# On Hyperbolicity of Domains with Strictly Pseudoconvex Ends 

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#### Abstract

This article establishes a sufficient condition for Kobayashi hyperbolicity of unbounded domains in terms of curvature. Specifically, when $\Omega \subset \mathbb{C}^{n}$ corresponds to a sub-level set of a smooth, real-valued function $\Psi$, such that the form $\omega=\mathbf{i} \partial \bar{\partial} \Psi$ is Kähler and has bounded curvature outside a bounded subset, then this domain admits a hermitian metric of strictly negative holomorphic sectional curvature.


## 1 Introduction

It is well known that any domain in $\mathbb{C}^{n}$ biholomorphically equivalent to a bounded domain is Kobayashi-hyperbolic. The main result of this note, proven in Section 2, provides a sufficient condition for hyperbolicity of unbounded domains in terms of curvature. In general, a complex space is Kobayashi-hyperbolic if it can be shown to possess a hermitian metric, the holomorphic sectional curvature of which is bounded by a negative constant (cf. [3]). With this in mind, we assume that

$$
\Omega=\left\{Z \in \mathbb{C}^{n} \mid \Psi(Z)<1\right\}
$$

for a smooth function $\Psi: \Omega^{\prime} \rightarrow[0,+\infty)\left(\bar{\Omega} \subset \Omega^{\prime}\right)$ that is strongly plurisubharmonic outside a bounded subset of $\Omega$.

Theorem 1.1 Suppose there exists a bounded subset of $\Omega$ outside which the real form $\omega:=\mathbf{i} \partial \bar{\partial} \Psi$ has bounded curvature. Then $\Omega$ is Kobayashi-hyperbolic.

The proof of this result is not significantly altered if the role of $\mathbb{C}^{n}$ (i.e., as the ambient space containing $\Omega$ ) is taken by an arbitrary Stein manifold. Section 3 provides an example of a weakly pseudoconvex unbounded domain satisfying the above hypotheses, which we introduce as follows. Note first that the orthogonal group acts holomorphically on vectors $Z=X+\mathbf{i} Y \in \mathbb{C}^{n}$ according to the natural rule

$$
\sigma(Z)=\sigma(X)+\mathbf{i} \sigma(Y) \text { for all } \sigma \in(\mathbb{O})(n)
$$

Beginning with a projection map

$$
\pi:\left(\mathbb{C}^{n} \rightarrow[0,+\infty) \times[0,+\infty)\right.
$$

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such that $\pi(Z)=(|\Re(Z)|,|\Im(Z)|)=(|X|,|Y|)$, where $|\cdot|$ denotes the Euclidean norm in $\mathbb{R}^{n}$, we compose with any smooth function $\psi:[0,+\infty) \times[0,+\infty) \rightarrow$ $[0,+\infty)$ to define domains of the form

$$
\Omega=\left\{Z \in \mathbb{C}^{n} \mid \Psi(Z)<c\right\}
$$

where $\Psi=\psi \circ \pi$, for some $c>0$. These domains clearly possess orthogonal symmetry, though in general they are not symmetric with respect to the full unitary group. They are introduced in [2] as the building blocks of a theory of cellular decomposition of Stein manifolds, based on a famous theorem of Andreotti and Frankel [1]. Our calculation of holomorphic sectional curvature will be carried out specifically for $\psi(r, s)=r^{2} s^{2}$ and when $n=2$. In this case $\Psi$ is seen to be strongly plurisubharmonic (hence $\omega$ is Kähler) outside a pair of transversely intersecting discs, which are extremally embedded with respect to the Kobayashi metric and are bounded by the unique $(0)(2)$-orbit of weakly pseudoconvex points on the boundary of $\Omega$.

## 2 Hyperbolicity and Pseudoconvex Ends

Let $\Psi: \mathbb{C}^{n} \rightarrow[0,+\infty)$ be a smooth function and let $B \subset \mathbb{C}^{n}$ be a bounded subset. Consider a domain $\Omega \subseteq \mathbb{C}^{n}$ defined, without significant loss of generality, by the inequality $\Psi<1$ and contained in a slightly larger domain $\Omega^{\prime}$ corresponding to $\Psi<1+\varepsilon$ (in particular, $\bar{\Omega} \subset \Omega^{\prime}$ ). If it is assumed that $\left.\Psi\right|_{\Omega^{\prime} \backslash B}$ is strongly plurisubharmonic, then the real closed form $\omega:=\mathbf{i} \partial \bar{\partial} \Psi$ is Kähler on $\Omega^{\prime} \backslash B$, and we may denote by $g$ the associated metric. As always, the curvature tensor of $g$ is defined by the formula

$$
R_{i, \bar{j}, k, \bar{l}}=-\frac{\partial^{2} g_{i, \bar{j}}}{\partial z_{k} \partial \bar{z}_{l}}+\sum_{\alpha, \beta} g^{\alpha, \bar{\beta}} \frac{\partial g_{i, \bar{\beta}}}{\partial z_{k}} \frac{\partial g_{\alpha, \bar{j}}}{\partial \bar{z}_{l}}
$$

away from $B$. In the following theorem, all closures are taken with respect to $\Omega^{\prime}$.
Theorem 2.1 Suppose $\left.\omega\right|_{\bar{\Omega} \backslash B}$ has bounded curvature. Then $\Omega$ is Kobayashi-hyperbolic.

Proof We begin by defining a metric $h_{K}$ associated with the form $e^{K \Psi} \partial \bar{\partial} \Psi$ on $\Omega \backslash B$. Given a holomorphic map $f(\zeta)$ from the unit disc $D$ into $\Omega$, such that $f(0)=Z \in$ $\Omega \backslash B$, we assume the presence of Kähler normal coordinates $W=\left(w_{1}, \ldots, w_{n}\right)$ in a neighbourhood of $Z(W=0)$ in which

$$
g_{i, \bar{j}}(0)=\delta_{i, j} \quad \text { and } \quad \frac{\left(\partial g_{i, \bar{j}}\right)}{\left(\partial w_{k}\right)}(0)=0, \quad 1 \leq k \leq n
$$

Now consider

$$
f^{*} h_{K}=\left(e^{K \Psi \circ f} \sum_{i, j} \frac{\partial^{2} \Psi}{\partial w_{i} \partial \bar{w}_{j}} f_{i}^{\prime} \bar{f}_{j}^{\prime}\right) d \zeta \wedge d \bar{\zeta}=\mu_{f} d \zeta \wedge d \bar{\zeta}
$$

(for $K$ an as yet unspecified constant) on $T_{\zeta}^{1,0} D \otimes T_{\zeta}^{0,1} D$. Now

$$
\log \left(\mu_{f}\right)=K \Psi \circ f+\log \left(\sum_{i, j} g_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}\right)
$$

implies

$$
\frac{\partial \log \left(\mu_{f}\right)}{\partial \bar{\zeta}}=K \frac{\partial \Psi \circ f}{\partial \bar{\zeta}}+\frac{1}{\sum_{i, j} g_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}} \sum_{l=1}^{n}\left(\frac{\partial g_{i, \bar{j}}}{\partial \bar{w}_{l}} \bar{f}_{l}^{\prime} f_{i}^{\prime} \bar{f}_{j}^{\prime}+\delta_{j, l} g_{i, \bar{j}} f_{i}^{\prime} \bar{f}_{j}^{\prime \prime}\right),
$$

and hence

$$
\begin{aligned}
& \frac{\partial^{2} \log \left(\mu_{f}\right)}{\partial \zeta \partial \bar{\zeta}} \\
& \quad= \\
& \quad K \frac{\partial^{2} \Psi \circ f}{\partial \zeta \partial \bar{\zeta}} \\
& \quad-\frac{1}{\left(\sum_{i, j} g_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}\right)^{2}} \sum_{k, l}\left(\frac{\partial g_{i, \bar{j}}}{\partial w_{k}} f_{k}^{\prime} f_{i}^{\prime} \bar{f}_{j}^{\prime}+\delta_{i, k} g_{i, j} f_{i}^{\prime \prime} \bar{f}_{j}^{\prime}\right)\left(\frac{\partial g_{i, j}}{\partial \bar{w}_{l}} \bar{f}_{l}^{\prime} f_{i}^{\prime} \bar{f}_{j}^{\prime}+\delta_{j, l} g_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime \prime}\right) \\
& \\
& \quad+\frac{1}{\sum_{i, j} g_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}} \sum_{i, j, k, l}\left(\frac{\partial^{2} g_{i, j}}{\partial w_{k} \partial \bar{w}_{l}} f_{k}^{\prime} \bar{f}_{l}^{\prime} f_{i}^{\prime} \bar{f}_{j}^{\prime}+g_{i, j} \delta_{j, l} \delta_{i, k} f_{i}^{\prime \prime} \bar{f}_{j}^{\prime \prime}\right) \\
& = \\
& =K\left|f^{\prime}\right|^{2}-\frac{1}{\left|f^{\prime}\right|^{4}} \sum_{k, l} f_{k}^{\prime \prime} \bar{f}_{k}^{\prime} f_{l}^{\prime} \bar{f}_{l}^{\prime \prime}+\frac{1}{\left|f^{\prime}\right|^{2}} \sum_{i, j, k, l}\left(\frac{\partial^{2} g_{i, j}}{\partial w_{k} \partial \bar{w}_{l}} f_{k}^{\prime} \bar{f}_{l}^{\prime} f_{i}^{\prime} \bar{f}_{j}^{\prime}+\left|f^{\prime \prime}\right|^{2}\right) \\
& \quad= \\
& \quad K\left|f^{\prime}\right|^{2}-\frac{\left|\left\langle f^{\prime \prime}, f^{\prime}\right\rangle\right|^{2}}{\left|f^{\prime}\right|^{4}}+\frac{\left|f^{\prime \prime}\right|^{2}}{\left|f^{\prime}\right|^{2}}-\frac{1}{\left|f^{\prime 2}\right|^{2}} \sum_{i, j, k, l} R_{i, \bar{j}, k, i} f_{k}^{\prime} \bar{f}_{l}^{\prime} f_{i}^{\prime} \bar{f}_{j}^{\prime} \\
&
\end{aligned}
$$

where $R_{i, \bar{j}, k, \bar{l}}$ as above denotes the curvature tensor of $\omega$. Since $R_{i, \bar{j}, k, \bar{l}}$ is uniformly bounded on $\Omega \backslash B$, we may choose the constant $K \gg \sup _{\Omega \backslash B}\|R\|$. To complete the estimation of holomorphic sectional curvature at $Z \in \Omega \backslash B$, it remains to note $\mu_{f}=e^{K \Psi \circ f}\left|f^{\prime}\right|^{2}$, and therefore

$$
-\frac{1}{\mu_{f}} \frac{\partial^{2} \log \left(\mu_{f}\right)}{\partial \zeta \partial \bar{\zeta}} \leq-\left(K-\sup _{\Omega \backslash B}\|R\|\right) e^{-K}
$$

given $\Psi(Z)<1$. Without loss of generality, let $B$ correspond to

$$
\bar{\Omega} \cap B_{M}(0)=\{Z \in \bar{\Omega}| | Z \mid<M\} .
$$

Consider also the slightly larger ball $B_{M+\varepsilon}(0)$. We now introduce a $C^{\infty}$ cut-off function $\chi$, such that

$$
\chi(|Z|)= \begin{cases}1 & \text { if } Z \in \overline{B_{M}(0)}, \\ 0 & \text { if } Z \in \mathbb{C}^{n} \backslash B_{M+\varepsilon}(0)\end{cases}
$$

Letting $\omega_{0}$ again denote the standard Kähler metric form on $\mathbb{C}^{n}$, we then define a hermitian metric $h$ on $\Omega$ associated with the form

$$
e^{K^{\prime} \Psi+\chi \cdot|Z|^{2}}\left((1-\chi) \omega+\chi \cdot \omega_{0}\right)
$$

(In particular, the adjusted constant $K^{\prime} \geq K$ will be defined below.) The upper bound on holomorphic sectional curvature remains essentially the same for $Z \in$ $\Omega \backslash B_{M+\varepsilon}(0)$, where the metric is identified with $\overline{h_{K^{\prime}}}$ above. Two further regions of $\Omega$ must now be examined. First, note that if $Z \in \overline{B_{M}(0)} \cap \Omega$, we may consider $h$ to be associated with the form $e^{K^{\prime} \Psi+|Z|^{2}} \omega_{0}$. Hence to $f^{*} h$ we associate the function

$$
\log \left(\mu_{f}\right)=K^{\prime} \Psi \circ f+|f|^{2}+\log \left(\left|f^{\prime}\right|^{2}\right)
$$

and obtain

$$
\begin{aligned}
-\frac{1}{\mu_{f}} \frac{\partial^{2} \log \left(\mu_{f}\right)}{\partial \zeta \partial \bar{\zeta}} & \leq-\frac{e^{-\left(K^{\prime} \Psi \circ f+|f|^{2}\right)}}{\left|f^{\prime}\right|^{2}}\left(K^{\prime} g_{i, \bar{j}}+\delta_{i, j}\right) f_{i}^{\prime} \bar{f}^{\prime}{ }_{j} \\
& \leq-e^{-\left(K^{\prime}+M^{2}\right)}
\end{aligned}
$$

if we recall that $g_{i, j} f_{i}^{\prime} \bar{f}^{\prime}{ }_{j} \geq 0$. It remains now to estimate the holomorphic sectional curvature in the region

$$
A_{\varepsilon}=\{Z \in \Omega|M<|Z|<M+\varepsilon\}
$$

For $Z \in A_{\varepsilon}$ we will simply write

$$
\mu_{f}=e^{K^{\prime} \Psi \circ f} \Sigma_{i, j}(G \circ f)_{i, j} f_{i}^{\prime} \bar{f}^{\prime}{ }_{j}
$$

where $G_{i, \bar{j}}=\left((1-\chi) g_{i, \bar{j}}+\chi \delta_{i, j}\right)$ is understood to be smooth, positive definite, and bounded, with bounded derivatives in the region $A_{\varepsilon}$, which is relatively compact in $\Omega^{\prime}$. The calculation of holomorphic sectional curvature is then formally carried out as in the case of $h_{K}$, producing a curvature tensor $R^{\prime}$ that is bounded on $\Omega \backslash B_{M}(0)$. The leading term of this calculation corresponds to

$$
-\frac{K^{\prime}}{\mu_{f}} \frac{\partial^{2} \Psi}{\partial \bar{z}_{j} \partial z_{i}} f_{i}^{\prime} \bar{f}_{j}^{\prime}=-K^{\prime} e^{-\left(K^{\prime} \Psi \circ f\right)} \frac{g_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}}{(G \circ f)_{i, \bar{j}} f_{i}^{\prime} \bar{f}_{j}^{\prime}} \leq-K^{\prime} e^{-K^{\prime}} \inf _{A_{\varepsilon}} \frac{g_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}}{(G \circ f)_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}}
$$

To see that the infimum above is strictly positive, we write

$$
\frac{g_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}}{(G \circ f)_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}}=\frac{e^{-\chi \cdot|Z|^{2}} g_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}}{\left((1-\chi) g_{i, j}+\chi \delta_{i, j}\right) f_{i}^{\prime} \bar{f}_{j}^{\prime}}=\frac{e^{-\chi \cdot|Z|^{2}} g\left(\frac{f^{\prime}}{\left|f^{\prime}\right|}, \frac{f^{\prime}}{\left|f^{\prime}\right|}\right)}{(1-\chi) g\left(\frac{f^{\prime}}{\left|f^{\prime}\right|}, \frac{f^{\prime}}{\left|f^{\prime}\right|}\right)+\chi}
$$

noting that $g(\mathbf{v}, \mathbf{v})$ is positive definite for all $Z \in \overline{A_{\varepsilon}}$, and all $\mathbf{v} \in \mathbb{S}^{2 n-1}$. It remains now to choose

$$
K^{\prime} \cdot \inf _{A_{\varepsilon}} \frac{g_{i, \bar{j}} f_{i}^{\prime} \bar{f}_{j}^{\prime}}{(G \circ f)_{i, j} f_{i}^{\prime} \bar{f}_{j}^{\prime}} \gg \sup _{\Omega \backslash B_{M}(0)}\left\|R^{\prime}\right\|,
$$

as in the construction of $h_{K}$, so that the holomorphic sectional curvature of $h$ is uniformly bounded above by a strictly negative constant on $\Omega$. It follows at once that $\Omega$ is Kobayashi-hyperbolic.

3 An Example: $\Psi(Z)=|\Re(Z)|^{2}|\Im(Z)|^{2}$
Let $Z=\left(z_{1}, z_{2}, \ldots, z_{n}\right)$ be coordinates in $\mathbb{C}^{n}(n \geq 2), z_{j}=x_{j}+\mathbf{i} y_{j}$, and let $\Re(Z)=$ $X=\left(x_{1}, x_{2}, \ldots, x_{n}\right), \Im(Z)=Y=\left(y_{1}, y_{2}, \ldots, y_{n}\right)$. Then $\pi: \mathbb{C}^{n} \rightarrow \mathbb{R}^{2}$ will denote the natural projection

$$
\pi(Z)=(|X|,|Y|),
$$

where $|\cdot|$ denotes the Euclidean norm. Given $\psi:[0,+\infty) \times[0,+\infty) \rightarrow[0,+\infty)$ smooth, we define special complex domains with $(\mathbb{O})(n)$-symmetry, of the form

$$
\Omega=\left\{Z \in \mathbb{C}^{n} \mid \Psi(Z)<c\right\} ; \quad \Psi=\psi \circ \pi .
$$

A sufficient condition for local plurisubharmonicity of a general function $\Psi=\psi \circ \pi$ is examined in [2]. If it is assumed that the critical locus of $\psi$ is contained in $\psi^{-1}(0)$, then $d \psi \neq \mathbf{0}$ at all points $(r, s)$ outside that locus, and the level set $\psi(r, s)=c$ passing through a given regular point locally admits an implicit function $s=\phi(r)$. Let $r=$ $|\Re(Z)|, s=|\Im(Z)|$, and $Z \in \Psi^{-1}(c), 0<c \leq 1$. The following statement is given in [2, Lemma 3.3.1].

Lemma ([2, Lemma 3.3.1]) If, in a neighbourhood of $\left(r_{0}, s_{0}\right)=(|\Re(Z)|,|\Im(Z)|)$, $Z \in \Psi^{-1}(c)$, we have
(a) $\phi^{\prime}(r)>0, \phi^{\prime \prime}(r) \leq 0, \phi^{\prime}(r)>\frac{r}{\phi(r)}$, and

$$
\phi^{\prime \prime}(r)+\frac{\left(\phi^{\prime}(r)\right)^{3}}{r}-\frac{1}{\phi(r)}\left(1+\left(\phi^{\prime}(r)\right)^{2}\right) \geq 0,
$$

then the hypersurface is pseudoconvex at $Z$, co-oriented from above;
(b) $\phi^{\prime}(r) \leq 0, \phi^{\prime \prime}(r) \geq 0$, and

$$
\phi^{\prime \prime}(r)+\frac{\left(\phi^{\prime}(r)\right)^{3}}{r}-\frac{1}{\phi(r)} \leq 0,
$$

then the hypersurface is pseudoconvex at $Z$, co-oriented from below.
In the following, we examine the case $\psi(r, s)=r^{2} s^{2}$, for which the pseudoconvexity condition above is easily checked to hold. More directly, we can perform some routine calculations based on the formula

$$
\Psi(Z)=|\Re(Z)|^{2}|\Im(Z)|^{2}=\frac{1}{16}\left(4|Z|^{4}-\left[\sum_{k=1}^{n} z_{k}^{2}+\bar{z}_{k}^{2}\right]^{2}\right) .
$$

In particular,

$$
\begin{aligned}
\frac{\partial^{2} \Psi}{\partial \bar{z}_{j} \partial z_{i}} & =\frac{1}{16} \frac{\partial}{\partial \bar{z}_{j}}\left(8|Z|^{2} \cdot \bar{z}_{i}-4 \sum_{k}\left(z_{k}^{2}+\bar{z}_{k}^{2}\right) \cdot z_{i}\right) \\
& =\frac{1}{2}\left(|Z|^{2} \delta_{i, j}+z_{j} \bar{z}_{i}-z_{i} \bar{z}_{j}\right) .
\end{aligned}
$$

Given any tangent vector $\mathbf{v} \in T_{Z} \mathbb{C}^{n}$, note that

$$
\begin{aligned}
\frac{\partial^{2} \Psi}{\partial \bar{z}_{j} \partial z_{i}}(\mathbf{v}, \overline{\mathbf{v}}) & =\sum_{i, j} \frac{\partial^{2} \Psi}{\partial \bar{z}_{j} \partial z_{i}} \bar{v}_{j} \cdot v_{i}=\frac{1}{2} \sum_{i, j}\left(|Z|^{2} \delta_{i, j}+z_{j} \bar{z}_{i}-z_{i} \bar{z}_{j}\right) \bar{v}_{j} \cdot v_{i} \\
& =\frac{1}{2}\left(|Z|^{2}|\mathbf{v}|^{2}+|\langle Z, \mathbf{v}\rangle|^{2}-|\langle\bar{Z}, \mathbf{v}\rangle|^{2}\right) \\
& \geq \frac{1}{2}|\langle Z, \mathbf{v}\rangle|^{2} \geq 0
\end{aligned}
$$

where $\langle\cdot, \cdot\rangle$ denotes the standard hermitian inner product on $\mathbb{C}^{n}$. Hence $\Psi$ is plurisubharmonic. Note also that

$$
\left(\nabla^{1,0} \Psi\right)_{i}=\left(\overline{\nabla^{0,1} \Psi}\right)_{i}=\frac{1}{2}|Z|^{2} \cdot \bar{z}_{i}-\frac{1}{4} \sum_{k}\left(z_{k}^{2}+\bar{z}_{k}^{2}\right) \cdot z_{i}=0
$$

if and only if $\Psi(Z)=0,1 \leq i \leq n$, hence the critical locus of this plurisubharmonic function coincides with $\mathbb{R}^{n} \cup \mathbf{i} \mathbb{R}^{n}$. Let $\Gamma \subset \Omega$ denote the set of points at which the form $\omega=\mathbf{i} \partial \bar{\partial} \Psi$ is degenerate (i.e., $\omega$ restricted to $\Omega \backslash \Gamma$ is Kähler). Then $\omega$ induces the Levi form on restriction to the complex tangent space of the $\Psi$-level set through any point $Z \in \Omega$, hence weak pseudoconvexity of the level set at $Z$ implies $Z \in \Gamma$. Let $g_{i, j}$ denote the $i, j$-component of the associated Kähler metric corresponding to

$$
\frac{\partial^{2} \Psi}{\partial z_{i} \partial \bar{z}_{j}}=\frac{1}{2}\left(|Z|^{2} \delta_{i, j}+z_{j} \bar{z}_{i}-z_{i} \bar{z}_{j}\right)
$$

In the remainder we will specialize to the case $n=2$, where $\Gamma$ is defined simply by the vanishing of

$$
\operatorname{det}(g)=\frac{1}{4}\left(|Z|^{4}-\left|z_{1} \bar{z}_{2}-\bar{z}_{1} z_{2}\right|^{2}\right)=\frac{1}{4}\left|z_{1}^{2}+z_{2}^{2}\right|^{2}
$$

hence $\operatorname{det}(g)=0$ if and only if $z_{1}= \pm \mathbf{i} z_{2}$. Now $g_{i, j}=\frac{1}{2}\left(|Z|^{2} \delta_{i, j}+z_{j} \bar{z}_{i}-z_{i} \bar{z}_{j}\right)$, implies

$$
\frac{\partial^{2} g_{i, \bar{j}}}{\partial z_{k} \partial \bar{z}_{l}}=\frac{1}{2}\left(\delta_{k, l} \delta_{i, j}+\delta_{l, i} \delta_{k, j}-\delta_{k, i} \delta_{l, j}\right)
$$

Moreover,

$$
R_{i, \bar{j}, k, \bar{l}}=-\frac{\partial^{2} g_{i, \bar{j}}}{\partial z_{k} \partial \bar{z}_{l}}+\sum_{\alpha, \beta} g^{\alpha, \bar{\beta}} \frac{\partial g_{i, \bar{\beta}}}{\partial z_{k}} \frac{\partial g_{\alpha, \bar{j}}}{\partial \bar{z}_{l}}
$$

is bounded on $\Omega \backslash \Gamma_{\varepsilon}$, where $\Gamma_{\varepsilon}=\left\{Z| | z_{1}^{2}+z_{2}^{2} \mid<\varepsilon\right\}$ for $\varepsilon$ arbitrarily small and positive. A uniform bound can be determined reasonably explicitly in this case if we introduce a parameter $\xi=r e^{\mathrm{i} \theta}$ such that $z_{1}=\xi z_{2}$, and hence

$$
\operatorname{det}(g)=\frac{\left|z_{2}\right|^{4}}{4}\left(\left(r^{2}+1\right)^{2}-4 r^{2} \sin ^{2}(\theta)\right) \geq \frac{\left|z_{2}\right|^{4}}{4}\left(r^{2}-1\right)^{2}
$$

Then

$$
\begin{aligned}
\left|g^{\alpha, \bar{\beta}}\right| & =\frac{\left|z_{2}\right|^{2}}{2 \operatorname{det}(g)}\left(\left(r^{2}+1\right)^{2} \delta_{\alpha, \beta}+4\left(1-\delta_{\alpha, \beta}\right) r^{2} \sin ^{2}(\theta)\right)^{\frac{1}{2}} \\
& \leq \frac{2\left(r^{2}+1\right)}{\left|z_{2}\right|^{2}\left(r^{2}-1\right)^{2}}
\end{aligned}
$$

(if we note that $\left.\left(r^{2}+1\right)^{2} \geq 4 r^{2}\right)$. Similarly,

$$
\frac{\partial g_{i, \bar{\beta}}}{\partial z_{k}}=\frac{1}{2}\left(\bar{z}_{k} \delta_{i, \beta}+\bar{z}_{i} \delta_{k, \beta}-\bar{z}_{\beta} \delta_{i, k}\right)
$$

for which the substitution $z_{k}=z_{2}\left(\delta_{2, k}+\left(1-\delta_{2, k}\right) \xi\right)$, etc., and the inequality

$$
\left|\delta_{2, k}+\left(1-\delta_{2, k}\right) \xi\right| \leq \max \{1, r\}
$$

etc., yields

$$
\left|\frac{\partial g_{i, \bar{\beta}}}{\partial z_{k}}\right|,\left|\frac{\partial g_{\alpha, \bar{j}}}{\partial \bar{z}_{l}}\right| \leq \frac{3}{2} \max \{1, r\}\left|z_{2}\right| .
$$

Hence

$$
\left|\sum_{\alpha, \beta} g^{\alpha, \beta} \frac{\partial g_{i, \bar{\beta}}}{\partial z_{k}} \frac{\partial g_{\alpha, \bar{j}}}{\partial \bar{z}_{l}}\right| \leq 18 \frac{r^{2}+1}{\left(r^{2}-1\right)^{2}} \max \{1, r\}^{2}
$$

From the continuity in $r$ of this last expression for all $r \geq 1+\varepsilon$, and the fact that its limit as $r \rightarrow \infty$ is finite, we conclude that $R_{i, j, k, \bar{l}}$ is uniformly bounded for $|1-|\xi|| \geq \varepsilon$.

In the light of this example we make a final remark. Let $\Psi$ be any plurisubharmonic function on $\mathbb{C}^{2}$ satisfying the following:

- $\operatorname{det}\left(\frac{\partial^{2} \Psi}{\partial \bar{z}_{j} \partial z_{i}}\right)=|\langle Z, \bar{Z}\rangle|^{s} e^{f(Z, \bar{Z})},(s \in \mathbb{R} \backslash\{0\}, f$ smooth $)$ on $\Omega$;
- outside the locus corresponding to $\{\langle Z, \bar{Z}\rangle=0\}$, the form $\omega=\mathbf{i} \partial \bar{\partial} \Psi$ has bounded curvature on the domain $\Omega=\{\Psi<1\}$.
Then the pair of transversely intersecting discs corresponding to the locus above is extremally embedded in the sense that the Kobayashi distance between two points $( \pm \mathbf{i} \zeta, \zeta)$ and $\left( \pm \mathbf{i} \zeta^{\prime}, \zeta^{\prime}\right)$ in the same disc is equal to the Poincare distance between $\zeta$ and $\zeta^{\prime}$. In other words, the (0)(2)-orbit of weakly pseudoconvex points on $\partial \Omega$ bounds a pair of extremally embedded discs. Moreover, the Ricci form $\mathbf{i} \partial \bar{\partial} \log (\operatorname{det}(g))$ associated with the metric vanishes identically when $f$ is constant.

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