# A NOTE ON PERIODIC SOLUTIOANS OF SOME <br> NOINAUTONOHOUS DIFFERENTIAL EQUATIONS 

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#### Abstract

We prove the existence of nontrivial periodic solutions of some nonlinear ordinary differential equations with time-dependent coefficients using variational methods.


## 0 . Introduction and statement of the results

In this work we study the existence, under suitable conditions, of nontrivial $T$-periodic solutions of the following nonlinear equations:

$$
\begin{align*}
& \ddot{x}-\alpha(t) x-\beta(t) x^{2}+\gamma(t) x^{3}=0  \tag{0.1}\\
& \ddot{x}+\alpha(t) x-\beta(t) x^{2}+\gamma(t) x^{3}=0 \tag{0.2}
\end{align*}
$$

where $\alpha, \beta$ and $\gamma$ are measurable $T$-periodic functions such that if we denote by $\alpha, A, C$ and $C$ the infimum and supremum of $\alpha$ and $\gamma$, respectively, and $B=\|\beta\| L^{\infty}$, then
(0.3) $\quad 0<a \leqslant \alpha(t) \leqslant A<\infty, \quad 0<c \leqslant \gamma(t) \leqslant C<\infty$ and $B<\infty$.

The study of these equations was suggested by a paper of Cronin [4] which deals with an equation related with the biomathematical model of the aneurysm of the circle of Willis introduced by Austin [2]. In fact equation ( 0.2 ) is the homogeneous analogue of the equation studied in [4] for the case in which $\alpha, \beta$ and $\gamma$ are positive constants. Also in [4] there is a forcing term of the type $k \cdot \cos (\omega t)$.

Solutions of (0.1) and (0.2) will always be considered in the sense

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of $C^{1}$-functions, $x$, such that $\ddot{x}$ exists almost everywhere and $x(t)=x(t+T)$.

For the sake of simplification of computation we shall deal with the case of period $2 \pi$. The results, however, apply to every period $T$ with obvious modifications.

Abstract results concerning the existence of nontrivial periodic solutions of nonlinear problems have recently appeared in the literature (see [3], [5] and [7]). However they do not apply to equations (0.1) or (0.2). We also observe that, if $\alpha, \beta$ and $\gamma$ are constants, (0.1) and (0.2) can be studied using the phase plane.

We prove the following theorems:
THEOREM 0.1. Let $\alpha, B$ and $\gamma$ be measurable $2 \pi$-periodic functions that satisfy (0.3). Then equation (0.1) has a nontrivial $2 \pi$-periodic solution.

THEOREM 0.2. Let $\alpha, \beta$ and $\gamma$ be measurable $2 \pi$-periodic functions that satisfy (0.3) and such that (0.4) $m^{2}<a \leqslant A<(m+1)^{2}$ for some integer $m \geqslant 0$,
and

$$
\begin{equation*}
B^{2} \leqslant \frac{9}{2} \delta c, \text { where } \delta=a-m^{2} \tag{0.5}
\end{equation*}
$$

Then equation (0.2) has a nontrivial $2 \pi$-periodic solution.
This work is divided into two parts in which we prove theorems 0.1 and 0.2, respectively. Those proofs have an analogous structure and we use similar notations in both. The arguments and computations of section 1 are, however, much simpler than those of section 2 . We consider the interval $[0,2 \pi]$ and first we solve the projections of (0.1) and (0.2) onto spaces of finite dimension using in the case of (0.1) a "Mountain Pass Lemma" ([11) and in (0.2) a generalization of that lemma ([6]). Then, after adequate estimates, we pass to the limit. We observe that working with the finite dimensional approach the boundary conditions are well defined (see (1.2) and (2.2)) and the variational principles that we use are simpler.

NOTATIONS. Throughout this paper we use the standard spaces $L^{p}(0,2 \pi)$, $C^{p}(\overline{0,2 \pi})$, and the Sobolev spaces $H^{p}(0,2 \pi)=h^{p, 2}(0,2 \pi)$ which we denote
simply by $I^{p}, C^{p}$ and $H^{p}$. We use the symbols $|\cdot|_{p}$ and $\|\cdot\|_{p}$ to denote the usual norms of $L^{p}$ and $H^{p}$, respectively.

1. Equation (0.1)

Consider the interval $[0,2 \pi]$ and the equation

$$
\begin{equation*}
\ddot{u}-\alpha(t) u-\beta(t) u^{2}+\gamma(t) u^{3}=0 \tag{1.1}
\end{equation*}
$$

with the periodic conditions

$$
\begin{equation*}
u(0)=u(2 \pi), \quad \dot{u}(0)=\dot{u}(2 \pi) \tag{1.2}
\end{equation*}
$$

where $\alpha, \beta$ and $\gamma$ are measurable $2 \pi$-periodic functions that satisfy (0.3).

Proof of Theorem 0.1. The proof is divided in three steps:
Step 1: Approximate solution in finite dimension
For each positive integer $N$ consider the finite dimensional space

$$
Y_{N}=\left\{\sum_{k=-N}^{N} c_{k} e^{i k t}: c_{k} \in \mathrm{C} \text { and the sum is real }\right\}
$$

and the functional $J_{N}: Y_{N} \rightarrow \mathrm{R}$ defined by

$$
J_{N}(u)=\int_{0}^{2 \pi}\left(\frac{1}{2}|\dot{u}|^{2}+\frac{1}{2} \alpha(t) u^{2}+\frac{1}{3} B(t) u^{3}-\frac{1}{4} \gamma(t) u^{4}\right) d t
$$

It is easy to see that every critical point of $J_{N}$ is a solution of the equation

$$
\begin{equation*}
\ddot{x}-\alpha(t) x=P_{N}\left(\beta(t) x^{2}-\gamma(t) x^{3}\right) \tag{1.3}
\end{equation*}
$$

where $P_{N}$ is the orthogonal projection of $L^{2}(0,2 \pi)$ onto $Y_{N}$. So we are going to look for critical points of $J_{N}$. The tool we use is the "Mountain Pass Lemma" contained in [1].

By Hölder's inequality and (0.3) we have for $u \in Y_{N}$

$$
J_{I V}(u) \geqslant \frac{a}{2}|u|_{2}^{2}-\frac{B}{3}|u|_{3}^{3}-\frac{C}{4}|u|_{4}^{4} \geqslant|u|_{2}^{2}\left(a_{1}-a_{2}|u|_{2}-a_{3}|u|_{2}^{2}\right)
$$

(where $a_{1}, a_{2}, a_{3}$ are positive constants) which implies that there exist $r, \rho>0$ such that $J(u)>0$ if $0<|u|_{2} \leqslant r$ and $J(u) \geqslant \rho>0$ if
$|u|_{2}=r$. On the other hand, denoting by $\|\cdot\|_{\alpha}$ the norm $\|u\|_{\alpha}=\left(\int_{0}^{2 \pi}\left(|\dot{u}|^{2}+\alpha(t) u^{2}\right) d t\right)^{\frac{1}{2}}$, since $Y_{N}$ has finite dimension, we derive

$$
J_{N}(u) \leqslant \frac{1}{2}\|u\|_{\alpha}^{2}+\frac{B}{3}|u|_{3}^{3}-\frac{c}{4}|u|_{4}^{4} \leqslant|u|_{2}^{2}\left(a_{4}+a_{5}|u|_{2}-a_{6}|u|_{2}^{2}\right)
$$

(where $a_{4}, a_{5}$ and $a_{6}$ are positive constants) which implies that

$$
\begin{equation*}
J_{N}(u)<0 \text { for }|u|_{2} \text { large enough . } \tag{1.4}
\end{equation*}
$$

and therefore there is $w_{N} \in Y_{N}$ such that

$$
\begin{equation*}
\left|w_{N}\right|_{2}>r \text { and } J_{N}\left(w_{N}\right)=0 . \tag{1.5}
\end{equation*}
$$

In particular there is $w_{N}$ in $Y_{N}$ such that

$$
\begin{equation*}
\left|w_{N}\right|_{2}>r \text { and } J_{N}\left(w_{N}\right)=0, \tag{1.5}
\end{equation*}
$$

and obviously we can take $w_{N}=w \in Y_{1}$ since $J_{N \mid Y_{1}}=J_{1}$. This fact will be useful in step 2. Inequality (1.4) and the fact that $Y_{N}$ has finite dimension show that the Palais-Smale condition is trivially satisfied in $[0,+\infty)$. Then by the "Mountain Pass Lemma" ([1]) $J_{N}$ has a nonzero critical point, $u_{N}$, and a corresponding critical value $b_{N}$ characterized by:

$$
\begin{equation*}
b_{N}=\inf _{g \in \Gamma} \max _{y \in[0,1]} J_{N}(g(y))>0, \tag{1.6}
\end{equation*}
$$

where

$$
\Gamma_{N}=\left\{g \in C\left([0,1], Y_{N}\right): g(0)=0, g(1)=w\right\}
$$

Step 2: Estimates for $u_{\text {IV }}$
Now we are going to prove that there are positive constants independent of $N, A_{1}$ and $A_{2}$, such that

$$
\begin{gather*}
\left\|u_{N}\right\|_{2} \leqslant A_{1}  \tag{1.7}\\
\left|u_{N}\right|_{\infty} \geqslant A_{2}>0 . \tag{1.8}
\end{gather*}
$$

Let $b_{N}$ be the critical value that corresponds to the critical point of $J_{N}, u_{N}$. From its variational characterization (1.6) and from the fact that, for $t \in[0,1], g(t)=t w \in \Gamma_{N}$ we derive

$$
\begin{equation*}
0<b_{N}=\inf _{g \in \Gamma_{N}} \max _{t \in[0,1]} J_{N}(g(t)) \leqslant \max _{t \in[0,1]} J_{N}(t w)=J_{1}(\bar{\omega}) \tag{1.9}
\end{equation*}
$$

where $\bar{w}$ is an element of $Y_{1}$ in which $\max _{t \in[0,1]} J_{1}(t w)$ is achieved. If we denote by $\left\langle.\right.$, 〉 the duality bracket between $Y_{N}$ and its dual, by (1.9) and by the fact that $u_{N}$ is a critical point of $J_{N}$ we have

$$
\begin{aligned}
0<b_{N} & =J\left(u_{N}\right)-\frac{1}{3}\left\langle J^{\prime}\left(u_{N}\right), u_{N}\right\rangle \\
& =\frac{1}{6} \int_{0}^{2 \pi}\left(\left|\dot{u}_{N}\right|^{2}+\alpha u_{N}^{2}\right)+\frac{1}{12} \int_{0}^{2 \pi} \gamma u_{N}^{4} \leqslant J_{1}(\bar{w}),
\end{aligned}
$$

which implies the existence of a constant $B_{1}$ (independent of $N$ ) such that

$$
(1.10) \quad\left\|u_{N}\right\|_{\alpha}<B_{1}
$$

and, since $\|\cdot\|_{\alpha}$ is equivalent to the standard $H^{1}$-norm, by the Sobolev imbedding theorem, together with (0.3) and (1.3) it follows easily that

$$
\left|\ddot{u}_{N}\right|_{2} \leqslant\left|\ddot{u}_{N}-\alpha u_{N}\right|_{2}+\left|\alpha u_{N}\right|_{2} \leqslant\left|P_{N}\left(\beta u_{N}^{2}-\gamma u_{N}^{3}\right)\right|_{2}+B_{2} \leqslant B_{3}
$$

where $B_{2}$ and $B_{3}$ are positive constants independent of $N$. Hence (1.7) holds.

As for (1.8), we observe that, since $u_{N}$ is a critical point of $J_{N}$ in $Y_{N}$, we have

$$
\left\langle J^{\prime}\left(u_{N}\right), u_{N}\right\rangle=\int_{0}^{2 \pi}\left(\left|\dot{u}_{N}\right|^{2}+\alpha u_{N}^{2}+\beta u_{N}^{3}-\gamma u_{N}^{4}\right)=0
$$

Therefore by (0.3)

$$
0<a\left|u_{N}\right|_{2}^{2} \leqslant B\left|u_{N}\right|_{3}^{3}+C\left|u_{N}\right|_{4}^{4} \leqslant\left(B\left|u_{N}\right|_{\infty}+C\left|u_{N}\right|_{\infty}^{2}\right) \quad\left|u_{N}\right|_{2}^{2}
$$

which implies (1.8) since $u_{N} \neq 0$.

Step 3: Passing to the limit
By (1.7) and imbedding theorems, $\left(u_{N}\right)$ has a subsequence (which we still denote $\left(u_{N}\right)$ ) such that

$$
u_{N}-u \text { in } H^{2}
$$

and

$$
u_{N} \rightarrow u \text { in } C^{1}
$$

Estimate (1.8) ensures that $u$ is nonzero. Since

$$
P_{N}\left(\beta u_{N}^{2}-\gamma u_{N}^{3}\right) \rightarrow \beta u^{2}-\gamma u^{3} \text { in } L^{2}
$$

(1.3) shows that $u$ satisfies the equation

$$
\begin{equation*}
\ddot{u}-\alpha u-\beta u^{2}+\gamma u^{3}=0 \tag{1.9}
\end{equation*}
$$

in the $L^{2}$ sense. The periodicity of $u_{N}$ and $\dot{u}_{N}$ and the $C^{l}$-convergence obviously imply

$$
\begin{equation*}
u(0)=u(2 \pi), \quad \dot{u}(0)=\dot{u}(2 \pi) \tag{1.10}
\end{equation*}
$$

(1.9) and (1.10) show that $u$ can be extended to $(-\infty,+\infty)$ as a $C^{1}$-function, with period $2 \pi$, solving equation (0.1). This ends the proof.

REMARK 1. We observe that if $\alpha, \beta$ and $\gamma$ are $c^{P}$-functions, solutions of equation (0.1) are classical solutions, more precisely, solutions of class $c^{P+2}$.
2. Equation (0.2)

Consider the interval $[0,2 \pi]$ and the equation

$$
\begin{equation*}
\ddot{u}+\alpha(t) u-\beta(t) u^{2}+\gamma(t) u^{3}=0 \tag{2.1}
\end{equation*}
$$

with the periodic conditions

$$
\begin{equation*}
u(0)=u(2 \pi), \dot{u}(0)=\dot{u}(2 \pi) \tag{2.2}
\end{equation*}
$$

where $\alpha, \beta$ and $\gamma$ are functions that satisfy (0.3).
In the proof of Theorem (0.2) we use the following result, which follows easily from a theorem of Rabinowitz [6], and which is a generalization of the "Mountain Pass Lemma"

THEOREM 2.1. Let $H$ be a finite dimensional Hilbert space and $H_{1}, H_{2}$ subspaces of $H$ such that $H_{2}=H_{1}^{\perp}$. Suppose that $J \in C^{1}(H, R)$ and $J(x) \leqslant 0$ for every $x \in H_{1}$. If there are constants $r_{1}>r_{2}>0$ such that $J(x)>0$ in $\left.\left(B_{r_{2}} \backslash\{0\}\right) \cap H_{2}\right)$ and $J(x) \leqslant 0$ in $H \backslash B_{r_{1}}$, then $J$ has a critical point in $\{x \in H: J(x)>0\}$ and a corresponding critical value characterized by

$$
c=\inf _{h \in \Gamma} \max _{x \in H_{3} \cap \bar{B}_{r_{1}}} J(h(x))>0
$$

where, choosing $y \in H_{2}, H_{3}=H_{1} \oplus \operatorname{spon}\{y\}$ and

$$
\Gamma=\left\{h \in C\left(H_{3} \cap \bar{B}_{r_{1}}, H\right): h(x)=x \text { if } J(x) \leqslant 0\right\} .
$$

Proof of Theorem 0.2. As in Theorem 0.1 we divide the proof into three steps:

Step 1: Approximate solution in finite dimension
For each positive integer $N$ such that $N^{2}>A$ consider the finite dimensional space defined in the first step of the proof of Theorem 0.1, $Y_{N}$, and the functional $J_{N}: Y_{N} \rightarrow \mathrm{R}$ defined by

$$
J_{N}(u)=\int_{0}^{2 \pi}\left(\frac{|\dot{u}|^{2}}{2}-\frac{\alpha(t)}{2} u^{2}+\frac{\beta(t)}{3} u^{3}-\frac{\gamma(t)}{4} u^{4}\right) d t
$$

Analogously to (1.3) every critical point of $J_{N}$ is a solution of the equation

$$
\begin{equation*}
\ddot{x}+\alpha(t) x=P_{N}\left(\beta(t) x^{2}-\gamma(t) x^{3}\right), \tag{2.3}
\end{equation*}
$$

and therefore we are going to look for critical points of $J_{N}$.
Let $m$ be as in the statement of Theorem (0.2) and consider the following subspaces of $Y_{N}$ :

$$
Y_{N}^{+}=\left\{\sum_{k=-N}^{N} c_{k} e^{i k t} \in Y_{N}: k^{2} \geqslant(m+1)^{2}\right\}
$$

and

$$
Y_{N}^{-}=\left\{\sum_{k=-m}^{m} c_{k} e^{i k t} \in Y_{N}\right\}
$$

It is clear that $Y_{N}=Y_{N}^{+} \oplus Y_{N}^{-}$.
We claim that:

$$
\begin{equation*}
J_{N}(u) \leqslant 0 \quad \text { for every } \quad u \in Y_{N}^{-} \tag{2.4}
\end{equation*}
$$

and
(2.5) there are constants $r_{1}>r_{2}>0$ such that $J_{N}(u)>0$

$$
\text { in } Y_{N}^{+} \cap\left(B_{r_{2}} \backslash\{0\}\right) \text { and } J_{N}(u) \leqslant 0 \text { in } Y_{N} \backslash B_{r_{1}}
$$

In fact if $u \in Y_{N}^{-}$we have

$$
\begin{aligned}
J_{N}(u) & \leqslant \pi \sum_{k=-m}^{m}\left|c_{k}\right|^{2}\left(k^{2}-a\right)+\int_{0}^{2 \pi}\left(\frac{\beta}{3} u^{3}-\frac{\gamma}{4} u^{4}\right) \\
& \leqslant \pi\left(m^{2}-a\right) \sum_{k=-m}^{m}\left|c_{k}\right|^{2}+\int_{0}^{2 \pi}\left(\frac{B}{3}|u|^{3}-\frac{c}{4} u^{4}\right) \\
& =\int_{0}^{2 \pi} u^{2}\left(\frac{m^{2}-a}{2}+\frac{B}{3}|u|-\frac{c}{4} u^{2}\right)
\end{aligned}
$$

which, by condition (0.5), implies (2.4). As for (2.5), take $u \in Y_{N}^{+}$. Computing $J_{N}(u)$ we have

$$
\begin{aligned}
& J_{N}(u) \geqslant \pi \sum_{k=-N}^{N}\left|c_{k}\right|^{2}\left(k^{2}-A\right)+\int_{0}^{2 \pi}\left(\frac{\beta}{3} u^{3}-\frac{\gamma}{4} u^{4}\right) \\
& \geqslant \frac{k^{2} \geqslant(m+1)^{2}}{2}-A \\
& \geqslant|u|_{2}^{2}-\frac{B}{3}|u|_{3}^{3}-\frac{C}{4}|u|_{4}^{4} \\
&\left(q_{1}-q_{2}|u|_{2}-q_{3}|u|_{2}^{2}\right)
\end{aligned}
$$

This inequality implies the first statement of (2.5). As for the second one it is obtained as in the proof of Theorem 0.1 (see (1.4)). Hence, by Theorem 2.1, $J_{N}$ has a critical point in $Y_{N}$ and corresponding critical value $c_{N}$ characterized by
(2.6)

$$
c_{N}=\inf _{h \in \Gamma_{N}} \max _{u \in\left(Y_{N}^{-} \operatorname{span}(y\}\right) \cap \bar{B}_{r_{1}}} J_{N}(h(u))>0,
$$

where $y \in Y_{N}^{+}$and

$$
\left.\Gamma_{N}=\left\{h \in C\left(Y_{N}^{-} \oplus \operatorname{span}\{y\}\right) \cap \bar{B}_{r_{1}}, Y_{N}\right): h(u)=u \text { if } J_{N}(u) \leqslant 0\right\}
$$

Step 2: Estinates for $U_{N}$
As in the proof of Theorem 0.1 we are going to prove that there are constants independent of $N, A_{1}$ and $A_{2}$, such that

$$
\begin{equation*}
\left\|u_{N}\right\|_{2} \leqslant A_{1} \tag{2.7}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|u_{N}\right|_{\infty} \geqslant A_{2}>0 . \tag{2.8}
\end{equation*}
$$

Let $c_{N}$ be the critical value that corresponds to the critical point of $J_{N}, u_{N}$. We claim that

$$
\begin{equation*}
c_{N} \leqslant B_{1} \tag{2.9}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|u_{N}\right|_{4} \leqslant B_{2}, \tag{2.10}
\end{equation*}
$$

where $B_{1}$ and $B_{2}$ are positive constants independent of $N$. In fact, using the characterization of $c_{N}(2.6)$ and the fact that $h(u)=u \in \Gamma_{N}$, we have

$$
\begin{aligned}
0<c_{N} & =\inf _{h \in \Gamma_{N}} \max _{u \in\left(Y_{N}^{-} \oplus \operatorname{span}\{y\}\right) n \bar{B}_{r_{1}}} J_{N}(h(u)) \\
& \leqslant \max _{u \in\left(Y_{N}^{-} \operatorname{span}\{y\}\right) \cap \bar{B}_{r_{1}}} J_{N}(u) \\
& \leqslant \max _{u \in Y_{N}^{-\oplus} \operatorname{span}\{y\}} J_{N}(u) .
\end{aligned}
$$

 Suppose that the maximum is achieved at $\bar{u}_{N}$. Take $y=\cos j t$ where $j$ is an integer such that $N^{2} \geqslant j^{2} \geqslant(m+1)^{2}$. Then $\bar{u}_{n}=\sum_{k=-m}^{m} c_{k} e^{i k t}+\lambda \cos j t$
and we have

$$
\begin{aligned}
0 \leqslant c_{N} \leqslant J_{N}\left(\bar{u}_{N}\right) & \leqslant \sum_{k=-m}^{m}\left(k^{2}-a\right)\left|c_{k}\right|^{2} \pi+\frac{\pi}{2} \lambda^{2}\left(j^{2}-a\right)+\int_{0}^{2 \pi}\left(\frac{\beta}{3} \bar{u}_{N}^{3}-\frac{\gamma}{4} \bar{u}_{N}^{4}\right) \\
& \leqslant \pi \lambda^{2}+\frac{B}{3} \int_{0}^{2 \pi}\left|\bar{u}_{N}\right|^{3}-\frac{c}{4} \int_{0}^{2 \pi}\left|\bar{u}_{N}\right|^{4} \\
& \leqslant\left|\bar{u}_{N}\right|_{2}^{2}+\frac{B}{3}\left|\bar{u}_{N}\right|_{3}^{3}-\frac{c}{4}\left|\bar{u}_{N}\right|_{4}^{4}
\end{aligned}
$$

Then

$$
\frac{c}{4}\left|\bar{u}_{N}\right|_{4}^{4} \leqslant\left|\bar{u}_{N}\right|_{2}^{2}+\frac{B}{3}\left|\bar{u}_{N}\right|_{3}^{3}
$$

which implies the existence of a constant $B_{3}$, independent of $N$, such that

$$
\begin{equation*}
\left|\bar{u}_{N}\right|_{4} \leqslant B_{3} \tag{2.11}
\end{equation*}
$$

Let $\bar{c}_{N}=J_{N}\left(\bar{u}_{N}\right)$. If we get a bound for $\bar{c}_{N}$, (2.9) is proved since $c_{i V} \leqslant \bar{c}_{N}$. Using the fact that $\bar{u}_{N}$ is a critical point of $J_{N}$ in $Y_{N}^{-} \oplus \operatorname{span}\{y\}$ and Hölder's inequality, we get

$$
\begin{aligned}
c_{N} \leqslant \bar{c}_{N} & =\bar{c}_{N}-\frac{1}{2}\left\langle J_{N}^{\prime}\left(\bar{u}_{N}\right), \bar{u}_{N}\right\rangle \\
& =\int_{0}^{2 \pi}\left(\left(\frac{B}{3}-\frac{\beta}{2}\right) \bar{u}_{N}^{3}+\left(\frac{\gamma}{2}-\frac{\gamma}{4}\right) \bar{u}_{N N}^{4}\right) \\
& \leqslant \frac{B}{6}\left|\bar{u}_{N}\right|_{3}^{3}+\frac{c}{4}\left|\bar{u}_{N}\right|_{4}^{4} \leqslant B_{1}
\end{aligned}
$$

because of (2.11). Then (2.9) is proved. The fact that $u_{N}$ is a critical point of $J_{N}$ in $Y_{N}$ and (2.9) imply

$$
\begin{aligned}
B_{1} \geqslant c_{N} & =J_{N}\left(u_{N}\right)-\frac{1}{2}\left(J_{N}^{\prime}\left(u_{N}\right), u_{N}\right\rangle \\
& =\int_{0}^{2 \pi}\left(-\frac{\beta}{6} u_{N}^{3}+\frac{\gamma}{4} u_{N}^{4}\right) \\
& \geqslant-\frac{B}{6}\left|u_{N}\right|_{3}^{3}+\frac{c}{4}\left|u_{N}\right|_{4}^{4}
\end{aligned}
$$

$$
\frac{c}{4}\left|u_{N}\right|_{4}^{4} \leqslant B_{1}+\frac{B}{6}\left|u_{N}\right|_{3}^{3}
$$

which implies (2.10).
Now using (2.10), a Galiardo-Nirenberg interpolation and the Sobolev imbedding theorem we are going to prove (2.7). Throughout this proof we denote by $k$ several positive constants independent of $N$. By (2.3) , (2.10) and Hölder's inequality we have

$$
\begin{aligned}
\left|u_{N}\right|_{2} & \leqslant\left|u_{N}+\alpha u_{N}\right|_{2}+\left|\alpha u_{N}\right|_{2} \\
& \leqslant\left|P_{N}\left(\beta u_{N}^{2}-\gamma u_{N}^{3}\right)\right|_{2}+k \\
& \leqslant B\left|u_{N}^{2}\right|_{2}+C\left|u_{N}^{3}\right|_{2}+k \\
& \leqslant k\left|u_{N}\right|_{\infty}\left|u_{N}\right|_{4}^{2}+k \\
& \leqslant k\left|u_{N}\right|_{\infty}+k
\end{aligned}
$$

## then

$$
\begin{equation*}
\left|u_{N}\right|_{2} \leqslant k+k\left|u_{N}\right|_{\infty} . \tag{2.12}
\end{equation*}
$$

Write $u_{N}$ in the form $u_{N}=a_{N}+\tilde{u}_{N}$ where $\int_{0}^{2 \pi} \tilde{u}_{N}=0$ and $a_{N}$ is a constant. Since $\tilde{u}_{N}$ vanishes for some $t_{0} \in[0,2 \pi]$ we can apply the Gagliardo-Nirenberg interpolation to $\tilde{u}_{N}$ :

$$
\begin{equation*}
\left|\tilde{u}_{N}\right|_{\infty} \leqslant M\left|\tilde{u}_{N}\right|_{4}^{\frac{2}{3}}\left|\dot{\tilde{u}}_{N}\right|_{2}^{\frac{1}{3}} \tag{2.13}
\end{equation*}
$$

By (2.10) and Hölder's inequality

$$
\begin{equation*}
\left|a_{N}\right| \leqslant \frac{1}{2 \pi}\left|u_{N}\right|_{1} \leqslant k \tag{2.14}
\end{equation*}
$$

By (2.13), (2.14) and (2.10) we have

$$
\begin{align*}
\left|u_{N}\right|_{\infty} & \leqslant k+\left|\tilde{u}_{N}\right|_{\infty}  \tag{2.15}\\
& \leqslant k+M\left|\tilde{u}_{N}\right|_{4}^{\frac{2}{3}}\left|\dot{\tilde{u}}_{N}\right|_{2}^{\frac{1}{3}} \\
& \leqslant k+M\left(k+\left|u_{N}\right|_{4}\right)^{\frac{2}{3}}\left|\dot{u}_{N}\right|_{2}^{\frac{1}{3}} \\
& \leqslant k+k\left|\dot{u}_{N}\right|_{2}^{\frac{1}{3}}
\end{align*}
$$

By (2.12) and (2.15) we get

$$
\begin{equation*}
\left|u_{N}\right|_{2} \leqslant k+k\left|\dot{u}_{N}\right|_{2}^{\frac{1}{3}} \tag{2.16}
\end{equation*}
$$

Since $\left|\dot{u}_{N}\right|_{2} \leqslant\left|\ddot{u}_{N}\right|_{2}$, from (2.16), (2.10) and Hölder's inequality we conclude that (2.7) holds.

As for (2.8), we observe that from the variational characterization of the eigenvalues and by (0.4) it follows that $\lambda=0$ is not an eigenvalue of the linear problem

$$
\left\{\begin{array}{l}
\ddot{u}+\alpha(t) u=\lambda u, \\
u(0)-u(2 \pi)=\dot{u}(0)-\dot{u}(2 \pi)=0 .
\end{array}\right.
$$

From this fact, it is easy to see, arguing by contradiction, that there is $k>0$ such that

$$
|\ddot{u}+\alpha u|_{2} \geqslant k|u|_{2}
$$

for every $u \in H^{2}$ that satisfies the periodic conditions (2.2). Then, using (2.3),

$$
\begin{aligned}
k\left|u_{N}\right|_{2} & \leqslant\left|P_{N}\left(\beta u_{N}^{2}-\gamma u_{N}^{3}\right)\right|_{2} \\
& \leqslant\left(B\left|u_{N}\right|_{\infty}+C\left|u_{N}\right|_{\infty}^{2}\right)\left|u_{N}\right|_{2}
\end{aligned}
$$

which implies $(2.8)$, since $u_{N} \neq 0$.
Step 3: Passing to the limit
This is analogous to step 3 of the proof of Theorem 0.1.
This ends the proof.

REMARK 2.1. As in the case of remark 1.1 , the regularity of the solution thus obtained depends on the regularity of the coefficient functions $\alpha, \beta$ and $\gamma$.

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