# GHROMATIC SOLUTIONS, II 

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1. Introduction. This paper is a continuation of the Waterloo Research Report CORR 81-12, (see [1]) referred to in what follows as I. That Report is entitled "Chromatic Solutions". It is largely concerned with a power series $h$ in a variable $z^{2}$, in which the coefficients are polynomials in a "colour number" $\lambda$. By definition the coefficient of $z^{2 r}$, where $r>0$, is the sum of the chromatic polynomials of the rooted planar triangulations of $2 r$ faces. (Multiple joins are allowed in these triangulations.) Thus for a positive integral $\lambda$ the coefficient is the number of $\lambda$-coloured rooted planar triangulations of $2 r$ faces. The use of the symbol $z^{2}$ instead of a simple letter $t$ is for the sake of continuity with earlier papers.

In I we consider the case
(1) $\lambda=2+2 \cos (2 \pi / n)$,
where $n$ is an integer exceeding 4 . For each $n$ a set of parametric equations is exhibited. In principle, and sometimes in actuality, these permit the determination of the coefficients in $h$.

In the present paper we carry the theory further and obtain a differential equation for $h$. The proof applies to all values of $\lambda$ of the form (1), with $n$ at least 5 . Because the coefficients in $h$ are polynomials in $\lambda$ it can be inferred that the equation holds for all values of $\lambda$ except 4 . The value 4 is excluded because it makes some of the coefficients in the equation infinite. The paper concludes with a consideration of the limiting case, leading to a special differential equation for $h$ valid when $\lambda=4$.
2. The functions $\beta$ and $\gamma$. The theory of I is expressed in terms of a power series $\beta$ in $z^{2}$. There are equations in the Report permitting us to relate $\beta$ to $h$. From Equation I (50) with $r=2$ we have

$$
\begin{equation*}
\lambda p_{3}{ }^{\prime}=\beta p_{2}{ }^{\prime}-(4-\theta) p_{2} \tag{2}
\end{equation*}
$$

Here, and hereafter in this Report, a prime denotes differentiation with respect to $z^{2}$. The functions $p_{2}$ and $p_{3}$ are defined in terms of other functions $q_{2}$ and $q_{3}$ by $\mathrm{I}(43)$. In the even case, when $\theta=0$, these are given in terms of $z^{2}$ and $h$ by $\mathrm{I}(21)$ and $\mathrm{I}(22)$. In the odd case $(\theta=1)$ they are given by $\mathrm{I}(23)$ and $\mathrm{I}(24)$. Hence we can, in both cases, find $\beta$ in

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terms of $z^{2}$ and $h$. In both cases the result is

$$
\begin{equation*}
\beta=\nu\left\{\lambda^{-1} z^{4} h+\nu z^{2}\right\}^{\prime}+3 z^{2} \tag{3}
\end{equation*}
$$

(For $\nu$, see $\mathrm{I}(5)$ and $\mathrm{I}(6)$. ) We note that $\beta$ has the integral

$$
\begin{equation*}
\gamma=\lambda^{-1} \nu z^{4} h+\nu^{2} z^{2}+3 z^{4} / 2 \tag{4}
\end{equation*}
$$

In working towards our differential equations we shall use $\gamma$ rather than $h$.
3. A partial differential equation. Let us introduce a new independent indeterminate $w$ and write also $u=w^{2}$. We make the definition
(5) $U=\sum_{r=\theta}^{M} p_{r} w^{2 r-\theta}$,
so that $U$ is a function of two independent variables $w$ and $z^{2}$. We can now rewrite Equation $I(54)$ as
(6) $(\alpha+\beta u) \frac{\partial U}{\partial\left(z^{2}\right)}=2 u^{2} \frac{\partial U}{\partial u}$.

The function $\alpha$ is evaluated in I and found to be a constant. In fact

$$
\begin{equation*}
\alpha=-\lambda \tag{7}
\end{equation*}
$$

by $\mathrm{I}(60)$.
Given appropriate boundary conditions we might hope to use the partial differential equation (6) to determine $U$ as a power series in $z^{2}$ and $u$, in terms of the coefficients in $\beta$. Such boundary conditions are available. Let us write $F_{0}$ or $[F]_{0}$ to indicate the coefficient of $z^{0}$ in a power series $F$. Then

$$
\begin{aligned}
U_{0} & =\sum_{r=\theta}^{M}\left[p_{r}\right]_{0} w^{2 r-\theta} \\
& =\sum_{r=\theta}^{M} \frac{n(-1)^{\tau}(M+r-\theta-1)!\nu^{2 r-\theta} w^{2 r-\theta}}{2 \cdot(2 r-\theta)!(M-r)!}
\end{aligned}
$$

by $\mathrm{I}(43)$ and Theorem $\mathrm{I}(3.1)$,

$$
=\sum_{r=0}^{M} \frac{n(-1)^{r+\theta}(M-r-1)!(\nu w)^{2 r+\theta}}{2 \cdot(2 r+\theta)!(M-r-\theta)!}
$$

Hence

$$
\begin{equation*}
U_{0}=(-1)^{M} \cos \left\{n \cos ^{-1}\left(\frac{1}{2} \nu w\right)\right\} \tag{8}
\end{equation*}
$$

by I(11).
From (8) we can find the values at $z=0$ of the derivatives of $U$ of all orders with respect to $u=w^{2}$. Using (6) we can express the derivatives of $U$, of all orders with respect to $u$ and $z^{2}$, in terms of the derivatives
with respect to $u$. Hence we can determine the values at $z=0$ of all the derivatives, in terms of the coefficients in $\beta$, and therefore we can determine $U$ as a power series in $u$ and $z^{2}$. It is important to notice that this procedure determines the power series $U$ uniquely.
We can treat similarly a power series $W$ in $u$ and $z^{2}$ defined as a solution of the partial differential equation

$$
\begin{equation*}
(\alpha+\beta u) \frac{\partial W}{\partial\left(z^{2}\right)}=2 u^{2} \frac{\partial W}{\partial u} \tag{9}
\end{equation*}
$$

satisfying the boundary condition
(10) $\quad[W]_{0}=n \cos ^{-1}\left(\frac{1}{2} \nu w\right)$.

The inverse cosine is to be defined as a power series in $w$ with initial term $\pi / 2$. As with $U$ we find that $W$ is uniquely determined.

The two power series $W$ and $U$ are related as follows.

$$
\begin{equation*}
U=(-1)^{M} \cos W \tag{11}
\end{equation*}
$$

To prove this we have only to observe that the function $U$ determined by (11) satisfies both (6) and (8). In what follows we shall work with $W$ rather than $U$.
4. The polynomial $Y$. We now study Equation I(44). Parameters $\psi$ corresponding to even values of $M+m$ are roots of the equation
(12) $\sum_{r=\theta}^{M} p_{r} w^{2 r-\theta}=1$
in $w$. Moreover by I(45) they are repeated roots of this equation.
Similarly the parameters $\psi$ corresponding to odd values of $M+m$ are repeated roots of the equation

$$
\begin{equation*}
\sum_{r=\theta}^{M} p_{r} w^{2 \tau-\theta}=-1 . \tag{13}
\end{equation*}
$$

Combining these observations we see that each of our $M-2$ distinct parameters $\psi$ is a repeated root of

$$
\begin{equation*}
\left\{\sum_{r=\theta}^{M} p_{r} w^{2 r}\right\}^{2}-w^{2 \theta}=0 . \tag{14}
\end{equation*}
$$

By the definitions of I each $\psi$ has a positive initial term, so no one of them can be the negative of another. Since the left of (14) is a polynomial in $w^{2}$ we can now recognize the negatives of the parameters $\psi$ as additional repeated roots of (14). Accordingly the squares of the parameters $\psi$ are $M-2$ distinct repeated roots for $u$ of the equation

$$
\begin{equation*}
\left\{\sum_{r=\theta}^{M} p_{r} u^{r}\right\}^{2}-u^{\theta}=0 \tag{15}
\end{equation*}
$$

But it is pointed out in Section 5 of I that these squares are the roots of

$$
\sum_{r=0}^{M-2} p_{r+2}^{\prime} u^{r}=0 .
$$

We therefore have a polynomial identity

$$
\begin{equation*}
\left\{\sum_{r=\theta}^{M} p_{r} u^{r}\right\}^{2}-u^{\theta}=Y\left\{\sum_{r=0}^{M-2} p_{r+2}^{\prime} u^{r}\right\}^{2}, \tag{16}
\end{equation*}
$$

where $Y$ is a polynomial of degree 4 in $u$, whose coefficients are functions of $z^{2}$.

If $\theta=0$ then $p_{0}=1$, by $\mathrm{I}(46)$. Since $p_{2}{ }^{\prime}$ is non-zero, by $\mathrm{I}(58)$ and $I(59)$, we deduce that $Y$ always has $u$ as a factor. Hence we can write

$$
\begin{equation*}
Y=u\left(A+B u+C u^{2}+D u^{3}\right) \tag{17}
\end{equation*}
$$

where $A, B, C$ and $D$ are functions of $z^{2}$.
Writing (16) in terms of $U$ we find that
(18) $u^{4}\left(1-U^{2}\right)=-Y\left\{\frac{\partial U}{\partial\left(z^{2}\right)}\right\}^{2}$.

Hence, by (11),
(19) $\quad u^{4}=-Y\left\{\frac{\partial W}{\partial\left(z^{2}\right)}\right\}^{2}$.

We now try to eliminate $W$ from (19). First we write the equation in the form
(20) $\frac{\partial W}{\partial\left(z^{2}\right)}=u^{2}(-Y)^{-\frac{1}{2}}$.

Hence, by (9),
(21) $\frac{\partial W}{\partial u}=\frac{1}{2}(\alpha+\beta u)(-Y)^{-\frac{1}{2}}$.

Accordingly
(22) $\frac{\partial}{\partial u}\left\{u^{2}(-Y)^{-\frac{1}{2}}\right\}=\frac{\partial}{\partial\left(z^{2}\right)}\left\{\frac{1}{2}(\alpha+\beta u)(-Y)^{-\frac{1}{2}}\right\}$,
from which we can deduce

$$
\begin{equation*}
\left(8-2 \beta^{\prime}\right) u Y-2 u^{2} \frac{\partial Y}{\partial u}+(\alpha+\beta u) \frac{\partial Y}{\partial\left(z^{2}\right)}=0 \tag{23}
\end{equation*}
$$

Using (17) we can write this as

$$
\begin{aligned}
& \left(8-2 \beta^{\prime}\right)\left(A u^{2}+B u^{3}+C u^{4}+D u^{5}\right) \\
& \quad-2\left(A u^{2}+2 B u^{3}+3 C u^{4}+4 D u^{5}\right) \\
& \quad+(\alpha+\beta u)\left(A^{\prime} u+B^{\prime} u^{2}+C^{\prime} u^{3}+D^{\prime} u^{4}\right)=0
\end{aligned}
$$

By equating coefficients of like powers of $u$ in this equation we obtain the following differential equations.
(24) $\alpha A^{\prime}=0$,
(25) $\left(6-2 \beta^{\prime}\right) A+\alpha B^{\prime}+\beta A^{\prime}=0$,
(26) $\left(4-2 \beta^{\prime}\right) B+\alpha C^{\prime}+\beta B^{\prime}=0$,
(27) $\left(2-2 \beta^{\prime}\right) C+\alpha D^{\prime}+\beta C^{\prime}=0$,
(28) $-2 \beta^{\prime} D+\beta D^{\prime}=0$.

In order to solve these equations we need to know $A_{0}, B_{6}, C_{0}$ and $D_{0}$, the coefficients of $z^{0}$ in $A, B, C$ and $D$ respectively. But from (9) and (19) we have

$$
u^{4}\left(\alpha+\beta_{0} u\right)^{2}=-[Y]_{0} \cdot 4 u^{4}\left\{\frac{\partial W_{0}}{\partial u}\right\}^{2} .
$$

We have $\beta_{0}=\nu^{2}$, by (3), and of course $\alpha=-\lambda$. We can now deduce from (10) that

$$
\begin{align*}
& \left(\lambda-\nu^{2} u\right)^{2}=-n^{2} \nu^{2}[Y]_{0} u^{-1}\left(4-\nu^{2} u\right)^{-1}  \tag{29}\\
& n^{2} \nu^{2}[Y]_{0}=-u\left(4-\nu^{2} u\right)\left(\lambda-\nu^{2} u\right)^{2} .
\end{align*}
$$

As consequences of (29) we have

$$
\begin{align*}
& n^{2} \nu^{2} A_{0}=-4 \lambda^{2}, n^{2} B_{0}=8 \lambda+\lambda^{2} \\
& n^{2} C_{0}=-(4+2 \lambda) \nu^{2}, n^{2} D_{0}=\nu^{4} \tag{30}
\end{align*}
$$

5. A differential equation for $\gamma$. From (24) we see that $A$ is constant. Hence, by (30),
(31) $n^{2} \nu^{2} A=-4 \lambda^{2}$.

We can now integrate (25) as

$$
6 A z^{2}-2 A \beta+\alpha B+c_{1}=0
$$

where $c_{1}$ is a constant. Taking coefficients of $z^{0}$ in this and applying (30) we find that $n^{2} c_{1}=\lambda^{3}$. We can deduce that
(32) $n^{2} \nu^{2} B=\lambda\left\{\lambda \nu^{2}-24 z^{2}+8 \beta\right\}$,
(33) $n^{2} \nu^{2} B^{\prime}=\lambda\left\{-24+8 \beta^{\prime}\right\}$.

Substituting from these equations in (26) we find

$$
\begin{align*}
& n^{2} \nu^{2} \lambda C^{\prime}=\left(4-2 \beta^{\prime}\right) \lambda\left(\lambda \nu^{2}-24 z^{2}+8 \beta\right)+\beta \lambda\left(-24+8 \beta^{\prime}\right)  \tag{34}\\
& n^{2} \nu^{2} C^{\prime}=4 \lambda \nu^{2}-96 z^{2}+8 \beta+48 z^{2} \beta^{\prime}-2 \lambda \nu^{2} \beta^{\prime}-8 \beta \beta^{\prime}
\end{align*}
$$

Integrating this we obtain

$$
n^{2} \nu^{2} C=4 \lambda \nu^{2} z^{2}-48 z^{4}+8 \gamma+48 z^{2} \beta-48 \gamma-2 \lambda \nu^{2} \beta-4 \beta^{2}+c_{2},
$$

where $c_{2}$ is a constant. We now take coefficients of $z^{0}$, noting that $\gamma_{0}$ is zero by (4). By (30)

$$
-(4+2 \lambda) \nu^{1}=-2 \lambda \nu^{4}-4 \nu^{4}+c_{2} .
$$

Hence $c_{2}=0$. We can write

$$
\begin{equation*}
n^{2} \nu^{2} C=4 \lambda \nu^{2} z^{2}-48 z^{4}-40 \gamma+48 z^{2} \beta-2 \lambda \nu^{2} \beta-4 \beta^{2} \tag{35}
\end{equation*}
$$

Solving (28) we have

$$
D=c_{3} \beta^{2}
$$

for some constant $c_{3}$. Using (30) we find that $n^{2} c_{3}=1$. Hence
(36) $n^{2} D=\beta^{2}$,
(37) $n^{2} D^{\prime}=2 \beta \beta^{\prime}$.

So far we have not used (27). Let us multiply that equation by $n^{2} \nu^{2}$ and substitute in it from (34), (35) and (37). We find that

$$
\begin{array}{r}
\left(2-2 \beta^{\prime}\right)\left(4 \lambda \nu^{2} z^{2}-48 z^{4}-40 \gamma-2 \lambda \nu^{2} \beta+48 z^{2} \beta-4 \beta^{2}\right)-2 \lambda \nu^{2} \beta \beta^{\prime} \\
+\beta\left(4 \lambda \nu^{2}-96 z^{2}+8 \beta+48 z^{2} \beta^{\prime}-2 \lambda \nu^{2} \beta^{\prime}-8 \beta \beta^{\prime}\right)=0 .
\end{array}
$$

This simplifies to

$$
\left(1-\beta^{\prime}\right)\left(\lambda \nu^{2} z^{2}-12 z^{4}-10 \gamma+6 z^{2} \beta\right)=6 z^{2} \beta
$$

that is

$$
\begin{equation*}
\left(1-\gamma^{\prime \prime}\right)\left(-\lambda \nu^{2} z^{2}+12 z^{4}+10 \gamma-6 z^{2} \gamma^{\prime}\right)+6 z^{2} \gamma^{\prime}=0 \tag{38}
\end{equation*}
$$

In view of (4) this differential equation for $\gamma$ can be regarded as the promised differential equation for $h$. It can be written also as

$$
\begin{equation*}
\left(1-\gamma^{\prime \prime}\right)\left(-\lambda \nu^{2} z^{2}+12 z^{4}+10 \gamma\right)+6 z^{2} \gamma^{\prime} \gamma^{\prime \prime}=0 \tag{39}
\end{equation*}
$$

It should be emphasized that our proof of Equations (38) and (39) applies only when $\lambda$ is of the form (1), with $n>4$. But the proof can be extended to other values of $\lambda$, and this is done in the following section.
6. The range of validity of the differential equation. Let us discuss Equation (38) without reference to any previous definition of $\gamma$. We do this for an arbitrary real or complex value of $\lambda$, excluding only the value $\lambda=4$ which makes $\nu$ infinite. We consider the possibility of the equation having a solution in the form of a power series

$$
\begin{equation*}
\gamma=\sum_{r=0}^{\infty} \gamma_{r} z^{2 r} \tag{40}
\end{equation*}
$$

where the $\gamma_{r}$ depend only on $\lambda$. For any such solution we have also
(41) $\quad \gamma^{\prime}=\sum_{r=1}^{\infty} r \gamma_{r} z^{2(r-1)}$
(42) $\quad \gamma^{\prime}=\sum_{r=2}^{\infty} r(r-1) \gamma_{r} z^{2(r-2)}$.

If this $\gamma$ is to be identical with the $\gamma$ of Equation (4) we must have

$$
\begin{equation*}
\gamma_{0}=0, \gamma_{1}=\nu^{2} \tag{43}
\end{equation*}
$$

Accordingly we impose (43) as an extra condition on $\gamma$.
6.1. There is one and only one power series $\gamma$ of the form (40) which satisfies both (38) and (43). Moreover the coefficients in $\gamma$ have the following property: each $\gamma_{j}$ with $j>1$ is of the form $\nu P_{j}$, where $P_{j}$ is a polynomial in $\nu^{-1}$ whose coefficients are rational numbers.

Proof. Let us begin by assuming a power series (40) to satisfy (38) and (43). Let us study the relations that must hold between its coefficients.

Equating coefficients of $z^{0}$ in (38) we find

$$
\left(1-2 \gamma_{2}\right)\left(10 \gamma_{0}\right)=0
$$

which is consistent with (43). For coefficients of $z^{2}$ we have

$$
\left(1-2 \gamma_{2}\right)\left(-\lambda \nu^{2}+10 \gamma_{1}-6 \gamma_{1}\right)+6 \gamma_{1}=0
$$

If we put $\gamma_{1}=\nu^{2}$ as required by (43) we have

$$
-\lambda \nu^{2}+10 \gamma_{1}-6 \gamma_{1}=(4-\lambda) \nu^{2}=\nu
$$

It follows that

$$
\begin{equation*}
2 \gamma_{2}=6 \nu+1=\nu\left(6+\nu^{-1}\right) \tag{44}
\end{equation*}
$$

Accordingly we say that $P_{2}$ is $3+\frac{1}{2} \nu^{-1}$.
We next equate coefficients of $z^{4}$.

$$
\left(1-2 \gamma_{2}\right)\left(12+10 \gamma_{2}-12 \gamma_{2}\right)+\left(-6 \gamma_{3}\right)(\nu)+12 \gamma_{2}=0
$$

Substituting from (44) we deduce that

$$
\begin{equation*}
\gamma_{3}=\nu\left(6-5 \nu^{-1}+\nu^{-2}\right) \tag{45}
\end{equation*}
$$

We say therefore that $P_{3}$ is $6-5 \nu^{-1}+\nu^{-2}$.
For an integer $m$ exceeding 2 we can equate coefficients of $z^{2(m-1)}$ and have

$$
\begin{align*}
& \left\{-(m+1) m \gamma_{m+1}\right\}\{\nu\}+\left\{-m(m-1) \gamma_{m}\right\}\left\{12-2 \gamma_{2}\right\}  \tag{46}\\
& +\sum_{r=2}^{m-2}\left\{-(m+1-r)(m-r) \gamma_{m+1-r}\right\}\left\{10 \gamma_{r+1}-6(r+1) \gamma_{r+1}\right\} \\
& \\
& +\left\{1-2 \gamma_{2}\right\}\left\{10 \gamma_{m}-6 m \gamma_{m}\right\}+6 m \gamma_{m}=0 .
\end{align*}
$$

Starting with the known values of $\gamma_{0}, \gamma_{1}, \gamma_{2}$ and $\gamma_{3}$ we can use this equation to determine $\gamma_{4}, \gamma_{5}$ and so on. The coefficients $\gamma_{j}$ are thus uniquely determined by (46), in a recursive manner. We infer that there is at most one power series $\gamma$ satisfying both (38) and (43). On the other hand the coefficients $\gamma_{j}$ determined by (43), (44), (45) and (46) do
specify a formal power series $\gamma$ that satisfies (38). For when this power series is substituted on the left of (38) the coefficients of all powers of $z^{2}$ vanish.

To complete the proof we have to show that when $j>1$ the coefficient $\gamma_{j}$ is of the required form $\nu P_{j}$. We have already verified this in the cases $j=2$ and $j=3$. Assume as an inductive hypothesis that it is true up to $j=m$, where $m$ is some integer exceeding 2 , and consider the case $j=m+1$. Then by (46) and the inductive hypothesis

$$
\nu \gamma_{m+1}=\nu^{2} P_{m+1}
$$

where $P_{m+1}$ is a polynomial in $\nu^{-1}$ with rational numbers as coefficients. Hence $\gamma_{m+1}=\nu P_{m+1}$ and the induction succeeds.

We have $\nu^{-1}=4-\lambda$. Hence we can assert the following variant of 6.1.
6.2. There is one and only one power series $\gamma$ of the form (40) which satisfies both (38) and (43). Moreover the coefficients in $\gamma$ have the following property: each $\gamma_{j}$ with $j>1$ is of the form $\nu P_{j}$, where $P_{j}$ is a polynomial in $\lambda$ whose coefficients are rational numbers.

Consider the power series $\gamma$ defined in terms of $h$ by Equation (4). Like $h$ it is defined for all values of $\lambda$ other than 4 . ( $h$ is defined for the value 4 also.) The coefficients in this $\gamma$ have the following property: If $j>1$ the coefficient of $z^{2 j}$ is of the form $\nu R_{j}$, where $R_{j}$ is a polynomial in $\lambda$ whose coefficients are rational numbers. Here we use the fact that each chromatic polynomial appearing in $h$ divides by $\lambda$. (The rational numbers mentioned are integers if $j>2$.) For each such $j$ the effect of the foregoing theory is to show that $P_{j}=R_{j}$ for infinitely many values of $\lambda$, of the form (1). Since $P_{j}$ and $R_{j}$ are polynomials in $\lambda$ it follows that they are identical. We conclude that the power series $\gamma$ defined in terms of $h$ by Equation (4) is identical with the power series $\gamma$ exhibited in the proof of 6.1 , for all values of $\lambda$ other than 4 . Let us state this result as a theorem.
6.3. The power series $\gamma$ defined in terms of $h$ by Equation (4) satisfies the differential equations (38) and (39) for all values of $\lambda$ other than 4.
7. The case $\lambda=4$. Let us write

$$
\begin{equation*}
H=z^{4} h \tag{47}
\end{equation*}
$$

Then

$$
\begin{align*}
& \gamma=\lambda^{-1} \nu H+\nu^{2} z^{2}+3 z^{4} / 2,  \tag{48}\\
& \gamma^{\prime}=\lambda^{-1} \nu H^{\prime}+\nu^{2}+3 z^{2}, \\
& \gamma^{\prime \prime}=\lambda^{-1} \nu H^{\prime \prime}+3 .
\end{align*}
$$

We can substitute from these equations in (38), thereby obtaining a differential equation for the power series $H$. The author found that

$$
\begin{align*}
\lambda^{-1} H^{\prime \prime}\left\{z^{2}+9 \nu^{-1} z^{4}\right. & \left.+10 \lambda^{-1} H-6 \lambda^{-1} z^{2} H^{\prime}\right\}  \tag{51}\\
& =-2 \nu^{-1} z^{2}+6 z^{2}-20 \lambda^{-1} \nu^{-1} H+18 \lambda^{-1} \nu^{-1} z^{2} H^{\prime} .
\end{align*}
$$

Considering the limiting case $\lambda \rightarrow 4, \nu^{-1} \rightarrow 0$, we arrive at the following theorem.
7.1. In the case $\lambda=4$ the power series $H$ satisfies the differential equation (52) $H^{\prime \prime}\left\{2 z^{2}+5 H-3 z^{2} H^{\prime}\right\}=48 z^{2}$.

## Reference

1. W. T. Tutte, Chromatic solutions, Can. J. Math. 34 (1982), 741-758.

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