

A REMARK ON BASES IN HARDY SPACES

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ABSTRACT. The Franklin spline system in $[0, 1]$ has been generalized by Strömberg to a system in \mathbb{R}^n which is an unconditional basis in $H^p(\mathbb{R}^n)$ for $p > n/(n + m + 1)$. Here the natural number m is the order of the system. For some of these values of p , it was known that the H^p quasi-norm is equivalent to a certain expression containing the coefficients of the function with respect to this basis. We prove this equivalence for all $p > n/(n + m + 1)$.

1. **Introduction.** In [4] J.-O. Strömberg constructs an orthonormal basis $(f_\nu^\omega)_{\nu, \omega}$ in $L^2(\mathbb{R}^n)$ which generalizes the Franklin system in $L^2[0, 1]$. Here $\nu = (j, k) = (j, k_1, \dots, k_n)$ ranges over \mathbb{Z}^{n+1} and ω over a finite set. Each f_ν^ω is a tensor product of one-dimensional spline functions of order m and takes its largest values near the cube

$$Q_\nu = \{x : 2^{-i}k_i \leq x_i \leq 2^{-i}(k_i + 1), \quad i = 1, \dots, n\},$$

whose characteristic function is denoted by χ_ν .

Strömberg proves that this system is an unconditional basis in the Hardy space $H^p(\mathbb{R}^n)$, $p > n/(n + m + 1)$, and compares the quasinorm of f in H^p to its coefficients $c_\nu^\omega = (f, f_\nu^\omega)$. For all $p > n/(n + m + 1)$ he shows that

$$(1.1) \quad \left\| \left(\sum_{\nu, \omega} |c_\nu^\omega|^2 2^{nj} \chi_\nu \right)^{1/2} \right\|_{L^p} \leq C \|f\|_{H^p}.$$

Here and in the sequel, C denotes various constants. At least in one dimension, (1.1) holds only for these values of p , see [3]. The converse inequality

$$(1.2) \quad \|f\|_{H^p} \leq C \left\| \left(\sum_{\nu, \omega} |c_\nu^\omega|^2 2^{nj} \chi_\nu \right)^{1/2} \right\|_{L^p}$$

is proved only for $p > n/(n/2 + m + 2)$. We shall prove (1.2) for all $p > n/(n + m + 1)$, by modifying Strömberg's proof. This answers a question asked at the end of [4]. It is enough to prove the following.

THEOREM. *Let $n/(n + m + 1) < p \leq 1$. There is a constant $C = C(n, m, p)$ such that for any finite set of numbers c_ν^ω*

$$(1.3) \quad \left\| \sum c_\nu^\omega f_\nu^\omega \right\|_{H^p} \leq C \left\| \left(\sum_{\nu, \omega} |c_\nu^\omega|^2 2^{nj} \chi_\nu \right)^{1/2} \right\|_{L^p}.$$

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We shall need some properties of f_ν^ω . Strömberg's definition reads

$$(1.4) \quad f_\nu^\omega(x) = 2^{ni/2} \tau^\omega(2^j x - k)$$

with $\nu = (j, k) \in \mathbb{Z} \times \mathbb{Z}^n$, where τ^ω is a tensor product

$$\tau^\omega(x) = \prod_{i=1}^n \omega_i(x_i).$$

Here each ω_i is either of two fixed spline functions τ and ρ in \mathbb{R} , and at least one is τ . These two functions are in $C^m(\mathbb{R})$ and they are polynomials in the complementary intervals of a countable discrete set. Moreover,

$$(1.5) \quad |D^k \tau(t)| \leq Cr^{|k|}, \quad k = 0, \dots, m + 1,$$

and similarly for ρ , for some $r < 1$. The moments $\int \tau(t) t^\alpha dt$ vanish for $\alpha = 0, \dots, m + 1$.

2. Proof of the theorem. Choose a radial nonzero $\psi \in C^\infty(\mathbb{R}^n)$ with support in $|x| \leq 1$ and vanishing moments up to order m . Writing $\psi_s(x) = s^{-n} \psi(x/s)$ and $G(x, s) = \psi_s * f(x)$ for $f \in \mathcal{S}'$, we define a Lusin function

$$A_\gamma G(z) = \left(\iint_{|x-z| < \gamma s} |G(x, s)|^2 \frac{dx ds}{s^{n+1}} \right)^{1/2}.$$

Here $z, x \in \mathbb{R}^n$ and $s > 0$. Then

$$(2.1) \quad \|A_\gamma G\|_{L^p} \sim \|f\|_{H^p},$$

in the sense that the quotient between these two quantities stays away from 0 and ∞ , see [1, Theorems 6.6 and 6.9].

Let a be an atom in H^p with vanishing moments up to order m . The arguments on p. 492 of [4] show that $\sum \pm(a, f_\nu^\omega) f_\nu^\omega$ belongs to H^p with quasi-norm at most C , for any sign combination. Here molecules could also be used. Hence, the operator which maps $f \in H^p$ onto $\sum \pm(f, f_\nu^\omega) f_\nu^\omega$ is bounded in H^p , uniformly over all sign choices, and the H^p quasi-norms of all these sums are comparable.

We consider the probability measure on the set of all sign choices making the signs independent and each sign of probability 1/2. Let E denote the corresponding expectation. We conclude from (2.1) with $\gamma = 1$ that

$$\begin{aligned} \left\| \sum c_\nu^\omega f_\nu^\omega \right\|_{H^p}^p &\sim E \left\| \sum \pm c_\nu^\omega f_\nu^\omega \right\|_{H^p}^p \sim E \int dz \left(\iint_{|x-z| < s} \left| \psi_s * \sum \pm c_\nu^\omega f_\nu^\omega(x) \right|^2 \frac{dx ds}{s^{n+1}} \right)^{p/2} \\ &\leq \int dz \left(\iint_{|x-z| < s} E \left| \sum \pm c_\nu^\omega \psi_s * f_\nu^\omega(x) \right|^2 \frac{dx ds}{s^{n+1}} \right)^{p/2}, \end{aligned}$$

by Hölder’s inequality. But $E |\sum \pm a_j|^2 = \sum |a_j|^2$, so

$$(2.2) \quad \left\| \sum c_\nu^\omega f_\nu^\omega \right\|_{H^p}^p \leq C \int dz \left(\iint_{|x-z|<s} \sum |c_\nu^\omega|^2 |\psi_s * f_\nu^\omega(x)|^2 \frac{dx ds}{s^{n+1}} \right)^{p/2}.$$

We must thus estimate $\psi_s * f_\nu^\omega$. Our inequalities are similar to those of [4, p. 488–9].

LEMMA. For each $N > 0$ there exists $C = C(N)$ such that for $x \in \mathbb{R}^n$

$$\begin{aligned} |\psi_s * f_\nu^\omega(x)| &\leq C 2^{nj/2} (2^j s)^{m+1} (1 + 2^j |x - 2^{-j}k|)^{-N}, \quad s \leq 2^{-j} \\ &\leq C 2^{nj/2} (2^j s)^{-n-m-2} (1 + s^{-1} |x - 2^{-j}k|)^{-N}, \quad s > 2^{-j}. \end{aligned}$$

Proof. We let $t > 0$ and estimate $\psi_s * \tau_t^\omega$. Denote by P_x the Taylor polynomial of τ^ω at x of degree m . For $s \leq t$ we use the vanishing moments of ψ to get

$$|\psi_s * \tau_t^\omega(x)| = \left| \int_{|y| \leq s} \psi_s(y) t^{-n} \left(\tau^\omega \left(\frac{x-y}{t} \right) - P_{x/t} \left(-\frac{y}{t} \right) \right) dy \right|.$$

Because of (1.5), the parenthesis here is dominated by $C(|y|/t)^{m+1}(1 + t^{-1}|x|)^{-N}$. Thus,

$$|\psi_s * \tau_t^\omega(x)| \leq C t^{-n} \left(\frac{s}{t} \right)^{m+1} \left(1 + \frac{|x|}{t} \right)^{-N}.$$

For $s > t$, we use instead the Taylor polynomial Q_x of ψ at x of degree $m + 1$. Since $\tau^\omega(x)$ contains at least one factor $\omega_i(x_i) = \tau(x_i)$, its moments of order up to $m + 1$ vanish. Hence,

$$\psi_s * \tau_t(x) = s^{-n} \int \left(\psi \left(\frac{x-y}{s} \right) - Q_{x/s} \left(-\frac{y}{s} \right) \right) \tau_t^\omega(y) dy.$$

If $|x| \leq 2s$, this is easily estimated by $Cs^{-n}(t/s)^{m+2}$. If $|x| > 2s$, then $Q_{x/s} = 0$ and we may assume $|y| > |x|/2$ since otherwise $|x - y| > s$ and the integrand vanishes. Then (1.5) implies

$$|\psi_s * \tau_t^\omega(x)| \leq Cs^{-n} \int_{|y|>|x|/2} |\tau_t^\omega(y)| dy \leq Cs^{-n} \left(\frac{|x|}{t} \right)^{-N}.$$

So for $s > t$ and $N > m + 2$,

$$|\psi_s * \tau_t^\omega(x)| \leq Cs^{-n} \left(\frac{t}{s} \right)^{m+2} \left(1 + \frac{|x|}{s} \right)^{-N}.$$

The estimates obtained imply the lemma, as seen from (1.4).

Because of this lemma, (2.2) implies

$$\left\| \sum c_\nu^\omega f_\nu^\omega \right\|_{H^p}^p \leq C \int dz \left(\sum |c_\nu^\omega|^2 2^{nj} (I_1 + I_2) \right)^{p/2},$$

where

$$I_1 = \iint_{|x-z| < s \leq 2^{-i}} (2^i s)^{2(m+1)} \left(1 + \frac{|x-2^{-i}k|}{2^{-i}}\right)^{-2N} \frac{dx ds}{s^{n+1}}$$

and

$$I_2 = \iint_{\substack{|x-z| < s \\ s > 2^{-i}}} (2^i s)^{-2(n+m+2)} \left(1 + \frac{|x-2^{-i}k|}{s}\right)^{-2N} \frac{dx ds}{s^{n+1}}.$$

It is easily seen that

$$I_1 \leq C \int_0^{2^{-i}} (2^i s)^{2(m+1)} \frac{ds}{s} \left(1 + \frac{|z-2^{-i}k|}{2^{-i}}\right)^{-2N} \leq C \left(1 + \frac{|z-2^{-i}k|}{2^{-i}}\right)^{-2N}$$

and

$$I_2 \leq C \int_{2^{-i}}^\infty (2^i s)^{-2(n+m+2)} \left(1 + \frac{|z-2^{-i}k|}{s}\right)^{-2N} \frac{ds}{s} \leq C \left(1 + \frac{|z-2^{-i}k|}{2^{-i}}\right)^{-2(n+m+2)},$$

if N is large. Thus,

$$(2.3) \quad \left\| \sum c_\nu^\omega f_\nu^\omega \right\|_{H^p}^p \leq C \int dz \left(\sum |c_\nu^\omega|^2 2^{nj} \left(1 + \frac{|z-2^{-i}k|}{2^{-i}}\right)^{-2(n+m+2)} \right)^{p/2}.$$

Now define a function F in \mathbb{R}^{n+1} as in [2, p. 118], setting $F = c_\nu^\omega 2^{nj/2}$ on $\frac{1}{2}Q_\nu \times [2^{-i}, 2^{-i+1}]$ for each ν and $F = 0$ elsewhere. Here $\frac{1}{2}$ of course means concentric scaling of the cube in \mathbb{R}^{n+1} . Then (2.3) implies

$$\left\| \sum c_\nu^\omega f_\nu^\omega \right\|_{H^p} \leq C \|g_\lambda^*(F)\|_{L^p}, \quad \lambda = n + m + 2,$$

where

$$g_\lambda^*(F)(z) = \left(\iint \left(1 + \frac{|x-z|}{s}\right)^{-2N} |F(x, s)|^2 \frac{dx ds}{s^{n+1}} \right)^{1/2}$$

is a Littlewood-Paley function. Further,

$$\|g_\lambda^*(F)\|_{L^p} \leq C \|A_\gamma F\|_{L^p}, \quad \gamma > 0,$$

by [1, Theorem 3.5], since $\lambda > n/p$ here. But $A_\gamma F$ is dominated by $C(\sum |c_\nu^\omega|^2 2^{nj} \chi_\nu)^{1/2}$ if γ is small enough, and (1.3) and the theorem follow.

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