

THE NEUTRON STAR POPULATION IN THE GALAXY

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ABSTRACT. Properties of neutron stars are discussed and their density in Galaxy is roughly estimated. Their input in average galactic density is not large and they cannot be a reasonable dark matter candidates.

1. Introduction

Neutron stars (NS) are observed as radiopulsars, strong X-ray sources in close binaries: X-ray pulsars, X-ray bursters, X-ray sources with quasi-periodical oscillations (QPO); most transient X-ray sources also contain NS.

As the lifetime of radiopulsars and most of the strong X-ray sources is much less than the age of Galaxy, overwhelming majority of NS will be neither radiopulsars nor strong X-ray sources. Gamma-ray bursters (GRB) have been proposed as a candidates for old nearby NS (Bisnovatyi-Kogan et.al.,1975). Discovery of hard transient X-ray pulsar in the strongest GRB of 05 March 1979 (Mazetz et.al.,1979) and results of spatial distribution of faint GRB (Mazetz et.al.,1981) give evidences in favor of this interpretation. Here we estimate spatial density of NS on the base of theoretical analysis and observational data.

2. Stellar evolution and supernovae

Estimations of the number of NS in Galaxy from theory of stellar evolution with birth rate function (theoretically) and from the pulsar and supernovae statistics (observationally) were presented by Shapiro and Teukolsky (1983). Assume Salpeter birth rate function

$$\psi, dm = 2 \times 10^{-12} m^{-2.35} dm \text{ stars } pc^{-3} yr^{-1}, \quad m = M/M_{\odot}. \quad (1)$$

This function is valid for mass interval $0.4 \leq m \leq 10$ and is steeper for larger masses. The birthrate of stars with masses in the interval $m_1 \leq m \leq m_2$ follows from (1) after integration

$$\Psi(m_1, m_2) = \int_{m_1}^{m_2} \psi dm = 1.48 \times 10^{-12} (m_1^{-1.35} - m_2^{-1.35}) \text{ stars } pc^{-3} yr^{-1}. \quad (2)$$

The average mass of new born stars in mass interval $m_1 \leq m \leq m_2$ for Salpeter function is

$$\overline{m(m_1, m_2)} = 3.86 \frac{m_1^{-0.35} - m_2^{-0.35}}{m_1^{-1.35} - m_2^{-1.35}}. \quad (3)$$

Shapiro and Teukolsky (1983) obtained $\rho_{ns}/\rho_T = 0.02$ (ρ_{ns} is average density of the NS, $\rho_T = 0.14 M_{\odot} pc^{-3}$ is total barionic mass density in Galaxy) adopting that NS were born by stars with masses on main sequence (MMS) $4 \leq MMS \leq 10 M_{\odot}$ and NS mass is equal to $1.4 M_{\odot}$. Now the lower MMS for NS birth is taken equal to about $8 M_{\odot}$ (see i.e. Bisnovatyi-Kogan, 1989). The upper limit for MMS is not known. Taking into account progenitor mass of SN 1987A ($\sim 20 M_{\odot}$) and lower metallicity of stars in LMC we adopt for estimations the value $M_{lim} = 35 M_{\odot}$, for larger MMS collapse lead to formation black hole (BH). Masses of BH - stellar remnants, Cyg X-1 and A0620-00, are estimated between 4 and $15 M_{\odot}$. For $m_1=8$ and $m_2=35$ we obtain with account of (2)

$$\rho_{ns}/\rho_T = 9.55 \times 10^{-3}. \quad (4)$$

For BH density these authors get $\rho_{bh}/\rho_T=0.22$, using $m_{lim}=10$ and $M_{bh}=MMS$ with $\overline{m_{bh}}=38.6$, following from (3) at $m_2 = \infty$. Here we get $\rho_{bh}/\rho_T = 0.011$ for $m_{lim}=35$ and $\overline{m_{bh}} = 10$. Taking into account uncertainties, connected with stellar evolution in binaries, we obtain estimations

$$\rho_{ns}/\rho_T = 0.01 \div 0.03, \quad \rho_{bh}/\rho_T = 0.01 \div 0.03. \quad (5)$$

Corresponding birthrates in Galaxy will be

$$\dot{n}_{ns} = 0.01 \div 0.03 \text{ yr}^{-1}, \quad \dot{n}_{bh} = (1.5 \div 5) \times 10^{-3} \text{ yr}^{-1} \quad (6)$$

instead of values 0.021 and 0.0085 by Shapiro and Teukolsky (1983).

SN frequency in Galaxy is estimated as $R_{sn} = 1/28 \div 1/60 \text{ yr}^{-1}$ by different authors (Shapiro and Teukolsky, 1983) in accordance with theoretical estimations (6). In some SN explosions star is totally disrupted (presumably in SN I), so lower limit in (6) better coincide with SN statistics. It indicates that fate of massive stars in close binaries does not differ much from single stars.

3. Pulsar birthrate

Obtaining of pulsar birthrate from observations require a knowledge of luminosity evolution, form of the pulsar beam, magnetic field decay function. For exponential luminosity decay with time constant $4 \times 10^6 \text{ yr}$, the minimum pulsar birthrate was found to be one every 230 years in Galaxy. Observational data show that beams are elongated in the latitude direction and beaming factor may be about 2. With account of uncertainties the birthrate of pulsars in Galaxy is estimated as one over $30 \div 120 \text{ years}$ (Taylor and Stinebring, 1986) in excellent agreement with theoretical estimation (6).

4. X-ray sources and neutron star formation in binary systems

X-ray searches from the satellites lead to discovery of more than 100 bright sources, identified as NS in close binary systems. Total number of accreting neutron stars in Galaxy is of the same order, so their number is much less than total number of radiopulsars (about 10^5). Reminding that almost half of stars on the main sequence belong to binaries, we come to conclusion, that disruption of binary during formation of NS occurs frequently and the number of NS remaining in pairs is about 10^{-3} of all NS in Galaxy.

Large percent of disrupted pairs and existence of pulsar velocities much larger than orbital in pairs show that Blaauw effect cannot explain this phenomena and indicate to their recoil origin. Bisnovatyi-Kogan and Moiseenko (1991) proposed mechanism of violation of mirror symmetry during formation of rotating NS with toroidal and poloidal components of magnetic field in progenitor. Consider dipole field in combination with symmetrical toroidal, having the same sign in both hemispheres. Nonuniform contraction during collapse lead to differential rotation and generation of additional toroidal field from poloidal one. Newly generated poloidal field have different signs in two hemispheres and the sum of original and generated toroidal fields have no mirror symmetry. Amplification of toroidal field by twisting lead to magnetorotational explosion (Bisnovatyi-Kogan, 1989),

whose mirror asymmetry is determined by initial value of toroidal field. In order to get pulsar velocity ~ 300 km/s symmetrical part of toroidal field must be about 5×10^{14} Gs (Bisnovatyi-Kogan and Moiseenko, 1991). This value exceed the dipole pulsar magnetic fields ~ 500 times what is like the ratio of spot and large scale magnetic fields in Sun.

5. Recycled pulsars and LMXB

About 1/3 of bright X-ray sources belong to low-mass X-ray-binaries (LMXB), consisting of NS and low-mass star with $M \lesssim M_{\odot}$. These objects are rather old (about 10^8 years) and have small magnetic fields. Accretion disk reaches a surface of NS and accelerates its rotation up to millisecond periods. After ceasing of mass transfer NS transforms into radiopulsar, which has a name of "recycled radiopulsar" (RR). The transformation of X-ray sources into radiopulsars was first considered by Bisnovatyi-Kogan and Komberg (1974).

Most numerous sample of RR originates from LMXB. About half of LMXB are situated in globular clusters leading to suggestion of the origin of LMXB in tidal capture of a red dwarf by NS (Fabian, Pringle and Rees, 1975). Among 42 known RR only 11 lay outside globular clusters. This is connected mainly with selection effect inherent to the search procedure. RR contain about 10 % of all known radiopulsars so their input in average galactic density is small. Radiopulsars are related to RR by one of three main properties:

1. belonging to globular cluster
2. belonging to close binary system
3. very small (less than 10 msec) period.

Usually RR obtain 2 or 3 of these properties. In addition all RR have magnetic fields 1 - 3 order of magnitude less than the smallest magnetic fields of ordinary radiopulsars. This confirms the suggestion of magnetic field decay during accretion made by Bisnovatyi-Kogan and Komberg (1974).

Recent discoveries of 5 RR in M 15 (Anderson et al., 1990; Manchester, 1990) and 11 RR in 47 Tuc (Manchester et al., 1991) arises problems of evolution of LMXB and RR. The number of RR exceeds number of LMXB in one cluster about 10 times so the lifetime of LMXB must be 10 times less than of RR if their origin is in common. This explanation arises difficulties, connected with formation of large number of binaries in globular cluster by tidal capture. Calculations of Bisnovatyi-Kogan and Romanova (1982) give the number of possible binary much less than that observed in 47 Tuc. The possible explanation may be related with nonmonotonic evolution. In past cores of globular clusters could be more dense with larger binary birthrate. Formation of several binaries could reverse evolution and lead to expansion of the core. RR in 47 Tuc could be remnants of this dense core phase.

6. Gamma-ray burst sources

The number of registered GRB is about 500 and is of the order of number of radiopulsars. If GRB originate on old single nearby NS there is a direct genetic connection between radiopulsars and GRB sources. In models of starquakes with or without nuclear explosion (Bisnovatyi-Kogan et al., 1975; Tsygan, 1975) all NS, including living and dead radiopulsars are the sources of GRB. If energy of explosion leading to GRB is limited by $10^{36} \div 10^{39}$ ergs then they could be observed by existing detectors with sensitivity not better 3×10^{-7} ergs/cm² = F_{lim} from distances less than $200 \div 5000$ pc. Taking into account that sensitivity of most detectors is less than F_{lim} it follows that average distance to observed GRB is about half thickness of the galactic disk in accordance with their isotropic distribution over the sky.

Assume that all NS in Galaxy ($2 \div 7$) $\times 10^8$ stars, are uniformly distributed in the disk with thickness 400 pc and radius 15 kpc and all observed GRB arrive from the part of the disk with radius 200 pc around Sun. Then the number of NS taking part in visible GRB production is equal to $(0.3 \div 1) \times 10^5$. Adopting the average frequency of visible GRB about one per day we obtain that old NS must give birth to GRB every $(0.3 \div 1) \times 10^5$ days or every 100 \div 350 years.

7. Upper limits on the number of neutron stars in Galaxy

Estimate a density on NS not using extrapolation of modern birthrate to earliest stages. Consider two kinds of estimations:

1. Production of heavy elements.

It is established that elements beginning from carbon are ejected into interstellar medium in SN explosions and amount of heavy elements ejected in one explosion is not less than $0.1 M_{\odot}$. For the density of heavy elements $\rho_H \leq 0.03 \rho_{\Gamma}$ and NS mass $1.4 M_{\odot}$ we get

$$\rho_{ns}/\rho_{\Gamma} \leq 0.4 \quad \text{in all SN explosions.} \quad (7)$$

2. Neutrino background.

Formation of NS leaves a signature in form of neutrino background (Bisnovatyi-Kogan and Seidov, 1982). Binding energy of NS $\sim 0.1 M_{\odot}$ is transformed into middle-energy neutrino with $E_{\nu} = 5 \div 30$ Mev. It follows from estimations of Bisnovatyi-Kogan and Seidov (1982), that NS with average density $\rho_{ns} = 0.03 \rho_{\Gamma}$ at constant production rate equal to the present one give the background in Cl - Ar Solar detector about 2×10^{-3} SNU. If most NS were produced in early epoch of violent starbirth, the energy of neutrino in background is less due to redshift and a possibility of detection falls. Using of Ga - Ge neutrino detectors may give better constrains on early formed neutrino background and lower absolute upper limit on density of NS. After estimation of spectra and amount of energy produced during formation of BH in relativistic collapse, one may get similar constrains on density of BH from stellar collapses.

8. References

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DISCUSSION

HORVATH: In currently studied models of type II SN the final fate of the compact remnant turns out to be essentially determined by the mass of the iron collapsing core, which appears to be a non-monotonic function of the total progenitor mass. Would you predict a clear progenitor mass forming BH above and NS below from your magnetorotational driven explosions?

BISNOVATYI-KOGAN: The mentioned main sequence mass of $35 M_{\odot}$ for the boundary between NS and BH is very indefinite and may be considered as a personal opinion, based on experience in stellar evolution and supernova theory.