead of trivializing $V$ by the section $e^{t} \in \Gamma(V)$, we made the transformation $e^{t} \rightarrow f e^{t}$, for some nowhere vanishing $f \in C^{\infty}(M)$, then the appropriate changes to the Jacobi structure would be $\Lambda \rightarrow f \Lambda, E \rightarrow f E-\breve{\iota}_{d f} \Lambda$, and the transformation for the contact form would be $\theta \rightarrow \frac{1}{f} \theta$. Thus it is clear that the freedom to modify the trivializing section of $V$ by a scalar multiple does not change the contact distribution and fully accounts for the freedom to change the contact form by a scalar multiple. Indeed the generalized CR structure is defined intrinsically.

## A Appendix: Proof of Proposition 2.7

Suppose that $M$ is a manifold, $A$ is a Lie algebroid over $M, V$ is an $A$-module over $M$, and $A$ is an $A V$-Courant algebroid over $M$.

For $X, Y \in \Gamma(A)$, we have the following identities:

- $\left[\iota_{X}, \iota_{Y}\right]=0$;
- $\left[d, \iota_{X}\right]=\mathcal{L}_{X} ;$
- $\left[\mathcal{L}_{X}, \iota_{Y}\right]=\iota_{[X, Y]} ;$
- $[d, d]=0$;
- $\left[\mathcal{L}_{X}, d\right]=0$;
- $\left[\mathcal{L}_{X}, \mathcal{L}_{Y}\right]=\mathcal{L}_{[X, Y]}$.

We will provide the proof we promised for Proposition 2.7, which we restate here.
Proposition A. 1 Let $\phi: A \rightarrow \mathbb{A}$ be an isotropic splitting. Then under the isomorphism $\phi \oplus j: A \oplus\left(V \otimes A^{*}\right) \rightarrow \mathbb{A}$, the bracket is given by

$$
\llbracket X+\xi, Y+\eta \rrbracket_{\phi}=[X, Y]+\mathcal{L}_{X} \eta-\iota_{Y} d \xi+\iota_{X} \iota_{Y} H_{\phi}
$$

where $X, Y \in \Gamma(A), \xi, \eta \in \Gamma\left(V \otimes A^{*}\right)$, and $H_{\phi} \in \Gamma\left(V \otimes \wedge^{3} A^{*}\right)$, with $d H_{\phi}=0$.
Furthermore, if $\psi: A \rightarrow \mathbb{A}$ is a different choice of isotropic splitting, then $\psi(X)=$ $\phi(X)+j\left(\iota_{X} \beta\right)$, and $H_{\psi}=H_{\phi}-d \beta$, where $\beta \in \Gamma\left(V \otimes \wedge^{2} A^{*}\right)$.

Proof The proof will follow immediately from the following lemmas.
Lemma A. 2 If $\xi \in \Gamma\left(V \otimes A^{*}\right)$ and $e \in \Gamma(\mathbb{A})$, then $\llbracket e, j(\xi) \rrbracket=j\left(\mathcal{L}_{\pi(e)} \xi\right)$.
Proof Let $e_{1}, e_{2} \in \Gamma(\mathbb{A}), \xi \in \Gamma\left(V \otimes A^{*}\right)$,

$$
\begin{aligned}
\left\langle\llbracket e_{1}, j(\xi) \rrbracket, e_{2}\right\rangle & =\mathcal{L}_{\pi\left(e_{1}\right)}\left\langle j(\xi), e_{2}\right\rangle-\left\langle j(\xi), \llbracket e_{1}, e_{2} \rrbracket\right\rangle \\
& =\mathcal{L}_{\pi\left(e_{1}\right)} \iota_{\pi\left(e_{2}\right)} \xi-\iota_{\pi\left(\left[e_{1}, e_{2}\right]\right)} \xi=\mathcal{L}_{\pi\left(e_{1}\right)} \iota_{\pi\left(e_{2}\right)} \xi-\iota_{\left[\pi\left(e_{1}\right), \pi\left(e_{2}\right)\right]} \xi \\
& =\mathcal{L}_{\pi\left(e_{1}\right)} \iota_{\pi\left(e_{2}\right)} \xi-\left[\mathcal{L}_{\pi\left(e_{1}\right)}, \iota_{\pi\left(e_{2}\right)}\right] \xi=\iota_{\pi\left(e_{2}\right)} \mathcal{L}_{\pi\left(e_{1}\right)} \xi=\left\langle j\left(\mathcal{L}_{\pi\left(e_{1}\right)} \xi\right), e_{2}\right\rangle .
\end{aligned}
$$

Lemma A. 3 If $\xi \in \Gamma\left(V \otimes A^{*}\right)$ and $e \in \Gamma(\mathbb{A})$, then $\llbracket j(\xi), e \rrbracket=-j\left(\iota_{\pi(e)} d \xi\right)$.
Proof

$$
\begin{aligned}
\llbracket j(\xi), e \rrbracket & =D\langle j(\xi), e\rangle-\llbracket e, j(\xi) \rrbracket=j\left(d \iota_{\pi(e)} \xi\right)-j\left(\mathcal{L}_{\pi(e)} \xi\right) \\
& =j\left(d \iota_{\pi(e)} \xi-\left(\iota_{\pi(e)} d \xi+d \iota_{\pi(e)} \xi\right)\right)=-j\left(\iota_{\pi(e)} d \xi\right)
\end{aligned}
$$

Lemma A. 4 If $\phi: A \rightarrow \mathbb{A}$ is an isotropic splitting and if $X, Y \in \Gamma(A)$, then

$$
\llbracket \phi(X), \phi(Y) \rrbracket-\phi([X, Y])=j\left(\iota_{X} \iota_{Y} H\right)
$$

where $H \in \Gamma\left(V \otimes \wedge^{3} A^{*}\right)$.
Proof Let $\phi$ be an isotropic splitting, and $X, Y, Z \in \Gamma(A)$. Then

$$
\pi(\llbracket \phi(X), \phi(Y) \rrbracket-\phi([X, Y]))=0
$$

so by exactness of the sequence (2.2), $\llbracket \phi(X), \phi(Y) \rrbracket-\phi([X, Y]) \in j\left(\Gamma\left(V \otimes A^{*}\right)\right)$. We define $H$ by

$$
H(X, Y, Z)=\langle\phi(Z), \llbracket \phi(X), \phi(Y) \rrbracket-\phi([X, Y])\rangle=\langle\phi(Z), \llbracket \phi(X), \phi(Y) \rrbracket\rangle
$$

where the second equality follows since $\phi$ is an isotropic splitting. It is obvious that $H$ is tensorial in $Z$. Furthermore, making repeated use of the fact that $\phi$ is an isotropic splitting, we check that $H$ is skew-symmetric:

$$
\begin{aligned}
\langle\phi(Z), \llbracket \phi(X), \phi(Y) \rrbracket\rangle & =\langle\phi(Z),-\llbracket \phi(Y), \phi(X) \rrbracket+D\langle\phi(X), \phi(Y)\rangle\rangle \\
& =-\langle\phi(Z), \llbracket \phi(Y), \phi(X) \rrbracket\rangle
\end{aligned}
$$

and

$$
0=\mathcal{L}_{X}\langle\phi(Z), \phi(Y)\rangle=\langle\llbracket \phi(X), \phi(Z) \rrbracket, \phi(Y)\rangle+\langle\phi(Z), \llbracket \phi(X), \phi(Y) \rrbracket\rangle
$$

It follows that $H \in \Gamma\left(V \otimes \wedge^{3} A^{*}\right)$.
Lemma A. 5 Using the notation of the previous lemmas, $d H=0$.
Proof Using the fact that $\left[\mathcal{L}_{X}, \iota_{Y}\right]=\iota_{[X, Y]}$, it is easy to show that

$$
d \iota_{Z} \iota_{Y} \iota_{X}+\iota_{Z} \iota_{Y} \iota_{X} d=\mathcal{L}_{Z} \iota_{Y} \iota_{X}+\mathcal{L}_{Y} \iota_{X} \iota_{Z}+\mathcal{L}_{X} \iota_{Z} \iota_{Y}+\iota_{Z} \iota_{[Y, X]}+\iota_{Y} \iota_{[X, Z]}+\iota_{X} \iota_{[Z, Y]} .
$$

Let $\phi: A \rightarrow \mathbb{A}$ be an isotropic splitting. We shall use the identification

$$
A \oplus\left(V \otimes A^{*}\right) \xrightarrow{\phi \oplus j} \mathbb{A}
$$

explicitly throughout this section. We have, for $X, Y, Z \in \Gamma(A)$,

$$
\llbracket X, Y \rrbracket_{\phi}=[X, Y]+\iota_{X} \iota_{Y} H
$$

Then using Axiom (AV-1) from the definition of an $A V$-Courant algebroid, we see that

$$
\begin{aligned}
0= & \llbracket Z, \llbracket Y, X \rrbracket_{\phi} \rrbracket_{\phi}-\llbracket \llbracket Z, Y \rrbracket_{\phi}, X \rrbracket_{\phi}-\llbracket Y, \llbracket Z, X \rrbracket_{\phi} \rrbracket_{\phi} \\
= & \llbracket Z,[Y, X]+\iota_{Y} \iota_{X} H \rrbracket_{\phi}-\llbracket[Z, Y]+\iota_{Z} \iota_{Y} H, X \rrbracket_{\phi}-\llbracket Y,[Z, X]+\iota_{Z} \iota_{X} H \rrbracket_{\phi} \\
= & \llbracket Z,[Y, X] \rrbracket_{\phi}+\mathcal{L}_{Z} \iota_{Y} \iota_{X} H-\llbracket[Z, Y], X \rrbracket_{\phi}+\iota_{X} d \iota_{Z} \iota_{Y} H-\llbracket Y,[Z, X] \rrbracket_{\phi}-\mathcal{L}_{Y} \iota_{Z} \iota_{X} H \\
= & \llbracket Z,[Y, X] \rrbracket_{\phi}-\llbracket[Z, Y], X \rrbracket_{\phi}-\llbracket Y,[Z, X] \rrbracket_{\phi} \\
& +\mathcal{L}_{Z} \iota_{Y} \iota_{X} H+\mathcal{L}_{X} \iota_{Z} \iota_{Y} H+\mathcal{L}_{Y} \iota_{X} \iota_{Z} H-d \iota_{Z} \iota_{Y} \iota_{X} H \\
= & {[Z,[Y, X]]+\iota_{Z} \iota_{[Y, X]} H-[[Z, Y], X]-\iota_{[Z, Y]} \iota_{X} H-[Y,[Z, X]]-\iota_{Y} \iota_{[Z, X]} H } \\
& +\mathcal{L}_{Z} \iota_{Y} \iota_{X} H+\mathcal{L}_{X} \iota_{Z} \iota_{Y} H+\mathcal{L}_{Y} \iota_{X} \iota_{Z} H-d \iota_{Z} \iota_{Y} \iota_{X} H \\
= & {[Z,[Y, X]]-[[Z, Y], X]-[Y,[Z, X]]+\iota_{Z} \iota_{[Y, X]} H+\iota_{X} \iota_{[Z, Y]} H+\iota_{Y} \iota_{[X, Z]} H } \\
& +\mathcal{L}_{Z} \iota_{Y} \iota_{X} H+\mathcal{L}_{X} \iota_{Z} \iota_{Y} H+\mathcal{L}_{Y} \iota_{X} \iota_{Z} H-d \iota_{Z} \iota_{Y} \iota_{X} H \\
= & \iota_{Z} \iota_{Y} \iota_{X} d H .
\end{aligned}
$$

Lemma A. 6 Let $\phi: A \rightarrow \mathbb{A}$ and $\psi: A \rightarrow \mathbb{A}$ be two isotropic splittings, and let $H_{\phi}$ and $H_{\psi}$ be the elements of $\Gamma\left(V \otimes \wedge^{3} A^{*}\right)$ associated with the corresponding splittings. Namely, if $X, Y \in \Gamma(A)$, then $\llbracket \phi(X), \phi(Y) \rrbracket=\phi([X, Y])+j \iota_{X} \iota_{Y} H_{\phi}$, and similarly for $H_{\psi}$.

Then there exists $\beta \in \Gamma\left(V \otimes \wedge^{2} A^{*}\right)$ such that $\psi(X)=\phi(X)+j\left(\iota_{X} \beta\right)$ and $H_{\psi}=$ $H_{\phi}-d \beta$.

Proof Since $\phi$ and $\psi$ are splittings, we see that

$$
\pi((\phi-\psi)(X))=0
$$

Thus, by the exactness of the sequence (2.2), $(\phi-\psi)(X)=j \circ S(X)$ for some linear map $S: A \rightarrow V \otimes A^{*}$.

However since the splittings are isotropic,

$$
\begin{aligned}
0 & =\langle\phi(X), \phi(Y)\rangle \\
& =\langle\psi(X)+j \circ S(X), \psi(Y)+j \circ S(Y)\rangle \\
& =S(X)(Y)+S(Y)(X),
\end{aligned}
$$

so we can define $\beta \in \Gamma\left(V \otimes \wedge^{2} A^{*}\right)$ by $\iota_{X} \beta=S(X)$. Then, we see that

$$
\begin{aligned}
\psi([X, Y])+\iota_{X} \iota_{Y} H_{\psi} & =\llbracket \phi(X)+j\left(\iota_{X} \beta\right), \phi(Y)+j\left(\iota_{Y} \beta\right) \rrbracket \\
& =\phi([X, Y])+j\left(\mathcal{L}_{X} \iota_{Y} \beta-\iota_{Y} d \iota_{X} \beta+\iota_{X} \iota_{Y} H_{\phi}\right) \\
& =\phi([X, Y])+j\left(\iota_{X} \iota_{Y} H_{\phi}\right)+j\left(\mathcal{L}_{X} \iota_{Y} \beta-\iota_{Y} \mathcal{L}_{X} \beta+\iota_{Y} \iota_{X} d \beta\right) \\
& =\phi([X, Y])+j\left(\iota_{X} \iota_{Y} H_{\phi}\right)+j\left(\iota_{[X, Y]} \beta+\iota_{Y} \iota_{X} d \beta\right) \\
& =\phi([X, Y])+j\left(\iota_{[X, Y]} \beta\right)+j\left(\iota_{X} \iota_{Y} H_{\phi}-\iota_{X} \iota_{Y} d \beta\right) \\
& =\psi([X, Y])+j\left(\iota_{X} \iota_{Y} H_{\phi}-\iota_{X} \iota_{Y} d \beta\right),
\end{aligned}
$$

so we have $H_{\psi}=H_{\phi}-d \beta$.
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## References

[1] M. Audin, A. C. da Silva, and E. Lerman, Symplectic geometry of integrable Hamiltonian systems. Lectures delivered at the Euro Summer School held in Barcelona, July 10-15, 2001. Advanced Courses in Mathematics, CRM Barcelona, Birkhäuser Verlag, Basel, 2003.
[2] P. Bressler and A. Chervov, Courant algebroids. Geometry J. Math. Sci. 128(2005), no. 4, 3030-3053. doi:10.1007/s10958-005-0251-7
[3] Z. Chen and Z.-J. Liu, Omni-Lie alegebroids. J. Geom. Phys. 60(2010), no. 5, 799-808. doi:10.1016/j.geomphys.2010.01.007
[4] Z. Chen, Z.-J. Liu, and Y.-H. Sheng, Dirac structures of omni-Lie algebroids. arXiv:0802.3819v2 [math.DG].
[5] $\longrightarrow$, E-Courant algebroids. arXiv:0805.4093v2 [math.DG].
[6] T. Courant and A. Weinstein, Beyond Poisson structures. In: Action hamiltoniennes de groupes. Troisième théorème de Lie (Lyon, 1986), Travaux en Cours, 27, Hermann, Paris, 1988, pp. 39-49.
[7] T. Courant Dirac manifolds. Trans. Amer. Math. Soc. 319(1990), no. 2, 631-661. doi:10.2307/2001258
[8] M. Crainic and R. L. Fernandes, Integrability of Lie brackets. Ann. of Math. 157(2003), no. 2, 575-620. doi:10.4007/annals.2003.157.575
[9] A. C. da Silva and A. Weinstein, Geometric models for noncommutative algebras. Berkeley Mathematical Lecture Notes, 10, American Mathematical Society, Providence, RI, 2000.
[10] J. Grabowski and G. Marmo, The graded Jacobi algebras and (co)homology. J. Phys. A 36(2003), no. 1, 161-181. doi:10.1088/0305-4470/36/1/311
[11] F. Guedira and A. Lichnerowicz, Géometrie des algèbres de Lie locales de Kirillov. J. Math. Pures Appl. 63(1984), 407-484.
[12] M. Gualtieri, Generalized complex geometry. arXiv:math/0703298v2 [math.DG].
[13] N. Hitchin, Generalized Calabi-Yau manifolds. Q. J. Math. 54(2003), no. 3, 281-308. doi:10.1093/qmath/hag025
[14] D. Iglesias-Ponte and A. Wade, Contact manifolds and generalized complex structures. J. Geom. Phys. 53(2005), no. 3, 249-258. doi:10.1016/j.geomphys.2004.06.006
[15] , Integration of Dirac-Jacobi structures. J. Phys. A 39(2006), no. 16, 4181-4190. doi:10.1088/0305-4470/39/16/006
[16] A. A. Kirillov, Local Lie algebras. (Russian) Uspehi Mat. Nauk 31(1976), no. 4(190), 57-76.
[17] U. Lindström, R. Minasian, A. Tomasiello, and M. Zabzine, Generalized complex manifolds and supersymmetry. Comm. Math. Phys. 257(2005), no. 1, 235-256. doi:10.1007/s00220-004-1265-6
[18] K. Mackenzie, General theory of Lie groupoids and Lie algebroids. London Mathematical Society Lecture Note Series, 213, Cambridge University Press, Cambridge, 2005.
[19] K. C. H. Mackenzie and P. Xu, Integration of Lie bialgebroids. Topology 39(2000), no. 3, 445-467. doi:10.1016/S0040-9383(98)00069-X
[20] C.-M. Marle, The Schouten-Nijenhuis bracket and interior products. J. Geom. Phys. 23(1997), no. 3-4, 350-359. doi:10.1016/S0393-0440(97)80009-5
[21] $\longrightarrow$, On Jacobi manifolds and Jacobi bundles. In: Symplectic geometry, groupoids, and integrable systems (Berkeley, CA, 1989), Math. Sci. Res. Inst. Publ., 20, Springer, New York, 1991, pp. 227-246.
[22] I. Moerdijk and J. Mrčun, On integrability of infinitesimal actions. Amer. J. Math. 124(2002), no. 3, 567-593. doi:10.1353/ajm.2002.0019
[23] J. M. Nunes da Costa and J. Clemente-Gallardo, Dirac structures for generalized Lie bialgebroids. J. Phys. A 37(2004), no. 7, 2671-2692. doi:10.1088/0305-4470/37/7/011
[24] P. Ševera and A. Weinstein, Poisson geometry with a 3-form background. Prog. Theor. Phys. Suppl. 144(2001), 145-154.
[25] K. Shiga, Cohomology of Lie algebras over a manifold I. J. Math Soc. Japan 26(1974), 324-61. doi:10.2969/jmsj/02620324
[26] I. Vaisman, Generalized CRF-structures. Geom. Dedicata 133(2008), 129-154. doi:10.1007/s10711-008-9239-z
[27] A. Wade, Conformal Dirac structures. Lett. Math. Phys. 53(2000), no. 4, 331-348. doi:10.1023/A:1007634407701
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[^1]:    ${ }^{1}$ See, for instance, $[8,19,22]$ for more details. Here $\mathbf{G L}(V)$ is the Lie groupoid of linear isomorphisms of the fibres of $V$, namely $\mathbf{G L}(V)_{x}^{y}=\operatorname{Hom}\left(V_{x}, V_{y}\right)$.
    ${ }^{2}$ Since $A$ is transitive and $M$ is connected, $t: \Gamma_{x} \rightarrow M$ is a surjective submersion. Let $y \in M$, and let $\sigma: U \rightarrow \Gamma_{x}$ be a section (so that $t \circ \sigma=\mathrm{id}$ ). Then $(z, v) \rightarrow \sigma(z)(v): U \times V_{x} \rightarrow V_{U}$ is a diffeomorphism.

[^2]:    ${ }^{3}$ An element of $\mathcal{G}$ is an equivalence class, which we may view as a subset of $\nu^{\prime-1}(y) \times \nu^{\prime-1}(z)$ that is $G$ invariant. As such, we may view it as the graph of an equivariant diffeomorphism $\nu^{\prime-1}(y) \rightarrow \nu^{\prime-1}(z)$. The multiplication in $\mathcal{G}$ is simply the composition of these diffeomorphisms. See [9] for details.

[^3]:    ${ }^{4}$ Here we use the identifications $P=\Gamma_{x}$ and $P^{\prime}=\mathcal{G}_{x}$. It is a covering map, since the right invariant vector fields, which are identified with the sections of $A$, span the tangent space of the source fibres.
    ${ }^{5}$ Here we use the identifications $G=\Gamma_{x}^{x}$ and $G^{\prime}=\mathcal{Y}_{x}^{x}$.

[^4]:    ${ }^{6}$ This follows from the fact that for any $x, y \in F$ there is a Lie algebroid morphism of $A$ covering a diffeomorphism of $M$ that takes $x$ to $y$. In addition these morphisms can be assumed to come from flowing along a section of $A$, and hence extend to $V$.

