

S AND g_λ^* -FUNCTIONS ON COMPACT LIE GROUPS

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Abstract

We characterize the Hardy spaces $H^p(G)$ of a compact Lie group G by means of S -functions in analogy with the theorem of Fefferman-Stein for \mathbb{R}^n . We also characterize $H^p(G)$ by means of the g_λ^* -functions.

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1. Introduction

The characterization of $H^p(\mathbb{R}^n)$ by means of S -functions is a well-known result of Fefferman-Stein [4, Theorem 8]. Using previously obtained atomic characterizations of $H^p(G)$ [1], we prove an analogous result for compact connected semisimple Lie groups G . As an application, we show that $\|g_\lambda^*(f)\|_p \leq C \|f\|_{H^p(G)}$. This inequality gives us another characterization of $H^p(G)$ by means of the g_λ^* -function.

The Hardy space $H^p(G)$ of distributions on a connected simply-connected compact group G is defined to be $H^p(G) = \{f \in \mathcal{S}'(G) \mid u_f^* \in L^p(G)\}$ where $u_f^*(x) = \sup_{(y,t) \in \Gamma(x)} |P_t * f(y)|$, P_t is the Poisson kernel associated with the Casimir operator of G , and $\Gamma(x) = \{(y, t) \in G \times \mathbb{R}^+ \mid d(x, y) < t\}$ is the cone with vertex $x \in G$ defined by a bi-invariant metric d on G . For suitable radial functions ϕ on the Lie algebra \mathfrak{t} of a maximal torus T of G (see (3.1) for a complete description), we define the S -function by

$$S_\phi f(x) = \left(\int_{\Gamma(x)} |(f * \phi_t)(y)|^2 t^{-(n+1)} dy dt \right)^{1/2}.$$

Our main result concerning the S -function is:

THEOREM 3.3. *For $f \in \mathcal{S}'(G)$, $f \in H^p(G)$ if and only if $S_\phi(f) \in L^p(G)$. Moreover, $\|u_f^*\|_p \cong \|S_\phi(f)\|_p$.*

For f a distribution on G and $\lambda > 1$, we define the g_λ^* -function of f by

$$g_\lambda^*(f)(x) = \left(\int_0^\infty \int_G \left[\frac{t}{t + d(x, y)} \right]^{\lambda n} |(f * \phi_t)(y)|^2 t^{-(1+n)} dy dt \right)^{1/2}.$$

The g_λ^* -function characterization of $H^p(G)$ that we obtain in Section 4 is contained in these two theorems:

THEOREM 4.1. *Suppose that $f \in \mathcal{S}'(G)$. For $0 < p \leq 1$ and $\lambda > 2/p$, $f \in H^p(G)$ if and only if $g_\lambda^*(f) \in L^p(G)$. Moreover $\|g_\lambda^*(f)\|_p \simeq \|S_\phi(f)\|_p \simeq \|u_f^*\|_p$.*

THEOREM 4.2. *For $p > 1$ and $\lambda > 2/p$, $\|g_\lambda^*(f)\|_p \leq C \|f\|_p$.*

In fact, in this paper we will show that these are characterizations of atomic Hardy space $H_a^p(G)$ as defined in Section 2. The authors have previously demonstrated the equivalence of atomic Hardy space $H_a^p(G)$ and $H^p(G)$.

2. Notation and definitions

Let G be a connected simply-connected compact Lie group of dimension n . Let \mathfrak{g} be the Lie algebra of G and let \mathfrak{t} be the Lie algebra of a fixed maximal torus T of G of dimension ℓ . Let A be a system of positive roots for the pair $(\mathfrak{g}, \mathfrak{t})$. Then $\text{Card}(A) = (n - \ell)/2$. Let $\delta = \sum_{\alpha \in A} \alpha/2$.

If $|\cdot|$ is the norm on \mathfrak{g} induced by the negative of the Killing form B on $\mathfrak{g}^\mathbb{C}$, the complexification of \mathfrak{g} , then $|\cdot|$ induces a bi-invariant metric d on G . Furthermore, since $B|_{\mathfrak{t}^\mathbb{C} \times \mathfrak{t}^\mathbb{C}}$ is non-degenerate, for each complex linear functional $\lambda \in \text{hom}_\mathbb{C}(\mathfrak{t}^\mathbb{C}, \mathbb{C})$ there is a unique $H_\lambda \in \mathfrak{t}^\mathbb{C}$ such that $\lambda(H) = B(H, H_\lambda)$ for $H \in \mathfrak{t}^\mathbb{C}$. The inner product and norm on \mathfrak{t} give rise to an inner product $\langle \cdot, \cdot \rangle$ and norm $\|\cdot\|$ on $\text{hom}_\mathbb{C}(\mathfrak{t}, i\mathbb{R})$ by means of this canonical isomorphism.

The weight lattice P is defined by $P = \{\lambda \in \text{hom}_\mathbb{C}(\mathfrak{t}, i\mathbb{R}) : \lambda(X) \in 2\pi i\mathbb{Z}\}$. The set Λ of dominant weights is defined by $\Lambda = \{\lambda \in P : \langle \lambda, \alpha \rangle \geq 0 \text{ for } \alpha \in A\}$. The set \widehat{G} of equivalence classes of irreducible unitary representations of G is parameterized by $\Lambda : \widehat{G} = \{[U_\lambda]\}_{\lambda \in \Lambda}$. The representation U_λ has dimension d_λ and character $\chi_\lambda(X)$ given by

$$d_\lambda = \prod_{\alpha \in A} \frac{\langle \lambda + \delta, \alpha \rangle}{\langle \delta, \alpha \rangle}; \quad \chi_\lambda(X) = \frac{\sum_{w \in W} \varepsilon(w) e^{i(w(\lambda + \delta), X)}}{\sum_{w \in W} \varepsilon(w) e^{i\langle w\delta, X \rangle}}, \quad (X \in \mathfrak{t})$$

where W is the Weyl group and $\varepsilon(w)$ is the signature of w . Let $\{X_1, \dots, X_n\}$ be an orthonormal basis of \mathfrak{g} . The Casimir operator

$$\Delta = \sum_{i=1}^n X_i^2$$

is an elliptic bi-invariant operator on G that is independent of the choice of basis. Let W_t and P_t be the Gauss-Weierstrass and Poisson kernels defined on $G^+ = G \times \mathbb{R}^+ = G \times (0, \infty)$ by

$$W_t(x) = \sum_{\lambda \in \Lambda} e^{-t(\|\lambda + \delta\|^2 - \|\delta\|^2)} d_\lambda \chi_\lambda(x) \quad (x, t) \in G^+$$

and

$$P_t(x) = \sum_{\lambda \in \Lambda} e^{-t\sqrt{\|\lambda + \delta\|^2 - \|\delta\|^2}} d_\lambda \chi_\lambda(x) \quad (x, t) \in G^+.$$

The solutions to the heat equation

$$\frac{\partial \varphi}{\partial t}(x, t) = \Delta \varphi(x, t) \quad \varphi(g, 0^+) = f(x)$$

and the Poisson equation

$$\frac{\partial^2 \varphi}{\partial t^2}(x, t) + \Delta \varphi(x, t) = 0 \quad \varphi(g, 0^+) = f(x)$$

for $f \in L^1(G)$ are given by $W_t * f$ and $P_t * f$ respectively. Here and elsewhere, Haar measures on compact groups are normalized to have total mass one. All Lebesgue spaces to be discussed will be with respect to such measures.

Let $\Gamma(x) = \{(y, t) \in G^+ \mid d(x, y) < t\}$. For a distribution f in $\mathcal{S}'(G)$, let

$$u_f(x, t) = P_t * f(x) \quad \text{and} \quad u_f^*(x) = \sup_{(y,t) \in \Gamma(x)} |u_f(y, t)|.$$

Then, for $0 < p < \infty$,

$$H^p(G) = \{f \in \mathcal{S}'(G) \mid u_f^* \in L^p(G)\}.$$

The ‘norm’ $\|f\|_{H^p(G)}$ of f in $H^p(G)$ is the Lebesgue norm $\|u_f^*\|_p$. Although $\|\cdot\|_{H^p(G)}$ is not a norm in general, it induces a complete metrizable topology on $H^p(G)$. Since $H^p(G) = L^p(G)$ for $p > 1$, we will restrict our attention to the case $0 < p \leq 1$.

We will also need the atomic Hardy spaces as originally defined by Coifman-Weiss [3] in the context of spaces of homogeneous type. We will actually use the modification for compact groups found in Clerc [2]. For each y in G , let L_y denote

left translation by y in G . Let ε_1 and δ_1 be positive numbers such that $\exp^{-1} \circ L_{x^{-1}}$ is a diffeomorphism from the G -ball $B(x, \varepsilon_1)$ into the ball $B(0, \delta_1)$ of \mathfrak{g} for all x in G . Let $T_x(G)$ be the tangent space of G at x . For a positive integer k and an element y of G , let

$$\mathcal{P}_k(y) = \{P : P = q \circ \exp^{-1} \circ L_{y^{-1}} \text{ for some polynomial } q \text{ on } \mathfrak{g} \text{ of degree } \leq k\}.$$

Let $0 < p \leq 1 \leq q \leq \infty$. Set $k(p) = [n(1/p - 1)]$. A regular (p, q) atom on G is a function $a(x)$ supported in some ball $B(y, \rho)$ ($0 < \rho < \varepsilon_1$) such that

- (i) $\|a\|_q \leq \rho^{n(1/q - 1/p)}$ (size condition);
- (ii) $\int_G a(x)P(x) dx = 0, P \in \mathcal{P}_{k(p)}(y)$ (cancellation condition).

An exceptional atom is a function bounded by 1. The atomic Hardy space $H_a^{p,q}(G)$ is the space of all $f \in \mathcal{S}'(G)$ of the form

$$f = \sum_k c_k a_k, \quad \sum_k |c_k|^p < \infty,$$

the decomposition being in terms of regular (p, q) and exceptional atoms. The ‘norm’ $\|f\|_{p,q,a}$ of f in $H_a^{p,q}(G)$ is defined to be $\inf \left\{ (\sum_k |c_k|^p)^{1/p} \right\}$ taken over all atomic decompositions of f . It is known in the more general context of spaces of homogeneous type that for fixed p , identical atomic Hardy spaces arise for all $q \in [1, \infty]$. We therefore need only consider the $q = \infty$ case. We denote $H_a^{p,\infty}(G)$ by $H_a^p(G)$. We will denote the norm of this space by $\|\cdot\|_{p,a}$.

3. The S -function characterization of $H^p(G)$

Let ϕ be a radial function in $\mathcal{S}(\mathbb{R}^\ell)$ which satisfies

$$(3.1) \quad \begin{aligned} & \text{(i)} \quad \hat{\phi}(0) = 0 \\ & \text{(ii)} \quad \int_0^\infty \phi(s)^2 ds/s = c(\phi) \neq 0. \end{aligned}$$

We define a central function in $C^\infty(G)$ by its restriction to T :

$$(3.2) \quad \phi_t(x) = \sum_{\lambda \in \Lambda} \hat{\phi}(t \|\lambda + \delta\|) d_\lambda \chi_\lambda(x).$$

Let R be defined as in [2] and let μ^R denote the number of singular positive roots (as defined in [2, p. 87]). Let $D^R(H) = \prod_\alpha \sin \alpha(H)/2$, the product being over all positive non-singular roots. For a multi-index $J = (j_1, \dots, j_n)$, let $X^J = X_1^{j_1} X_2^{j_2} \dots X_n^{j_n}$ and let $|J| = j_1 + \dots + j_n$.

LEMMA 3.1. *Suppose that $x \in G$ is conjugate to $\exp H$ for $H \in \mathfrak{t}$. Then there is a constant C independent of x and t such that for any multi-index J and $m \in \mathbb{N}$*

$$(3.3) \quad \begin{aligned} \text{(i)} \quad & |X^J \phi_t(x)| \leq Ct^{-m} \text{ if } t \geq \varepsilon_1, \\ \text{(ii)} \quad & |X^J \phi_t(x)| \leq Ct^{-|J|-n}, \\ \text{(iii)} \quad & |X^J \phi_t(x)| \leq Ct^m \left(\|H\|^{-m-n-|J|} + t^{-\mu^R} D^R(H)^{-1} \right) \text{ if } \|H\| > t. \end{aligned}$$

The proof of this lemma is the same as the proof of [2, Theorem 5.4]. We will continue to denote unimportant constant by C , without distinguishing between different constants, if they have no crucial dependence on objects under consideration.

For any $f \in S'(G)$, we define the S -function of f by

$$(3.4) \quad S_\phi f(x) = \left(\int_{\Gamma(x)} |(f * \phi_t)(y)|^2 t^{-(n+1)} dy dt \right)^{1/2}$$

and the g -function of f by

$$(3.5) \quad g(f)(x) = \left(\int_0^\infty |(f * \phi_t)(x)|^2 dt/t \right)^{1/2}.$$

LEMMA 3.2. $\|g(f)\|_2 = c(\phi) \cdot \|f\|_2$.

PROOF. Since $\|g(f)\|_2^2 = c(\phi) \int_0^\infty \int_G |(f * \phi_t)(x)|^2 dx dt/t$, the lemma follows from (3.1) and the Plancherel Theorem.

THEOREM 3.3. *For $f \in \mathcal{S}'(G)$, $u_f^* \in L^p(G)$ if and only if $S_\phi(f) \in L^p(G)$. Moreover, $\|u_f^*\|_p \cong \|S_\phi(f)\|_p$.*

PROOF. If $u_f^* \in L^p(G)$, then $f \in H_a^p(G)$ [1]. Therefore f has an atomic decomposition $f = \sum_j c_j a_j$ with $\sum |c_j|^p \leq C \|u_f^*\|_p^p$. Now

$$\begin{aligned} \|S_\phi f(x)\|_p^p &= \int_G \left(\int_{\Gamma(x)} \left| \sum_j c_j (a_j * \phi_t)(y) t^{-(n+1)/2} \right|^2 dy dt \right)^{p/2} dx \\ &\leq \sum_j |c_j|^p \int_G \left(\int_{\Gamma(x)} |(a_j * \phi_t)(y) t^{-(n+1)/2}|^2 dy dt \right)^{p/2} dx. \end{aligned}$$

If we show that

$$(3.6) \quad \int_G \left(\int_{\Gamma(x)} |(a * \phi_t)(y)|^2 t^{-(n+1)} dy dt \right)^{p/2} dx \leq C$$

for a constant C independent of the atom a , then

$$\|S_\phi f(x)\|_p^p \leq C \sum_j |c_j|^p \leq C \|f\|_{p,a}^p \leq C \|a_f^*\|_p^p.$$

As the proof of (3.6) for $p \in (0, 1)$ is the same as for the case $p = 1$, we will assume that $p = 1$ for notational simplicity. The proof for exceptional atoms a is an easy consequence of Hölder’s inequality and Lemma 3.2:

$$\begin{aligned} \|S_\phi(a)\|_1 &\leq C \|S_\phi(a)\|_2 \\ &\leq \left(\int_G \int_0^\infty \int_{d(x,y)<t} |(a * \phi_t)(y)|^2 t^{-(n+1)} dy dt dx \right)^{1/2} \\ &\leq C \|g(a)\|_2 \leq C \|a\|_2 \leq C. \end{aligned}$$

Now let $a(x)$ be a regular $(1, \infty)$ atom supported, without loss of generality, in $B(I, \rho)$. Using (i) of (3.3) we may assume that $t < \varepsilon_0$ for some fixed ε_0 . We break up $\int_G |S_\phi(a)(x)| dx$ into two pieces according to whether $d(x, I) \geq 8L\rho$ or $d(x, I) < 8L\rho$ where L is the largest root length. Then

$$\int_{d(x,I)<8L\rho} |S_\phi(a)(x)| dx \leq C\rho^{n/2} \|g(a)\|_2 \leq C\rho^{n/2} \|a\|_2 \leq C.$$

The remaining piece $\int_{d(x,I)\geq 8L\rho} |S_\phi(a)(x)| dx$ of $\|S_\phi(a)\|_1$ is itself broken into two pieces by partitioning each $\Gamma(x)$ into $\Gamma_1(x) = \{(y, t) : d(y, x) < t \leq 2L\rho\}$ and $\Gamma_2(x) = \{(y, t) : d(y, x) < t, 2L\rho < t\}$. We will show that each

$$J_i = \int_{d(x,I)\geq 8L\rho} \left| \int_{\Gamma_i(x)} \left(\int_{B(I,\rho)} a(\xi)\phi_t(\xi^{-1}y) d\xi \right)^2 t^{-(n+1)} dy dt \right|^{1/2} dx \quad (i = 1, 2)$$

is bounded independently of a . For ξ, x and (y, t) in the integration in J_1 ,

$$d(\xi, I) \geq d(x, I) - d(y, \xi) - d(x, y) \geq d(x, I)/4 \geq 2L\rho > t.$$

Therefore, by (iii) of (3.3),

$$\begin{aligned} |J_1| &\leq C \int_{d(x,I)\geq 2L\rho} \left\{ \int_{\Gamma_1(x)} \sup_{\xi \in B(y,\rho)} d(I, \xi)^{-2(n+1)} t^{1-n} dy dt \right\}^{1/2} dx \\ &\quad + C \|a\|_\infty \int_G \left\{ \int_0^{2L\rho} \int_G \int_G D^R(\xi^{-1}y)^{-2} d\xi t^{2n-1} dy dt \right\}^{1/2} dx. \end{aligned}$$

The second summand is obviously bounded; for the first,

$$\int_{d(x,I)\geq 2L\rho} d(I, x)^{-(n+1)} \left\{ \int_{d(x,y)<t} \int_0^{2L\rho} t^{1-n} dy dt \right\}^{1/2} dx \leq C\rho^{-1} \left\{ \int_0^{2L\rho} t dt \right\}^{1/2} \leq C.$$

We estimate J_2 by partitioning each $\Gamma_2(x)$ into two pieces

$$\gamma_1(x) = \{(y, t) : d(y, x) < t, d(y, B(I, \rho)) \geq 4L\rho, t > 2L\rho\}$$

and

$$\gamma_2(x) = \{(y, t) : d(y, x) < t, d(y, B(I, \rho)) < 4L\rho, t > 2L\rho\}.$$

Write $J_2 \leq I_1 + I_2$ where

$$I_i = \int_{d(x, I) \geq 8L\rho} \left| \int_{\gamma_i(x)} \left(\int_{B(I, \rho)} a(\xi) \phi_t(\xi^{-1}y) d\xi \right)^2 t^{-(n+1)} dy dt \right|^{1/2} dx \quad (i = 1, 2).$$

Since a is a $(1, \infty)$ -atom,

$$|I_2| \leq C \|a\|_\infty \rho^{n+1} \int_{B(I, 8L\rho)^c} \left(\int_{\gamma_2(x)} \mathcal{M}\phi(y, t) t^{-(n+1)} dy dt \right)^{1/2} dx$$

where

$$\mathcal{M}\phi(y, t) = \sup \left\{ |X_j \phi_t(\xi)|^2 : \xi \in B(y, \rho), 1 \leq j \leq n \right\}.$$

If $y \in \gamma_2(x)$, then $t > d(y, x) > d(x, I) - d(y, I) > d(x, I)/4$. Therefore, by (ii) and (iii) of (3.3), we have

$$\begin{aligned} |I_2| &\leq C\rho \int_{d(x, I) \geq 8L\rho} \left| \int_{d(x, I)/4}^{\epsilon_0} t^{-(2n+3)} dt \right|^{1/2} dx \\ &\quad + C\rho \int_G \left(\int_{2L\rho}^{\epsilon_0} t^{2n} \int_{d(x, y) < t} \sup \left\{ t^{-2\mu^R} D^R(\xi)^{-2} : \xi \in B(y, \rho) \right\} dy dt \right)^{1/2} dx. \end{aligned}$$

The first summand is easily seen to be bounded and the second is bounded by $C \int_G D^R(y)^{-2} dy \leq C$ (cf. [2, Lemma 6.4]).

In the first step in estimating I_1 , we also use (ii) and (iii) of (3.3) as well as [2, Lemma 6.4] to obtain

$$|I_1| \leq C\rho \int_{B(I, 8L\rho)^c} \left[\int_{\gamma_1(x)} \sup_{\xi \in B(y, \rho)} (t + \|\xi\|)^{-2(n+1)} t^{-(n+1)} dy dt \right]^{1/2} dx + C.$$

For any $y \in \gamma_1(x)$ and $\xi \in B(y, \rho)$, $d(\xi, I) \geq d(y, I)$ and $t + d(y, I) > (d(x, I) + t)/4$. Therefore

$$\begin{aligned} |I_1| &\leq C\rho \int_{B(I, 8L\rho)^c} \left[\int_{\gamma_1(x)} (t^{1/4}(t + d(y, I))^{-(n+1)})^2 t^{-n-3/2} dy dt \right]^{1/2} dx \\ &\leq C\rho \int_{B(I, 8L\rho)^c} \left[\int_{2L\rho}^{\epsilon_0} t^{-3/2} dt \right]^{1/2} d(x, I)^{-n-3/4} dx \leq C. \end{aligned}$$

This completes the proof that $\|S_\phi(a)\|_1 \leq C$ with C independent of the $(1, \infty)$ -atom a .

We turn to the other direction of the equivalence, assuming that $S_\phi(f) \in L^p(G)$. Let Ψ be a radial function in $\mathcal{S}'(\mathbb{R}^\ell)$ that satisfies

$$(3.7) \quad \begin{aligned} & \text{(i)} \quad \text{supp}(\Psi) \subseteq \{\theta : |\theta| \leq 1\} \\ & \text{(ii)} \quad \int_{\mathbb{R}^\ell} \theta^I \Psi(\theta) d\theta = 0 \quad (I \in \mathbb{N}^\ell, |I| \leq 3n + 3 + 2n(1/p - 1/2)) \\ & \text{(iii)} \quad \int_0^\infty \hat{\phi}(t) \hat{\Psi}(t) dt/t = 1. \end{aligned}$$

By the Calderon reproducing formula on G , any $f \in \mathcal{S}'$ has a reproducing transformation

$$(3.8) \quad f(x) = \int_{G^+} (f * \phi_t)(y) \Psi_t(xy^{-1}) dy t^{-1} dt$$

which we break up as the sum of I_1 and I_2 where

$$(3.9) \quad \begin{aligned} I_1(x) &= \int_0^\varepsilon \int_G (f * \phi_t)(y) \Psi_t(xy^{-1}) dy t^{-1} dt, \quad \text{and} \\ I_2(x) &= \int_\varepsilon^\infty \int_G (f * \phi_t)(y) \Psi_t(xy^{-1}) dy t^{-1} dt \end{aligned}$$

for a small ε that will be determined later. For this fixed ε , there is a constant C_ε such that

$$\|I_2\|_\infty \leq \left[\int_\varepsilon^\infty \int_G |(f * \phi_t)(y)|^2 t^{-(n+1)} dy dt \right]^{1/2}.$$

Since G is compact, there are elements x_1, \dots, x_N ($N = N(G, \varepsilon)$) such that G is covered by the open $\varepsilon/4$ -balls centered at these points. Let χ_i denote the characteristic function of $B(x_i, \varepsilon)$. Then

$$\begin{aligned} \|S_\phi(f)\|_p^p &\geq N^{-1} \int_G \sum_{i=1}^N \chi_i(x) \left[\int_\varepsilon^\infty \int_{d(x,y)<t} |(f * \phi_t)(y)|^2 t^{-(n+1)} dy dt \right]^{p/2} dx \\ &\geq C_{N,p} \int_G \left[\sum_{i=1}^N \int_\varepsilon^\infty \int_{d(x,y)<t} \chi_i(x) |(f * \phi_t)(y)|^2 t^{-(n+1)} dy dt \right]^{p/2} dx \\ &\geq C_{N,p} \int_G \left[\sum_{i=1}^N \int_\varepsilon^\infty \int_{B(x_i, \varepsilon/4)} |(f * \phi_t)(y)|^2 t^{-(n+1)} dy dt \right]^{p/2} dx \\ &\geq C_{N,p} \left[\int_\varepsilon^\infty \int_G \chi_i(x) |(f * \phi_t)(y)|^2 t^{-(n+1)} dy dt \right]^{p/2} \\ &\geq C_{p,\varepsilon} \|I_2\|_\infty^p. \end{aligned}$$

Thus, we can find a constant $C_{p,\varepsilon}$ depending only on p and ε such that $I_2(x) = C_{p,\varepsilon} a(x)$ where $a(x)$ is an exceptional atom and $|C_{p,\varepsilon}| \leq C \|S_\Phi(f)\|_p$.

To estimate $I_1(x)$, we let xy^{-1} be conjugate to $\exp\theta \in T$. Then $D(xy^{-1}) = \prod_{\alpha \in A} \sin \frac{1}{2}\alpha(\theta)$. There is a polynomial $P_{n(p)}$ of degree $2n + 2 + n(1/p - 1/2)$ such that

$$(3.10) \quad \begin{aligned} I_1(x) = & C \int_0^\varepsilon \int_G \left\{ \left(\prod_{\alpha \in A} \alpha(\theta) + \sum_{\alpha \in A} C_\alpha \alpha(\theta)^3 \prod_{\beta \in A} \beta(\theta) + \dots + P_{n(p)}(\theta) \right) \times \right. \\ & \left. D^{-1}(\theta) \Psi_t(xy^{-1})(f * \phi_t)(y) \right\} dy t^{-1} dt \\ & + C \int_0^\varepsilon \int_G (f * \phi_t)(y) R(\theta) D^{-1}(\theta) \Psi_t(xy^{-1}) dy t^{-1} dt \end{aligned}$$

where $R(\theta)$ is a $C^\infty(\mathbb{R}^\ell)$ -function such that $R(\theta)D(\theta)^{-1} = O(\|\theta\|^{n+3+n(1/p-1/2)})$ and $X_i R(\theta)D(\theta)^{-1} = O(\|\theta\|^{n+2+n(1/p-1/2)})$. As a consequence of these estimates, we have $\|R(x)D^{-1}(x)\Psi_t(x)\|_\infty \leq Ct^{3+n(1/p-1/2)}$.

To complete the proof of Theorem 3.1, we will prove that each term in (3.10) has a suitable atomic decomposition. There are two types of terms that we must deal with. We will show that

$$I_{1,1}(x) = \int_0^\varepsilon \int_G (f * \phi_t)(y) \Psi_t(xy^{-1}) \prod_{\alpha \in A} \frac{\alpha(\theta)}{\sin \alpha(\theta)} dy t^{-1} dt$$

and

$$I_R(x) = \int_0^\varepsilon \int_G (f * \phi_t)(y) \Psi_t(xy^{-1}) R(\theta) D(\theta)^{-1} dy t^{-1} dt$$

have atomic decompositions

$$(3.11) \quad I_{1,1}(x) = \sum_j \lambda_j a_j(x), \quad I_R(x) = \sum_j \nu_j b_j(x)$$

with each a_j a $(p, 2)$ -atom, each b_j an exceptional atom and $\sum_j |\lambda_j|^p \leq \|S_\Phi(f)\|_p^p$ and $\sum_j |\nu_j|^p \leq \|S_\Phi(f)\|_p^p$. All other terms in (3.10) are handled in the same way as $I_{1,1}$.

Let ε_1 be as in the definition of atoms. For a choice of $\varepsilon \in (0, \varepsilon_1/32)$, the ball $B(x, 16\varepsilon)$ is contained in a local coordinate chart $\{V_x, \eta\}$ with $\text{diam}(V_x) < \varepsilon_1$. Let $\{x_1, \dots, x_N\}$ be such that $G = \cup_{j=1}^N B(x_j, \varepsilon)$. Let $U_j = B(x_j, \varepsilon)$, let $\chi_j(x) = \chi_{U_j}(x)$ and set $\xi_j(x) = \chi_j(x) / \sum_{i=1}^N \chi_i(x)$. Let $M(\theta) = M(xy^{-1}) = \prod_{\alpha \in A} \alpha(\theta) / \sin \alpha(\theta)$. Then $I_{1,1}(x) = \sum_{j=1}^N F_j(x)$ and $I_R(x) = \sum_{j=1}^N G_j(x)$ where

$$F_j(x) = C \int_0^\varepsilon \int_G \xi_j(y) (f * \phi_t)(y) \Psi_t(xy^{-1}) M(\theta) dy t^{-1} dt$$

and

$$G_j(x) = C \int_0^\varepsilon \int_G \xi_j(y) (f * \phi_t)(y) \Psi_t(xy^{-1}) R(\theta) D(\theta)^{-1} dy t^{-1} dt.$$

It suffices to show that each F_j and G_j has an atomic decomposition of the indicated type. Henceforth we drop the index j .

Since $\{16 \cdot U, \eta\}$ is contained in a local coordinate chart, we may assume without loss of generality that $\eta(U)$ is the open cube of side length ε centered at $0 \in \mathbb{R}^n$ and that $d(x, y) = |\eta(x) - \eta(y)|$. We will write $\ell(B)$ for the sidelength of a dyadic cube in $\eta(U)$ and write $|B|$ for $|\eta^{-1}(B)|$. Let

$$\mathcal{B} = \{I_B : (y, t) \in I_B \text{ if and only if } y \in B \text{ and } \ell(B)/2 < t \leq \ell(B)\}.$$

For each $I_B \in \mathcal{B}$, we will write \tilde{I}_B for $(\ell(B)/2, \ell(B)) \times \eta^{-1}(B)$. If

$$(3.12) \quad f_B(x) = \int_{\tilde{I}_B} \xi(y) (f * \phi_t)(y) \Psi_t(xy^{-1}) M(\theta) dy t^{-1} dt$$

and

$$(3.13) \quad g_B(x) = \int_{\tilde{I}_B} \xi(y) (f * \Phi_t)(y) \Psi_t(xy^{-1}) R(\theta) D(\theta)^{-1} dy t^{-1} dt$$

then

$$(3.14) \quad F = \sum_{I_B \in \mathcal{B}} f_B \quad \text{and} \quad G = \sum_{I_B \in \mathcal{B}} g_B$$

in S' .

Observe first of all that f_B and g_B are C^∞ -functions supported in $4\eta^{-1}(B)$ since Ψ is C^∞ and the integrands in (3.13) and (3.14) vanish unless $B \cap B(x, t)$ is not an empty set for some $t \in \tilde{I}_B$. Also

$$(3.15) \quad \int_G f_B(x) P(x) dx = 0$$

for all polynomials P with degree at most $n[1/p - 1] + n$. In fact, if \log is the inverse of the local exponential map, then

$$\int_G f_B(x) P(x) dx = \int_{I_B} (f * \phi_t)(y) \xi(y) \int_G \Psi_t(xy^{-1}) M(xy^{-1}) P(\log(xy^{-1}y)) dx dy dt / t.$$

To prove (3.15), it therefore suffices to show that

$$\int_G \Psi_t(x) M(x) P(\log(xy)) dx = 0$$

for any fixed $t < \varepsilon$ and $y \in G$. Since $\Psi_t \cdot M$ is a central function, we need only prove that

$$\int_G \int_G \Psi_t(x) M(x) P(\log(zxz^{-1}y)) dx dz = 0$$

or

$$\int_G \Psi_t(x) M(x) \int_G P(\log(zxz^{-1}y)) dz dx = 0.$$

But $\int_G P(\log(zxz^{-1}y)) dz$ is a class function that is a polynomial of $\theta = \log(x)$ with degree at most $n[1/p - 1] + n$. Thus it suffices to show that

$$\int_G \Psi_t(x) M(x) (\log(x))^J dx = C \int_t \Psi_t(\exp \theta) M(\exp \theta) \theta^J D^2(\theta) d\theta = 0$$

for all multi-indices J with $|J| \leq n[1/p - 1] + n$. This follows by Poisson summation in view of the choice of Ψ_t .

From the preceding observations, we know that each f_B is a constant multiple of a (p, ∞) -atom. It does not yet follow, however, that the first equation of (3.14) is an atomic decomposition of F since the norms of the f_B 's do not sum properly. For each $I_B \in \mathcal{B}$ we define

$$S_B = \left(\int_{I_B} |(f * \phi_t)(y)|^2 dy t^{-1} dt \right)^{1/2}.$$

We claim that for all $I_B \in \mathcal{B}$ and all multi-indices J ,

$$(3.16) \quad \begin{aligned} (i) \quad & \|X^J f_B\|_\infty \leq C S_B |B|^{-1/2-|J|/n}, \text{ and} \\ (ii) \quad & \|g_B\|_\infty \leq C S_B |B|^{1/p+2/n} \end{aligned}$$

where C depends on J but not on B .

By Schwarz's inequality,

$$|X^J f_B(x)| \leq C S_B \left(\int_{I_B} |X^J(\Psi_t(y^{-1}x)M(yx^{-1}))|^2 dy t^{-1} dt \right)^{1/2}.$$

Therefore $\|\Psi_t M\|_\infty \leq C t^{-n} \leq C |B|^{-1}$, $\|X^J(\Psi_t M)\|_\infty \leq C |B|^{-1-|J|/n}$, and $\|M\|_\infty \leq C$; thus (i) of (3.16) follows. Similarly,

$$\begin{aligned} |g_B(x)| &\leq C S_B \left(\int_{I_B} |\Psi_t(y^{-1}x)D(y^{-1}x)^{-1}R(y^{-1}x)|^2 dy t^{-1} dt \right)^{1/2} \\ &\leq C S_B |B|^{1/2} \sup \{ \|\Psi_t D^{-1}R\|_\infty : \ell(B)/2 \leq t \leq \ell(B) \} \\ &\leq C S_B |B|^{1/p+2/n} \end{aligned}$$

which completes the proof of (3.16).

For each integer k , let $\Omega_k = \{x : S_\Phi f(x) > 2^k\}$ and let \mathcal{B}_k be defined by

$$\mathcal{B}_k = \{I_B \in \mathcal{B} : |\eta^{-1}(B) \cap \Omega_k| > |\eta^{-1}(B)/4| \text{ and } |\eta^{-1}(B) \cap \Omega_{k+1}| \leq |\eta^{-1}(B)/4|\}$$

where $B/2$ is any one of the 2^n subdyadic cubes of B . It is easy to see that $\Omega_{k+1} \subset \Omega_k$ and that each I_B must belong to precisely one \mathcal{B}_k . We claim that there is a C independent of k such that

$$(3.17) \quad \sum_{I_B \in \mathcal{B}_k} S_B^2 \leq C2^{2k} |\Omega_k|.$$

To see this, let M_{HL} denote the Hardy-Littlewood maximal function and let $\tilde{\Omega}_k = \{x : M_{HL}(\chi_{\Omega_k})(x) > 4^{-n}\}$. Observe that $\Omega_k \subset \tilde{\Omega}_k$ and that $|\tilde{\Omega}_k| \leq C|\Omega_k|$ by the Hardy-Littlewood maximal theorem. These imply that

$$(3.18) \quad \int_{\tilde{\Omega}_k \setminus \Omega_{k+1}} |S_\Phi f(x)|^2 dx \leq 2^{2k+2} |\tilde{\Omega}_k| \leq C2^{2k} |\Omega_k|.$$

Let

$$v_k(y, t) = \left| \left\{ x \in \tilde{\Omega}_k \setminus \Omega_{k+1} : d(x, y) < t \right\} \right|.$$

Notice that

$$\int_{\tilde{\Omega}_k \setminus \Omega_{k+1}} |S_\Phi f(x)|^2 dx = \int_0^\infty \int_G |(f * \Phi_t)(y)|^2 v_k(y, t) dy t^{-1-n} dt$$

and therefore

$$\int_{\tilde{\Omega}_k \setminus \Omega_{k+1}} |S_\Phi f(x)|^2 dx \geq \sum_{I_B \in \mathcal{B}_k} \int_{I_B} |(f * \Phi_t)(y)|^2 v_k(y, t) dy t^{-1-n} dt.$$

In view of this and (3.18), in order to obtain (3.17), it suffices to show that there is a constant C independent of k such that

$$(3.19) \quad v_k(y, t) \geq Ct^n \quad \text{for all } I_B \in \mathcal{B}_k \text{ and } (y, t) \in \tilde{I}_B.$$

Let $I_B \in \mathcal{B}_k$ and $(y, t) \in \tilde{I}_B$. Since $|\eta^{-1}(B) \cap \Omega_k| > |\eta^{-1}(B)/4|$, it follows that $M_{HL}(\chi_{\Omega_k})(x) > 4^{-n}$ for $x \in \eta^{-1}(B)$. Therefore $\eta^{-1}(B) \subset \tilde{\Omega}_k$. Also, since $|\eta^{-1}(B) \cap \Omega_{k+1}| \leq |\eta^{-1}(B)/4|$, it follows that $|\eta^{-1}(B) \setminus \Omega_{k+1}| \geq |3\eta^{-1}(B)/4|$. Thus,

$$\begin{aligned} \left\{ x \in \tilde{\Omega}_k \setminus \Omega_{k+1} : d(x, y) < t \right\} &\supseteq \left\{ x \in \tilde{\Omega}_k \setminus \Omega_{k+1} : d(x, y) < \ell(B)/2 \right\} \\ &\supseteq \eta^{-1}(B)/2 \cap \tilde{\Omega}_k \setminus \Omega_{k+1} \supseteq (\eta^{-1}(B)/2) \setminus \Omega_{k+1}. \end{aligned}$$

It follows that

$$\begin{aligned} v_k(y, t) &\geq |(\eta^{-1}(B)/2) \setminus \Omega_{k+1}| \\ &\geq |\eta^{-1}(B)/2| - |\eta^{-1}(B) \cap \Omega_{k+1}| \\ &\geq |\eta^{-1}(B)/4| \\ &\geq Ct^n \end{aligned}$$

proving (3.19) and, as noted, (3.17).

We can define a partial ordering on \mathcal{B}_k by inclusion. Let $\{B^i\}$ be an enumeration of the maximal elements of \mathcal{B}_k . Each $B \in \mathcal{B}_k$ satisfies $B \subset B^i$ for some i ; for every $B \in \mathcal{B}_k$ we choose such an $i = i(B)$. For every i , let $\mathcal{B}_k^i = \{B \in \mathcal{B}_k : i(B) = i\}$. Thus, $\mathcal{B}_k = \bigcup_i \mathcal{B}_k^i$ disjoint. Define

$$\varphi_k^i = \sum_{B \in \mathcal{B}_k^i} f_B \quad \text{and} \quad \gamma_k^i = \sum_{B \in \mathcal{B}_k^i} g_B.$$

We claim that there exists a C independent of k and i such that

$$(3.20) \quad \begin{aligned} (i) \quad &\|\gamma_k^i\|_\infty \leq C \sum_{B \in \mathcal{B}_k^i} S_B |B|^{1/p+2/n} \\ (ii) \quad &\|\varphi_k^i\|_\infty \leq C \sum_{B \in \mathcal{B}_k^i} S_B^2. \end{aligned}$$

The first estimate of (3.20) is immediate from (3.16). For the second estimate, we let B_1, B_2, \dots be an enumeration of \mathcal{B}_k^i ordered so that $|B_r| \geq |B_s|$ if $r \leq s$. In the proof of this estimate, we will write f_r and S_r for f_{B_r} and S_{B_r} . Then

$$\|\varphi_k^i\|_2^2 = \sum_r \|f_r\|_2^2 + 2 \operatorname{Re} \sum_{r < s} \int_G f_r \bar{f}_s.$$

Now $\|f_r\|_2^2 = \int_{4\eta^{-1}(B_r)} |f_j|^2 \leq C |4B_j| |B_j|^{-1} S_j^2$. To estimate the cross terms, we need only consider r and s such that $4B_r \cap 4B_s \neq \emptyset$, for $f_r \bar{f}_s$ vanishes identically otherwise. Therefore we suppose that $r < s$ and $4B_r \cap 4B_s \neq \emptyset$. We let x_s be the center of B_s and let $P_{r,s}$ be the Taylor polynomial of f_r at x_s of degree $a = n[1/p - 1] + n/2$. Then, by (3.15) we have

$$\begin{aligned} \left| \int_G f_r(x) \bar{f}_s(x) dx \right| &= \left| \int_G (f_r(x) - P_{r,s}(x)) \bar{f}_s(x) dx \right| \\ &= \left| \int_G \left(\int_{\bar{I}_{B_r}} \xi(y) (f * \Phi_t)(y) \Psi_t(y^{-1}x) M(y^{-1}x) dy t^{-1} dt - P_{r,s}(x) \right) \bar{f}_s(x) dx \right| \\ &\leq C \int_{4\eta^{-1}(B_s)} \sum_{|J| \leq a+1} \|X^J f_r\|_\infty \|f_s\|_\infty d(x, x_s)^{a+1} dx \\ &\leq C \sum_{|J| \leq a+1} |B_s|^{1+(a+1)/n} |B_r|^{-1/2-(a+1)/n} |B_s|^{-1/2} S_r S_s \\ &\leq C (|B_s| / |B_r|)^{1/2+(a+1)/n} S_r S_s. \end{aligned}$$

For these indices we set $\beta_{rs} = (|B_s|/|B_r|)^{1/2+(a+1)/n}$ and we set $\beta_{rs} = 0$ otherwise. We must show that $\sum_{rs} \beta_{rs} S_r S_s \leq C \sum_s S_s^2$ for some C . To do this it suffices to show that there is a constant C such that

$$(3.21) \quad \sum_r \beta_{rs} < C \text{ for all } s \quad \text{and} \quad \sum_s \beta_{rs} < C \text{ for all } r.$$

If so, $\sum_r (\sum_s \beta_{rs} S_s)^2 \leq \sum_r (\sum_s \beta_{rs}) (\sum_s \beta_{rs} S_s^2) \leq C \sum_s S_s^2$ and therefore

$$\sum_{rs} \beta_{rs} S_r S_s \leq (\sum_r S_r^2)^{1/2} \left(\sum_r (\sum_s \beta_{rs} S_s^2) \right)^{1/2} \leq C \sum_s S_s^2.$$

We turn to (3.21). For each $m \in \mathbb{N}$ there are at most $16^n 2^{mn}$ values of s such that $|B_s| = 2^{-mn} |B_r|$ and $4B_r \cap 4B_s \neq \emptyset$. For each s there are at most 16^n values of r such that $|B_r| = 2^{mn} |B_s|$ and $4B_r \cap 4B_s \neq \emptyset$. Therefore

$$\begin{aligned} \sum_r \beta_{rs} &\leq C \sum_{m=0}^{\infty} 2^{mn} 2^{-(mn/2)-m(a+1)} \leq C \sum_{m=0}^{\infty} 2^{m(n/2-a-1)} \leq C \\ \text{and} \quad \sum_s \beta_{rs} &\leq C \sum_{m=0}^{\infty} 2^{-m(n/2+a+1)} \leq C. \end{aligned}$$

Recall that $F(x) = \sum_{B \in \mathcal{B}} f_B(x) = \sum_{ik} \varphi_k^i(x)$. Let $\lambda_k^i = \|\varphi_k^i\|_2 / |4B^i|^{1/2-1/p}$ and $a_k^i(x) = \varphi_k^i(x)/\lambda_k^i$. Then, by [5, p. 240], $F(x) = \sum_{ik} \lambda_k^i a_k^i(x)$ is an atomic decomposition in which each a_k^i is a $(p, 2)$ -atom and $\sum_{ik} |\lambda_k^i|^p \leq \|S(f)\|_p^p$. Thus, (3.11) is finally proved.

Now let $v_k^i = C \sum_{B \in \mathcal{B}_k^i} S_B |B|^{1/p+2/n}$ and let $b_k^i(x) = C \gamma_k^i(x)/v_k^i$. By (3.20), $b_k^i(x)$ is an exceptional atom. Moreover, for $\kappa = 2/(2-p) > 1$,

$$\sum_{ik} |v_k^i|^p \leq C \sum_{ik} \sum_{B \in \mathcal{B}_k^i} S_B^p |B|^{1+2p/n} \leq C \sum_k \left(\sum_{B \in \mathcal{B}_k} S_B^2 \right)^{p/2} \left(\sum_i \sum_{B \in \mathcal{B}_k^i} |B|^\kappa \right)^{1/\kappa}.$$

Since there are, for each B^i , at most 2^{mn} cubes $B \in \mathcal{B}_k^i$ such that $|B| = 2^{-mn} |B^i|$ we conclude that

$$\sum_{B \in \mathcal{B}_k^i} |B|^\kappa = \sum_{m=1}^{\infty} 2^{mn} 2^{-mn\kappa} |B^i|^\kappa \leq C |B^i|.$$

Thus,

$$\sum_{ik} |v_k^i|^p \leq C \sum_k 2^{kp} |\Omega_k|^{p/2} |\Omega_k|^{1-p/2} = C \sum_k 2^{kp} |\Omega_k|$$

and the required atomic decomposition $I_R(x) = \sum_{ik} v_k^i b_k^i(x)$ has been proved.

4. The characterization of $H^p(G)$ by the g_λ^* -function

For $f \in \mathcal{S}'(G)$ and $\lambda > 1$, we define the g_λ^* -function of $f(x)$ by

$$g_\lambda^*(f)(x) = \left(\int_0^\infty \int_G \left[\frac{t}{t + d(x, y)} \right]^{\lambda n} |(f * \phi_t)(y)|^2 t^{-(1+n)} dy dt \right)^{1/2}.$$

THEOREM 4.1. *Suppose that $f \in \mathcal{S}'(G)$. For $0 < p \leq 1$ and $\lambda > 2/p$, $f \in H^p(G)$ if and only if $g_\lambda^*(f) \in L^p(G)$. Moreover $\|g_\lambda^*(f)\|_p \simeq \|S_\phi(f)\|_p \simeq \|u_f^*\|_p$.*

PROOF. Suppose that $0 < p \leq 1$ and $\lambda > 2/p$. By Theorem 3.1 we need only check that $\|S_\phi(f)\|_p \simeq \|u_f^*\|_p$. Since $S_\phi(f)(x) \leq Cg_\lambda^*(f)(x)$, only the estimate $\|g_\lambda^*(f)\|_p \leq C\|S_\phi(f)\|_p$ requires further proof. As in the proof of Theorem 3.3, it suffices to show that there is a constant C such that for any atom $a(x)$,

$$(4.1) \quad \int_G \left(\int_0^\infty \int_G \left[\frac{t}{t + d(x, y)} \right]^{\lambda n} |(a * \phi_t)(y)|^2 t^{-(1+n)} dy dt \right)^{p/2} dx \leq C.$$

We will supply details only for the case $p = 1$. If $a(x)$ is an exceptional atom, then

$$\begin{aligned} \|g_\lambda^*(a)\|_1 &\leq C \|g_\lambda^*(a)\|_2 \\ &= C \int_G \left(\int_0^\infty \int_G \left[\frac{t}{t + d(x, y)} \right]^{\lambda n} |(a * \phi_t)(y)|^2 t^{-(1+n)} dy dt \right) dx \\ &= C \int_G \left(\left(\int_0^\infty \int_{d(x,y) < t} + \int_0^\infty \int_{d(x,y) > t} \right) \left[\frac{t}{t + d(x, y)} \right]^{\lambda n} \times \right. \\ &\quad \left. |(a * \phi_t)(y)|^2 t^{-(1+n)} dy dt \right) dx. \end{aligned}$$

By Theorem 3.3, the first summand is bounded by $C\|a\|_2 \leq C$. The second summand is bounded by

$$C \int_G \int_0^\infty |(a * \phi_t)(y)|^2 t^{-1} dt dy \leq C \|g(a)\|_2 \leq C \|a\|_2 \leq C.$$

For a regular $(1, \infty)$ -atom a , we may assume that the support of a is contained in $B(I, \rho)$ with ρ sufficiently small. Our analysis will be based on Lemmas 2.4 and 6.4 of [2].

We write

$$\|g_\lambda^*(a)\|_1 = \int_{d(x,I) \geq 8\rho} |g_\lambda^*(a)(x)| dx + \int_{d(x,I) < 8\rho} |g_\lambda^*(a)(x)| dx = I_1 + I_2.$$

By Schwarz’s inequality,

$$|I_2| \leq C\rho^{n/2} \left(\int_G \int_0^\infty \int_G \left[\frac{t}{t+d(x,y)} \right]^{\lambda n} |(a * \Phi_t)(y)|^2 t^{-(1+n)} dy dt dx \right)^{1/2}$$

and therefore $|I_2| \leq C\rho^{n/2} \|a\|_2 \leq C$ as in the case of exceptional atoms. To estimate I_1 , note that

$$|I_1| \leq \int_{d(x,I) \geq 8\rho} \left(\int_{d(x,y) > t} \left[\frac{t}{t+d(x,y)} \right]^{\lambda n} |(a * \phi_t)(y)|^2 t^{-(1+n)} dy dt \right)^{1/2} dx + C \|S_\phi(a)\|_1.$$

It suffices, then, to estimate the double integral above. We do this by breaking it into three pieces:

$$\begin{aligned} L_1 &= \int_{d(x,I) \geq 8\rho} \left(\int_{\Delta_{1,x}} \left[\frac{t}{t+d(x,y)} \right]^{\lambda n} |(a * \phi_t)(y)|^2 t^{-(1+n)} dy dt \right)^{1/2} dx, \\ L_2 &= \int_{d(x,I) \geq 8\rho} \left(\int_{\Delta_{2,x}} \left[\frac{t}{t+d(x,y)} \right]^{\lambda n} |(a * \phi_t)(y)|^2 t^{-(1+n)} dy dt \right)^{1/2} dx, \\ L_3 &= \int_{d(x,I) \geq 8\rho} \left(\int_\rho^\infty \int_{d(x,y) > t} \left[\frac{t}{t+d(x,y)} \right]^{\lambda n} |(a * \phi_t)(y)|^2 t^{-(1+n)} dy dt \right)^{1/2} dx \end{aligned}$$

where

$$\begin{aligned} \Delta_{1,x} &= \{(y, t) : d(y, x) > t, 0 < t < \rho, d(y, B(I, \rho)) \geq 2\rho\} \text{ and} \\ \Delta_{2,x} &= \{(y, t) : d(y, x) > t, 0 < t < \rho, d(y, B(I, \rho)) < 2\rho\}. \end{aligned}$$

We start with L_2 . Notice that for any $(y, t) \in \Delta_{2,x}$ and $x \notin B(I, 8\rho)$, $d(x, y) > d(x, I) - d(y, I) \geq d(x, I)/2$. Combine this with the estimate $\|a * \phi_t\|_\infty \leq \|a\|_\infty \|\phi_t\|_1 \leq C \|a\|_\infty$ and [2, Lemma 3.4] to get

$$\begin{aligned} L_2 &\leq C\rho^{-n} \int_{d(x,I) \geq 8\rho} d(x, I)^{-\lambda n/2} \left(\int_0^\rho t^{\lambda n - n - 1} \int_{d(y,I) \leq 4\rho} dy dt \right)^{1/2} dx + C \\ &\leq C\rho^{-n - \lambda n/2 + n} \rho^{n/2} \rho^{-n/2 + \lambda n/2} + C. \end{aligned}$$

To estimate L_1 , note that for $(y, t) \in \Delta_{1,x}$ and $\xi \in B(y, \rho)$, $d(\xi, I) \geq d(y, I) - d(\xi, I) \geq d(y, I)/2$. Let

$$G(x, y, t) = \left(\frac{t}{t+d(x,y)} \right)^{\lambda n} \left(\frac{td(y, I)}{[t^2 + d(y, I)^2]^{(n+3)/2}} \right)^2 t^{1-n}.$$

By Lemma 3.1, we conclude that

$$\begin{aligned}
 |L_1| &\leq \int_{d(x,I) \geq 8\rho} \left(\int_0^\rho \int_{\Gamma(x)^c \cap B(I,\rho)^c} \left[\frac{t}{t+d(x,y)} \right]^{\lambda n} \left| \sup_{\xi \in B(y,\rho)} \phi_t(\xi) \right|^2 t^{-(1+n)} dy dt \right)^{1/2} dx \\
 &\quad + C \\
 &= L_{1,1} + L_{1,2} + C
 \end{aligned}$$

where

$$\begin{aligned}
 L_{1,1} &= \int_{d(x,I) \geq 8\rho} \left(\int_0^\rho \int_{\Gamma(x)^c \cap B(I,\rho)^c \cap \{y: d(y,I) > d(x,I)/2\}} G(x,y,t) dy dt \right)^{1/2} dx \\
 L_{1,2} &= \int_{d(x,I) \geq 8\rho} \left(\int_0^\rho \int_{\Gamma(x)^c \cap B(I,\rho)^c \cap \{y: d(y,I) < d(x,I)/2\}} G(x,y,t) dy dt \right)^{1/2} dx.
 \end{aligned}$$

Clearly

$$L_{1,1} \leq C \int_{d(x,I) \geq 8\rho} d(x,I)^{-(n+1)} dx \left(\int_0^\rho t dt \right)^{1/2} \leq C.$$

For $L_{1,2}$, note that in the region of integration $d(x,y) > d(x,I) - d(y,I) \geq d(x,I)/2$. Therefore,

$$\begin{aligned}
 L_{1,2} &\leq C \int_{d(x,I) \geq 8\rho} d(x,I)^{-\lambda n/2} \left(\int_0^\rho t^{\lambda n - n + 1} \int_{d(y,I) > \rho} d(y,I)^{-2n-2} dy dt \right)^{1/2} dx \\
 &\leq C \rho^{-\lambda n/2 + n} \rho^{\lambda n/2 - n/2 + 1} \rho^{-1 - n/2} \leq C.
 \end{aligned}$$

This completes the estimate of L_1 .

It remains to estimate L_3 . We divide the domain $\{(y,t) : d(x,y) > t, t > \rho\}$ into two pieces

$$\begin{aligned}
 \Omega_{x,1} &= \{(y,t) : d(y,x) > t, t > \rho, d(y, B(I, \rho)) \geq 2\rho\} \\
 \Omega_{x,2} &= \{(y,t) : d(y,x) > t, t > \rho, d(y, B(I, \rho)) < 2\rho\}
 \end{aligned}$$

and the integral L_3 into two terms $L_{3,1}$ and $L_{3,2}$ accordingly. For the latter, we argue as in Theorem 3.1 that

$$\begin{aligned}
 L_{3,2} &\leq C \rho \int_{d(x,I) \geq 8\rho} d(x,I)^{-\lambda n/2} \left(\int_\rho^\infty t^{-2n-3-n} \int_{d(y,I) \leq 3\rho} t^{\lambda n} dt dt \right)^{1/2} dx + C \\
 &\leq C \rho^{-\lambda n/2 + n + 1} \rho^{\lambda n/2 - n/2 - n - 1} \rho^{n/2} \leq C.
 \end{aligned}$$

Only $L_{3,1}$ remains. Let $\Theta_{x,1} = \{(y,t) \in \Omega_{x,1} : d(y,I) > \rho, d(y,I) > d(x,I)/2\}$ and $\Theta_{x,2} = \{(y,t) \in \Omega_{x,1} : d(y,I) > \rho, d(y,I) < d(x,I)/2\}$. Again using Lemmas

3.4 and [2, 6.4] we get

$$L_{3,1} \leq C + C\rho \int_{B(I,8\rho)^c} \left(\int_{\Theta_{x,1}} \frac{d(y, I)}{t^2 + d(y, I)^2} G(x, y, t) dy dt \right)^{1/2} dx$$

$$+ C\rho \int_{B(I,8\rho)^c} \left(\int_{\Theta_{x,2}} \frac{d(y, I)}{t^2 + d(y, I)^2} G(x, y, t) dy dt \right)^{1/2} dx.$$

The first of these integral summands is bounded by

$$C\rho \int_{B(I,8\rho)^c} d(x, I)^{-n-1/2} \left(\int_{\rho}^{\infty} t^{\lambda n - n - 2} \int_{d(x,y)>t} d(x, y)^{-\lambda n} dy dt \right)^{1/2} dx$$

$$\leq C\rho^{1/2} \left(\int_{\rho}^{\infty} t^{-2} dt \right)^{1/2} \leq C$$

and the second one is bounded by

$$C\rho \int_{B(I,8\rho)^c} d(x, I)^{-\lambda n/2} \left(\int_{\rho}^{\infty} t^{\lambda n - 2n - 2} \int_{B(I,\rho)^c} d(y, I)^{n+1} dy dt \right)^{1/2} dx \leq C.$$

Therefore, $\|g_{\lambda}^*(a)\|_1 \leq C$ for any atom $a(x)$, completing the proof of Theorem 4.1.

By an argument in [6], it is easy to prove that $\|g_{\lambda}^*(f)\|_p \leq C \|f\|_p$ for $p \geq 2$ and $\lambda > 2/p$. Interpolation (Theorem E of [3]) then gives

THEOREM 4.2. *For $p > 1$ and $\lambda > 2/p$, $\|g_{\lambda}^*(f)\|_p \leq C \|f\|_p$.*

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