Influence of the Neutron Star 1E 161348–5055 in RCW 103 on the Surrounding Medium

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Abstract: We have carried out a study of the neutral hydrogen in the direction of the X-ray source 1E 161348–5055, a compact central object (CCO) located in the interior of the supernova remnant (SNR) RCW 103. The Hi 21 cm line observations were carried out using the Australia Telescope Compact Array, complemented with single dish data from the Parkes radio telescope to recover information at all spatial scales. We derive a distance to RCW 103 of 3.1 kpc, in agreement with previous distance measurements. We have also detected a small hole in the Hi emission which is positionally and kinematically coincident with the location of the CCO which confirms the association between the SNR and the CCO. This is the third case of a depression in Hi emission seemingly associated with CCOs in SNRs. The characteristic parameters of the holes such as their size, eccentricity and evacuated mass are similar in all three cases. We estimate the absorbing HI column density towards 1E 161348–5055 to be ~6 × 10^{21} cm^{-2}, a value compatible with a blackbody solution for the CCO X-ray emission. However, the implied temperature and luminosity are very high compared to most neutron stars. Moreover, the strong long-term variability in X-rays favours the hypothesis that 1E 161348–5055 is an accreting binary source rather than an isolated, cooling neutron star. An analysis of the continuum image obtained at 1.4 GHz from these observations shows no trace of a pulsar wind nebula around 1E 161348–5055, in spite of it being a young object.

Keywords: stars: neutron — supernova remnants — ISM: individual: RCW 103 — X-rays: individual: 1E 161348–5055 — spectral lines: neutral hydrogen.

1 Introduction

X-ray observations from the past few years have revealed the existence of a large variety of Galactic point-like sources, many of them identified with neutron stars but presenting very different observational properties. These unresolved sources, which appear either isolated or in the interior of supernova remnants, have no radio counterpart and very high X-ray to optical flux ratios. About half of these sources show X-ray pulsations with periods between 6 and 12 seconds (anomalous X-ray pulsars, AXPs) and can even present sporadic strong γ-ray emission (soft gamma-ray repeaters, SGRs). However, Geppert, Page, & Zannias (1999) suggest that CCOs are fast-spinning, weakly-magnetised sources. The X-ray emission from CCOs is generally explained as thermal radiation from cooling NSs (e.g. Zavlin, Trümper, & Pavlov 1999), with typical temperatures of a few times 10^8 K, as inferred from their thermal-like spectra.

In a recent Hi study towards the bilateral SNR G296.5+10.0, Giacani et al. (2000) found that the associated CCO, 1E 1207.4–5209, lies near the centre of a small Hi depression located at the same systemic velocity as the SNR. The authors propose that the depression is due to self-absorption of a cool Hi cloud just to the foreground of a hotter volume of gas surrounding the CCO, and heated by its X-ray flux. We have begun a systematic search for similar traces in the neutral gas around other CCOs. The observations towards the X-ray point source RX J0822–4300 in Puppis A revealed an Hi structure consisting of a nearly circular minimum coincident with the CCO plus...
two aligned lobe-like depressions that appear to emerge from the CCO (Reynoso et al. 2003). The two lobes appear to have been formed by a combination of the proper motion of the CCO and the ejection of a collimated outflow. In this paper, we present the results obtained for 1E 161348–5055, the CCO associated with the SNR RCW 103 (G332.4–0.4).

At radio wavelengths, RCW 103 appears as an almost complete, circular, 8" diameter shell (Caswell et al. 1980). High-resolution radio polarimetric data reveal an approximately east-west alignment of the magnetic field which could be helping to limit the expansion north-south (Dickey et al. 1996). The brightening of the rim in the northern and southern sides, that hints for an incipient bilateral barrel shape for this SNR, could be related to this frozen-in magnetic field. Based on Hα absorption measurements at 21 cm, Caswell et al. (1975) suggest a distance of 3.3 kpc. Optical filaments are seen toward the brighter regions of the radio shell (van den Bergh, Marscher, & Terzian 1973; Ruiz 1983), and observations at infrared wavelengths show evidence for interaction with a dense interstellar medium (ISM), particularly to the south (Oliva, Moorwood, & Danziger 1990; Burton & Sypromilo 1993; Oliva et al. 1999). An optical expansion study of RCW 103 (Carter, Dickey, & Bomans 1997) indicates that this SNR is probably about 2000 years old, based on an assumed distance of 3.3 kpc. However, optical extinction studies suggest distances around 6.5 kpc (Leibowitz & Danziger 1983; Ruiz 1983).

Soft X-ray emission from RCW 103 was first detected with the Einstein Observatory (Tuohy et al. 1979), indicating a close correlation with the non-thermal radio emission. Also, a faint, point X-ray source, 1E 161348–5055, was located near the centre of the SNR (Tuohy & Garmire 1980). The lack of optical or radio counterparts led Tuohy et al. (1983) to propose that this point source was a thermally radiating CCO. However, subsequent observations with the Einstein IPC and ROSAT failed to confirm the existence of this source (Becker et al. 1993). Finally, Gotthelf, Petre, & Hwang (1997) detected hard X-rays from the elusive CCO using ASCA. They found that the spectral characteristics were incompatible with a simple cooling NS model. Follow-up observations (Gotthelf, Petre, & Vasisht 1999) confirmed that 1E 161348–5055 manifests long-term variability, explaining the non-detections after its discovery. There are strong indications that 1E 161348–5055 is part of a binary system (see section 4.1).

In this paper, we present radio continuum at 1.38 GHz and Hα 21 cm observations carried out with the Australia Telescope Compact Array (ATCA) towards the SNR RCW 103.

2 Observations and Data Reduction

Interferometric observations were obtained with the ATCA during one session of 12 h with the 750A array (baselines from 76.5 to 735 m), on 2002 January 22, and one session of 12 h with the EW 367 array (baselines from 46 to 367 m) on 2002 April 1. The antennas were pointed at RA = 16h41m30s, Dec = +51°0’0” (J2000).

A correlator configuration of 1024 channels, covering a depth of the intervening Hα gas, and T_e is the background continuum emission. Since we have subtracted the radio continuum from our Hα data, then Equation (1) can be reduced to

\[ T_e = (T_e - T_b) (1 - e^{-vT}) \]

where T_e and T_b are the spin temperature and the optical depth of the intervening Hα gas, and T_b is the background continuum emission. Since we have subtracted the radio continuum from our Hα data, then Equation (1) can be simplified to a single unknown of T_e = T_e. This condition can be accomplished by using only interferometric data, since in such way the extended Hα emission is filtered out.
We then computed \( e^{-\tau_v} \) over the same area, computed as described above [Equation (2)]. Strong absorption peaks appear at \(-43, -17, +1, \) and \(+34 \text{ km s}^{-1}\). All but the last feature were detected by Caswell et al. (1975) based on data obtained with the Parkes interferometer with a 2 km s\(^{-1}\) spectral resolution. We do not believe the absorption at \(+34 \text{ km s}^{-1}\) to be real because (a) absorption features should appear also from \(-43 \text{ km s}^{-1}\) to the tangent point at \(-115 \text{ km s}^{-1}\), and (b) the distance to the SNR would be approximately 20 kpc, which gives an unrealistic size and expansion velocity.

For comparison, we obtained ATCA spectra towards the strong Galactic radio sources present in the field G332.7–0.6 and G332.2–0.4 (Fig. 1), placed within 15' of RCW 103, and found that absorption features appear only in the velocity ranges from 0 to \(-50\) and 0 to \(-58 \text{ km s}^{-1}\) respectively.

To solve this puzzle, we followed the method employed by Dickey et al. (2003). We constructed a new Hi cube without subtracting the radio continuum and removing all baselines shorter than 1 k\(\lambda\) (210 m). In this way, all Hi structures larger than \(4'\) are filtered out. The term involving \( T_r \) in Equation (1) can thus be neglected. The data cube was further weighted by the radio continuum. We then computed \( e^{-\tau_v} \) using Equation (1). The resulting profile is plotted in Fig. 2c.

This filtering method succeeded in removing most of the Hi emission, as is apparent when comparing Figs. 2b and 2c. The absorption feature at \(-43 \text{ km s}^{-1}\) is confirmed, while the feature at \(+34 \text{ km s}^{-1}\) has disappeared. This implies that the latter was produced by self absorption on a scale between \(4'\) and \(-30'\).

The absorption feature at \(-43 \text{ km s}^{-1}\) corresponds to a lower distance limit of 3.1 kpc according to the Galactic rotation model of Fich, Blitz, & Stark (1989). We do not see any absorption against the emission peak at \(\sim-75 \text{ km s}^{-1}\), even though it has a brightness temperature of \(\sim 80 \text{ K}\). Thus the upper distance limit is 4.6 kpc. The lower distance coincides with the near side of the Scutum-Crux arm (Georgelin & Georgelin 1976), and will be assumed throughout the paper.

### 3.2 Hi Associated with RCW 103

We estimated the Hi column density towards 1E 161348–5055 by integrating the foreground Hi brightness temperature of the interferometric plus single-dish data up to \(-43 \text{ km s}^{-1}\). The value obtained depends on the lower limit adopted for the velocity interval in the integration since, although the local gas is supposed to lie at a systemic velocity of \(0 \text{ km s}^{-1}\), turbulence may cause departures of about \(7 \text{ km s}^{-1}\) from this value. Besides, due to the distance ambiguity, it is possible that some background gas is included in the integration, leading to an overestimation of \(N_{\text{Hi}}\). At \(v < -10 \text{ km s}^{-1}\), the line of sight crosses the far side of the Scutum-Crux arm. Taking into account these considerations, we integrate the brightness temperature between \(-3.5 \text{ and } -43 \text{ km s}^{-1}\) and estimate the Hi column density to be \(N_{\text{Hi}} \lesssim 6 \times 10^{21} \text{ cm}^{-2}\). The implications of this determination will be discussed in section 4.
The beam, 50′′ × 50′′, is plotted as a white open circle in the bottom left corner. The brightness-temperature scale is indicated on top of the image in units of K, while the contour levels vary from 68 K to 53 K. For comparison, a few representative contours of the radio continuum emission are included as white lines. The cross indicates the position of 1E 161348–5055 as given by Garmire et al. (1997). The interferometric data give \( T_{\text{c}} = T_{\text{R}} A_{\text{c}} \) (Schwarz et al. 1995). We averaged the Parkes data in a box around RCW 103 (excluding the contribution from the SNR) and found that \( T_{\text{c}} \) ≤ 100 K, hence \( T_{\text{R}} \) ≤ 20 K. Finally, at the location of 1E 161348–5055, the interferometric data give \( T_{\text{R}} = 44 \text{K} \) and \( T_{\text{c}} = -30 \text{K} \). Thus, applying Equation (3) to the emission towards the CCO, the hot Hi gas in the hole is found to have a temperature of \( < 130 \text{K} \). Such a temperature is unrealistically low at the interior of a SNR, therefore we discard the possibility that the Hi minima observed around CCOs are produced by self-absorption.

Figure 3. A grey-scale and contour image of the average H\(_i\) emission between -46.1 and -33.5 km s\(^{-1}\) towards RCW 103. The rms is 2 K. The beam, 50′′ × 50′′, is plotted as a white open circle in the bottom left corner. The brightness-temperature scale is indicated on top of the image in units of K, while the contour levels vary from 68 K to 53 K. For comparison, a few representative contours of the radio continuum emission are included as white lines. The cross indicates the position of 1E 161348–5055 as given by Garmire et al. (2000a).

An inspection of the Hi images at velocities around -43 km s\(^{-1}\) revealed that 1E 161348–5055 lies inside a local Hi depression which is present in all channels between -46.1 and -33.5 km s\(^{-1}\). In Figure 3, an image of the average Hi emission within this velocity interval is shown. The Hi depression, which attains its minimum at RA = 10^\(h\)17^\(m\)32^\(s\), Dec = -51°24′′ (J2000), is elongated, with a minor-to-major axis ratio of ~0.7, and has a mean diameter of ~64′′ (1 pc at a distance of 3.1 kpc). 1E 161348–5055 is 20″ (0.3 pc) away from the centre of the hole. We note that due to the proximity to the Galactic Plane, our data show several similar Hi depressions at different locations and velocities. However, the coincidence in position and velocity makes an association between the CCO and this Hi feature very likely.

In what follows, we analyse possible origins for this Hi void. If the Hi minimum is produced because of a real absence of Hi, then the missing mass is estimated to be 0.3 M\(_{\odot}\). If instead it is produced by hot Hi gas self-absorbed by a cooler foreground, as proposed for the CCO 1E 1207.4–5209 (Giacani et al. 2000), then we can follow the method described by Schwarz et al. (1995) to estimate an upper limit for the temperature of the hot neutral hydrogen gas.

If self-absorption is considered, then Equation (1) can be re-written as (cf. Equation (3) in Schwarz et al. 1995)

\[
T_{\text{c}} = (T_{\text{R}} - T_{\text{bg}}) A_{\text{c}} - T_{\text{c}} A_{\text{c}}, \tag{3}
\]

with \( A_{\text{c}} = (1 - e^{-\tau_{\alpha}}) \), and where \( T_{\text{bg}} \) is the temperature of the background hot Hi. Assuming that \( A_{\text{c}} \) is uniform across the continuum source, then \( T_{\text{c}} \) is a linear function of \( T_{\text{R}} \) with slope \((-A_{\text{c}})\) and a zero offset \((T_{\text{R}} - T_{\text{bg}})A_{\text{c}}\). The Hi hole attains its minimum emission at \( v = -46.1 \text{km s}^{-1}\).

We compared the brightness temperature of the ATCA Hi data at this velocity with the continuum emission and fitted a straight line to the distribution. To avoid confusion possibly introduced by regions of low continuum emission, all values of \( T_{\text{c}} \) under 350 mJy beam\(^{-1}\) were clipped. We determined that \( A_{\text{c}} = 0.195 \), which implies that \( \tau_{\alpha} = 0.2 \). To estimate the spin temperature, it must be noted that for a single dish measurement towards a region with no continuum emission, the brightness temperature of the line is \( T_{\text{c}} = T_{\text{R}} A_{\text{c}} \) (Schwarz et al. 1995). We averaged the Parkes data in a box around RCW 103 (excluding the contribution from the SNR) and found that \( T_{\text{c}} \) ≤ 100 K, hence \( T_{\text{R}} \) ≤ 20 K. Finally, at the location of 1E 161348–5055, the interferometric data give \( T_{\text{R}} = 44 \text{K} \) and \( T_{\text{c}} = -30 \text{K} \). Thus, applying Equation (3) to the emission towards the CCO, the hot Hi gas in the hole is found to have a temperature of \( < 130 \text{K} \). Such a temperature is unrealistically low at the interior of a SNR, therefore we discard the possibility that the Hi minima observed around CCOs are produced by self-absorption.

3.3 Radio Continuum Emission

The rotational energy of pulsars is dissipated via a magnetised relativistic wind composed of electrons and positrons. The shock front between this wind and the ambient medium can give rise to a synchrotron emitting bubble (known as a pulsar wind nebula, PWN). The detectability of such a PWN depends on the density of the ambient medium and the pulsar parameters. The majority of pulsars do not have detectable PWNs (Gaensler et al. 2000). For a very young pulsar inside a SNR, the luminosity \( L_{\text{c}} \) of the PWN is a strong function of the initial period of the pulsar (Gotthelf et al. 1997). We investigated different fits to the X-ray spectrum of 1E 161348–5055 using a nonequilibrium ionisation plasma model combined with a blackbody, a power-law, and a thermal bremsstrahlung.

4 Discussion

4.1 Spectral Model

Gotthelf et al. (1997) investigated different fits to the X-ray spectrum of 1E 161348–5055 using a nonequilibrium ionisation plasma model combined with a blackbody, a power-law, and a thermal bremsstrahlung.
of the image in units of Jy beam$^{-1}$ at continuum frequency of 1384 MHz. The gray scale is shown on top of the image in units of Jy beam$^{-1}$. The contours are plotted in steps of 7.5% of the peak intensity, 768 mJy beam$^{-1}$, starting at 15%. For clarity, white lines are used over regions of darker background. The cross shows the position of 1E 161348–5055 as given by Garmire et al. (2000a). The noise level is 5.5 mJy beam$^{-1}$.

In Table 1 we compare the main parameters derived for the present study reveals the third case in which a CCO associated with a SNR is located inside an H$\text{ii}$ depression. The other two cases are 1E 1207.4–5209 in G296.5+10.0, and RX J0822–4300 in Puppis A. In Table 1 we compare the main parameters derived for 1E 161348–5055 with those of other CCOs.

<table>
<thead>
<tr>
<th>CCO</th>
<th>1E 1207.4–5209$^a$ (G296.5+10.0)</th>
<th>RX J0822–4300$^b$ (Puppis A)</th>
<th>1E 161348–5055$^c$ (RCW 103)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Mean angular diameter (λ)</td>
<td>5.3</td>
<td>2.3</td>
<td>1.1</td>
</tr>
<tr>
<td>Mean linear diameter (pc)</td>
<td>3.2</td>
<td>1.3</td>
<td>1.0</td>
</tr>
<tr>
<td>Minor/major axis ratio</td>
<td>0.8</td>
<td>0.8</td>
<td>0.7</td>
</tr>
<tr>
<td>CCO offset from centre (λ)</td>
<td>30</td>
<td>37</td>
<td>20</td>
</tr>
<tr>
<td>CCO offset from centre (pc)</td>
<td>0.3</td>
<td>0.4</td>
<td>0.3</td>
</tr>
<tr>
<td>Missing mass (M$\odot$)</td>
<td>-</td>
<td>0.1</td>
<td>0.3</td>
</tr>
</tbody>
</table>

the three H\textsc{i} depressions. For G296.5+10.0, only the size and velocity width are given in Giacani et al. (2000) and thus the remaining parameters, where possible, have been estimated directly from the images presented in their paper. In computing the minor to major axis ratio, the beam elongation was taken into account.

A number of similarities can be found between the three cases. All H\textsc{i} cavities appear to have the same elongation, and the CCOs are off-centred by similar distances. In the two cases where the missing mass was computed, a similar value was obtained. On the other hand, the size of the H\textsc{i} hole in G296.5+10.0 is larger than the other two holes by more than a factor of two. This discrepancy may be related to the beam size of the observations, which is almost four times larger than those of the Puppis A and RCW 103 data. The main difference between the three cases is given by Puppis A, which presents two lobe-like H\textsc{i} depressions emerging from its associated CCO at the systemic velocity of the SNR (Reynoso et al. 2003). The present observations do not reveal any similar morphology around 1E 161348–5055.

In what follows we investigate whether the H\textsc{i} depression around 1E 161348–5055 can be a swept up hole. The rate of energy production needed to sweep up a cavity into motion with velocity $V = r/t$ is

$$E_k = 3 \times 10^{37} \text{ erg s}^{-1},$$

where $M$ is the evacuated mass in units of $M_\odot$, $r$ is the radius of the hole in pc, and $t$ is the age of the CCO, equal to the age of the host SNR, in years. Replacing in Equation (4) the radius and missing mass computed in Section 3, and assuming that the age of RCW 103 is 2000 yr (Carter et al. 1997), we obtain $E_k = 3.2 \times 10^{37} \text{ erg s}^{-1}$, where the quoted errors allow for ages of 1000 and 3000 yrs old. The spin down energy loss has not been measured for 1E 161348–5055, but if it is similar to the values observed in other CCOs (energy loss rates between $\sim 1 \times 10^{36} \text{ erg s}^{-1}$ and $1.5 \times 10^{37} \text{ erg s}^{-1}$; Slane et al. 1997; Brazier & Johnston 1999), then a complete conversion of the rotational energy into kinetic energy of the surrounding medium can account for the observed H\textsc{i} hole.

An alternative possibility is that the depression does not contain swept up low density gas, but is filled with H\textsc{i} gas heated up by the CCO at temperatures higher than the surroundings. In that case, the depression should contain enhanced ionised hydrogen (i.e. it would form a small H\textsc{ii} region), and this is not observed in the radiocontinuum image. We would, however, expect to see infra-red emission from such a region but a search in mid- and near-infrared wavelengths using data from the Midcourse Space Experiment (MSX) and the Two Micron All Sky Survey (2MASS) yielded negative results. In conclusion, we find that the swept-up cavity provides a more convincing explanation for the observed H\textsc{i} minimum.

5 Conclusions

In this paper, we present the third case in which a CCO lies at a local H\textsc{i} minimum at a velocity compatible with the systemic velocity of the host SNR. We have shown that self-absorption does not provide a satisfactory explanation for this kind of feature, as was proposed for 1E 1207.4–5209 (Giacani et al. 2000). Instead, it is possible that the H\textsc{i} depression is a swept up hole, where $\sim 0.3 M_\odot$ of neutral gas has been evacuated. We have found a number of similarities between the three H\textsc{i} holes around CCOs detected so far, such as the elongation, the missing mass, and the off-centred position of the CCOs. We did not detect any synchrotron nebula around the CCO down to a level of 1 mJy beam$^{-1}$.

The present data allow us to constrain the H\textsc{i} column density to be $N_{\text{HI}} = 6 \times 10^{21} \text{ cm}^{-2}$. This column density favours the blackbody model of Gotthelf et al. (1999), however the derived temperature is too high to be explained by standard NS cooling models. Instead, it appears most likely that 1E 161348–5055 is an accreting binary, and may constitute the first case in which such a system occurs in the interior of a SNR.

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