# THE SPIN REPRESENTATION OF THE SYMMETRIC GROUP 

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1. Let $\Gamma_{n}$ be the representation group or spin group $(9 ; 4)$ of the symmetric group $S_{n}$. Then the irreducible representations of $\Gamma_{n}$ can be allocated into two classes which we shall call (i) ordinary representations, which are the irreducible representations of the symmetric group, and (ii) spin or projective representations.

As is well known (3;5), there is an ordinary irreducible representation [ $\lambda$ ] corresponding to every partition $(\lambda)=\left(\lambda_{1}, \lambda_{2}, \ldots, \lambda_{m}\right)$ of $n$ with

$$
\lambda_{1} \geqslant \lambda_{2} \geqslant \ldots \geqslant \lambda_{m}>0 .
$$

Of fundamental importance in the study of the ordinary representation is the concept of a hook graph (2;6;8). Our aim in this note is to develop a similar concept for the spin representations of $\Gamma_{n}$.

There is an irreducible spin representation $\langle\lambda\rangle$ of $\Gamma_{n}$ corresponding to every partition $(\lambda)=\left(\lambda_{1}, \lambda_{2}, \ldots, \lambda_{m}\right)$ of $n$ with $\lambda_{1}>\lambda_{2}>\ldots>\lambda_{m}>0$. In the following, we shall say that a partition satisfies Condition A if it has no equal parts, where the parts are not necessarily in descending order.

Any partition $(\lambda)=\left(\lambda_{1}, \lambda_{2}, \ldots, \lambda_{m}\right)$ can be associated with a graph consisting of rows of symbols called nodes. $\lambda_{1}$ in the first row, $\lambda_{2}$ in the second row, ..., $\lambda_{m}$ in the last row. The node in the $i$ th row and $j$ th column of $\langle\lambda\rangle$ is called its $(i, j)$-node.

Definition 1. An $(i, j)_{(k)}$-bar $(k=i, i+1, \ldots, m)$ of $\langle\lambda\rangle$ will consist of the $(i, j)$-node together with the remaining $\lambda_{i}-j$ nodes to the right of it and
(i) the $\lambda_{k}-j+1$ nodes from the $k$ th row of $\langle\lambda\rangle$ if $k>i$,
(ii) no further nodes if $k=i$,
so that the resulting graph on deleting these nodes satisfies Condition A.
From this definition, it is clear that there is more than one $(i, j)_{(k)}$-bar attached to the $(i, j)$-node. In fact, if $r_{i j}$ denotes the number of $(i, j)_{(k)}$-bars attached to the $(i, j)$-node, then $0 \leqslant r_{i j} \leqslant m-i+1$. Further, since the resulting graphs on deleting the nodes have to satisfy Condition A, we have that

$$
\begin{aligned}
& r_{i 1}=m-i+1 \quad(i=1,2, \ldots, m), \\
& r_{i j}= \begin{cases}1, & \text { if } j>1 \text { and } j-1 \neq \lambda_{k} \text { for any } k, \\
0, & \text { if } j>1 \text { and } j-1=\lambda_{k} \text { for some } k .\end{cases}
\end{aligned}
$$

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Definition 2. The length of an $(i, j)_{(k)}$-bar is

$$
S_{i j}^{(k)}= \begin{cases}\lambda_{i}-j+1 & \text { if } k=i \\ \lambda_{i}+\lambda_{k}-2 j+2 & \text { if } k>i\end{cases}
$$

In an $(i, j)_{(k)}$-bar $(k \geqslant i)$, the number $q=\lambda_{k}-j+1$ is called the arm length of the $(i, j)_{(k)}$-bar. An $(i, j)_{(k)}$-bar with arm length $q$ will be called a $q$-bar. An $(i, j)_{(i)}$-bar will be called an $O$-bar.

We easily see that

$$
S_{i 1}^{(k)}= \begin{cases}\lambda_{i} & \text { if } k=i \\ \lambda_{i}+\lambda_{k} & \text { if } k>i\end{cases}
$$

Let

$$
S_{i 1}=\prod_{k=i}^{m} s_{i 1}^{(k)},
$$

and

$$
S_{i j}= \begin{cases}\lambda_{i}-j+1 & \text { if } j>1 \text { and } j-1 \neq \lambda_{k} \text { for any } k, \\ 1 & \text { if } j>1 \text { and } j-1=\lambda_{k} \text { for some } k .\end{cases}
$$

Definition 3. The graph obtained by replacing the $(i, j)$-node of $\langle\lambda\rangle$ by $S_{i j}$ is called the bar graph $S\langle\lambda\rangle$. The bar product $S^{\lambda}$ is the product of all the $S_{i j}$ 's in $S\langle\lambda\rangle$, that is,

$$
S^{\lambda}=\prod_{i=1}^{m} \prod_{j=1}^{\lambda_{i}} S_{i j} .
$$

2. The degree of an irreducible spin representation $\langle\lambda\rangle$ of $\Gamma_{n}$. Schur (9) has shown that the degree $f^{\lambda}$ of $\langle\lambda\rangle$ is given by the formula

$$
\begin{equation*}
f^{\lambda}=2^{\left[\frac{1}{2}(n-m)\right]} \frac{n!}{\lambda_{1}!\lambda_{2}!\ldots \lambda_{m}!} \prod_{1 \leqslant r<s \leqslant m} \frac{\lambda_{T}-\lambda_{s}}{\lambda_{T}+\lambda_{s}}, \tag{1}
\end{equation*}
$$

where $[x]$ denotes the greatest integer less than or equal to $x$. We now prove the following.

Theorem 1. Let $S^{\lambda}$ denote the bar product of an irreducible spin representation $\langle\lambda\rangle$ of $\Gamma_{n}$. Then the degree $f^{\lambda}$ of $\langle\lambda\rangle$ is given by the formula

$$
f^{\lambda}=\frac{2^{\left[\frac{1}{2}(n-m)\right]} n!}{S^{\lambda}} .
$$

The proof follows closely the proof of the corresponding theorem on the degree of the ordinary representations of $\Gamma_{n}$ in terms of the hook product (2).

From § 1, we have

$$
\prod_{j=1}^{\lambda_{i}} S_{i j}=\lambda_{i} \prod_{k=i+1}^{m}\left(\lambda_{i}+\lambda_{k}\right) \prod_{j}\left(\lambda_{i}-j+1\right)
$$

where $j=2, \ldots, \lambda_{m}, \lambda_{m}+2, \ldots, \lambda_{m-1}, \lambda_{m-1}+2, \ldots, \lambda_{i-1}, \lambda_{i-1}+2, \ldots, \lambda_{i}$. Thus, we see that

$$
\begin{aligned}
\prod_{j=1}^{\lambda_{i}} S_{i j} & =\lambda_{i} \prod_{k=i+1}^{m}\left(\lambda_{i}+\lambda_{k}\right) \prod_{j=2}^{\lambda_{i}}\left(\lambda_{i}-j+1\right) / \prod_{k=i+1}^{m}\left(\lambda_{i}-\lambda_{k}\right) \\
& =\lambda_{i}!\prod_{k=i+1}^{m} \frac{\lambda_{i}+\lambda_{k}}{\lambda_{i}-\lambda_{k}} .
\end{aligned}
$$

Hence, it follows that

$$
S^{\lambda}=\prod_{i=1}^{m} \prod_{j=1}^{\lambda_{i}} S_{i j}=\lambda_{1}!\lambda_{2}!\ldots \lambda_{m}!\prod_{1 \leqslant i<k \leqslant m} \frac{\lambda_{i}+\lambda_{k}}{\lambda_{i}-\lambda_{k}}
$$

and from (1)

$$
f^{\lambda}=\frac{2^{\left[\frac{1}{2}(n-m)\right]} n!}{S^{\lambda}}
$$

Example. We find the degree of the irreducible spin representation $\langle\lambda\rangle=\langle 863\rangle$ of $\Gamma_{17} \cdot\langle\lambda\rangle$ has the graph
and the bar graph

$$
\begin{array}{ccccccccc}
S^{\lambda} & 14 \times 11 \times 8 & 7 & 6 & 1 & 4 & 3 & 1 & 1 \\
& 9 \times 6 & 5 & 4 & 1 & 2 & 1 & & \\
& 3 & 2 & 1 & & & &
\end{array}
$$

Hence,

$$
\begin{aligned}
f^{\lambda} & =\frac{2^{7} \times 17!}{14 \times 11 \times 9 \times 8 \times 7 \times 6 \times 4 \times 3 \times 6 \times 5 \times 4 \times 2 \times 3 \times 2 \times 1} \\
& =5,657,600
\end{aligned}
$$

3. The removal of a bar from a graph. First, we consider the effect of removing an $O$-bar from the graph $\langle\lambda\rangle$ corresponding to the partition $(\lambda)=$ $\left(\lambda_{1}, \lambda_{2}, \ldots, \lambda_{m}\right)$ of $n$. Suppose that an $O$-bar is removed from the $(i, j)$-node of this graph, that is, the $\lambda_{i}-j+1$ nodes from the $i$ th row of this graph to the right of and including the ( $i, j$ )-node, assuming that $j-1 \neq \lambda_{k}$ for any $k$. Then $j-1>\lambda_{k}$ and $j-1<\lambda_{k-1}$ for some value of $k, i \leqslant k \leqslant m$. Move up the nodes which are below the empty spaces to obtain a graph with parts in descending order, that is, the graph of the partition

$$
(\lambda)^{\prime}=\left(\lambda_{1}, \lambda_{2}, \ldots, \lambda_{i-1}, \lambda_{i+1}, \ldots, \lambda_{k-1}, j-1, \lambda_{k}, \ldots, \lambda_{m}\right)
$$

of $n-\lambda_{i}+j-1$. Sometimes, it is more convenient to consider removing
the equivalent rim of the graph $\langle\lambda\rangle$, that is, $\langle\lambda\rangle$ ' is obtained from $\langle\lambda\rangle$ by removing $\lambda_{i}-\lambda_{i+1}$ nodes from the $i$ th row, $\lambda_{i+1}-\lambda_{i+2}$ nodes from the $(i+1)$ th row, $\ldots, \lambda_{k-1}-j+1$ nodes from the $k$ th row.

For example, we remove an $O$-bar from the ( 2,2 )-node of the graph $\langle 11,8,6,4,3,2\rangle$.

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. . . . . - 
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..
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We shall call the equivalent rim removed a skew $O$-bar. An $(i, j)$-node on the skew $O$-bar is called a leg node of the skew $O$-bar if the $(i-1, j+1$ )-node also lies on the skew $O$-bar. The total number of leg nodes is called the leg length of the skew $O$-bar. The leg length is clearly $k-i-1$. In the above example, the nodes indicated by $\times$ are leg nodes and thus the leg length of the skew $O$-bar is 4 .

Next, we consider the effect of removing a $q$-bar from the graph $\langle\lambda\rangle$, for some value of $q$. Suppose that we remove an $(i, 1)_{k}$-bar $(k>i)$ from $\langle\lambda\rangle$, that is, $q=\lambda_{k}$. The resulting graph is

$$
\langle\lambda\rangle^{\prime}=\left\langle\lambda_{1}, \lambda_{2}, \ldots, \lambda_{i-1}, \lambda_{i+1}, \ldots, \lambda_{k-1}, \lambda_{k+1}, \ldots, \lambda_{m}\right\rangle
$$

of $n-\lambda_{i}-\lambda_{k}$. Thus, $\langle\lambda\rangle^{\prime}$ is obtained from $\langle\lambda\rangle$ by removing $\lambda_{i}-\lambda_{i+1}$ nodes from the $i$ th row, $\lambda_{i+1}-\lambda_{i+2}$ nodes from the $(i+1)$ th row, $\ldots, \lambda_{k-1}-\lambda_{k+1}$ nodes from the $(k-1)$ th row, $\lambda_{k}-\lambda_{k+2}$ nodes from the $k$ th row, $\lambda_{k+1}+\lambda_{k+3}$ nodes from the $(k+1)$ th row, $\ldots, \lambda_{m-2}-\lambda_{m}$ nodes from the $(m-2)$ th row, $\lambda_{m-1}$ nodes from the $(m-1)$ th row, and $\lambda_{m}$ nodes from the $m$ th row. Hence, this is equivalent to removing two skew $O$-bars from $\langle\lambda\rangle$; first remove a skew $O$-bar of length $\lambda_{k}$ and then remove a skew $O$-bar of length $\lambda_{i}$. This equivalent rim removed is called a skew $\lambda_{k}$-bar. The total number of leg nodes in both parts of the skew $\lambda_{k}$-bar is called the leg length of the skew $\lambda_{k}$-bar. In a skew $\lambda_{k}$-bar, $\lambda_{k}$ is the arm length defined in § 1 .

This is now illustrated by an example. Suppose that we remove a $(1,1)_{(3)}$-bar from $\langle\lambda\rangle=\langle 11,8,6,4,3,2\rangle$. That is


The leg length of this skew 6 -bar is $4+3=7$, and the arm length is 6 .
4. The spin characters of $\Gamma_{n}$. Using the terminology of § 3, we can now prove the following theorem, which corresponds to the well-known MurnaghanNakayama formula (8) for the ordinary characters of $\Gamma_{n}$.

Theorem 2. Let $T$ denote the graph corresponding to the partition ( $\lambda$ ) of $n$ and $T-S_{i}$ the graph obtained by removing $a$ bar $S_{i}$ of length $i$ from $T$. Let $\zeta_{\pi}(T)$ denote the irreducible spin character of the positive class $(\pi)=\left(1^{\alpha 1} 3^{\alpha 3} 5^{\alpha} \ldots\right)$ of $\Gamma_{n}$ corresponding to the graph $T$, then

$$
\zeta_{\pi}(T)=\sum_{i}(-1)^{h_{i}+k_{i}} 2^{\left[\frac{1}{2}\left(\epsilon^{\prime}-\epsilon+1\right)\right]} \zeta_{\pi^{\prime}}\left(T-S_{i}\right),
$$

where the summation runs over all possible bars $S_{i}$ of length $i$ which can be removed from $T ; h_{i}$ and $k_{i}$ denote the arm length and leg length respectively of $S_{i} ; \pi^{\prime}$ is obtained from $\pi$ by deleting a cycle of length $i ; \epsilon$ and $\epsilon^{\prime}$ are 0 or 1 according as $\zeta_{\pi}(T)$ and $\zeta_{\pi^{\prime}}\left(T-S_{i}\right)$ are double or associate spin characters.

Schur (9) has shown that the irreducible spin characters of $\Gamma_{n}$ are generated by a certain class of symmetric functions $Q_{\lambda}$, known as $Q$-functions (5), with the property that

$$
Q_{\lambda}=\sum_{\pi} \frac{h_{\pi}}{2 . n!} 2^{\frac{1}{2}(p+m+\epsilon)} \zeta_{\pi}{ }^{\lambda} S_{\pi}
$$

where $(\lambda)=\left(\lambda_{1}, \lambda_{2}, \ldots, \lambda_{m}\right),(\pi)=\left(1^{\alpha_{1} 3^{\alpha} 5^{\sigma_{5}}} \ldots\right), p=\alpha_{1}+\alpha_{3}+\alpha_{5}+\ldots$, and $\epsilon=0$ or 1 according as $\zeta_{\pi}{ }^{\lambda}$ is a double spin character or an associate spin character.

Remark. If $\zeta_{\rho}{ }^{\lambda}$ is a spin character, then $\zeta_{\rho^{\lambda^{\prime}}}=(-1)^{\alpha} \zeta_{\rho}{ }^{\lambda}$, where $\alpha=0$ or 1 according as ( $\rho$ ) is a positive or negative class, is a second spin character. $\zeta_{\rho}{ }^{\lambda}$ and $\zeta_{\rho}{ }^{\lambda \prime}$ are known as associate spin characters. If $\zeta_{\rho}{ }^{\lambda}=\zeta_{\rho}{ }^{\lambda 1}$ for all classes ( $\rho$ ), then $\zeta_{\rho}{ }^{\lambda}$ is a double spin character.
In (4), certain rules have been proved whereby a $Q_{\lambda}$ corresponding to a partition ( $\lambda$ ) with some parts possibly negative or zero and not necessarily in descending order is written in terms of a $Q_{(\lambda)}$ with positive parts in descending order. These rules are
(i) if any two parts of $Q_{\lambda}$ are equal, then $Q_{\lambda} \equiv 0$;
(ii) $Q_{\left(\lambda_{1} \lambda_{2} \ldots \lambda_{r} \lambda_{r+1} \ldots \lambda_{m}\right)}=-Q_{\left(\lambda_{1} \lambda_{2} \ldots \lambda_{r+1} \lambda_{r} \ldots \lambda_{m}\right)}$,
(iii) $Q_{\left(\lambda_{1} \lambda_{2} \ldots \lambda_{m}\right)} \equiv 0$ if $\lambda_{i}<0$ for any $1 \leqslant i \leqslant m$ and $\left|\lambda_{1}\right|,\left|\lambda_{2}\right|, \ldots,\left|\lambda_{m}\right|$ are all different,
(iv) $Q_{\left(\lambda_{1}, \lambda_{2}, \ldots,-\lambda_{r}, \lambda_{r}, \ldots, \lambda_{m}\right)}=2(-1)^{\lambda_{r}} Q_{\left(\lambda_{1} \lambda_{2} \ldots \lambda_{r-1} \lambda_{r+1} \ldots \lambda_{m}\right)}$ and

$$
Q_{\left(\lambda_{1} \lambda_{2} \ldots \lambda_{r},-\lambda_{r}, \ldots, \lambda_{m}\right)} \equiv 0 .
$$

In (5), it has been proved that

$$
\zeta_{\pi}{ }^{\lambda}=\sum_{(\mu)} k_{(\mu)} 2^{\frac{1}{2}\left(m^{\prime}-m+\epsilon^{\prime}-\epsilon+1\right)} \zeta_{\pi^{\prime}},
$$

where the summation is taken over all partitions $(\mu)=\left(\mu_{1}, \mu_{2}, \ldots, \mu_{m^{\prime}}\right)$ corresponding to the $Q_{\mu}$ obtained from $Q_{\left(\lambda_{12}, \lambda_{1}, \ldots, \lambda_{j-i}, \ldots, \lambda_{m}\right)}(j=1,2, \ldots, m)$
by means of rules (i)-(iv) above, $k_{(\mu)}$ is the coefficient of $Q_{\mu}$ and $(\pi)^{\prime}$ is the class obtained from ( $\pi$ ) by deleting a cycle of length $i$. Thus, in order to prove this theorem, we must show that

$$
\begin{equation*}
k_{(\mu)} 2^{\frac{1}{2}\left(m^{\prime}-m+\epsilon^{\prime}-\epsilon+1\right)}=(-1)^{h_{i}+k_{i}} 2^{\frac{1}{2}\left(\epsilon^{\prime}-\epsilon+1\right)} . \tag{2}
\end{equation*}
$$

From the rules (i)-(iv) above, $k_{(\mu)} \neq 0$ if and only if
(a) $\lambda_{j}-i \neq \lambda_{k}$ for any $k>j$,
(b) $\lambda_{j}-i=0$,
(c) $\lambda_{j}-i+\lambda_{k}=0$, for some $k>j$.

Case (a) is equivalent to removing an $O$-bar of length $i$ from $\langle\lambda\rangle$. It $\lambda_{j}-1 \neq 0, m^{\prime}=m$, and if $k_{i}$ denotes the arm length of the $O$-bar, $k_{i}=0$. Use (i) repeatedly until the partition has parts in descending order to give the value of $k_{\mu}$. Clearly, by the definition of the leg length $h_{i}$ of the skew $O$-bar, $k_{\mu}=(-1)^{h_{i}}$, and thus (2) follows in this case.

In case (b), $m^{\prime}=m-1$ and thus $m^{\prime}-m+1+\epsilon^{\prime}-\epsilon=\epsilon^{\prime}-\epsilon=0$ and by the same argument as for case (a), $k_{(\mu)}=(-1)^{h_{i}}$, and thus (2) follows in this case.

Case (c) is equivalent to removing a $\lambda_{k}$-bar of length $i$ from $\langle\lambda\rangle$. Now $m^{\prime}-m=-2$, and thus

$$
k_{\mu} 2^{\frac{1}{2}\left(m^{\prime}-m+1+\epsilon^{\prime}-\epsilon\right)}=\frac{1}{2} k_{\mu} 2^{\frac{1}{2}\left(1+\epsilon^{\prime}-\epsilon\right)}
$$

and by rules (i) and (iv) above, $k_{\mu}=(-1)^{h_{i}+k_{i}}$, where $h_{i}$ is the length of the skew $\lambda_{k}$-bar and $k_{i}=\lambda_{k}$ is the arm length of the skew $\lambda_{k}$-bar. Thus (b) follows in this case again, and the proof of the theorem is complete.

It would be of interest to obtain a direct proof of this theorem of the same type as the proof of the Murnaghan-Nakayama recursion formula given by Robinson (7). In order that this may be done, a theory corresponding to the theory of Young diagrams must be developed for the graph $\langle\lambda\rangle$.
5. In a future publication, the theory of bar graphs will be applied to the study of the modular representations of the group $\Gamma_{n}$. It will be shown that bar graphs play a similar role for the modular representations of $\Gamma_{n}$ as hook graphs play for the modular representation of the group $S_{n}$. For instance, we shall prove a result corresponding to the well-known Nakayama conjecture (6) first proved by Brauer and Robinson (1). If we define the $p$-core of a graph $\langle\lambda\rangle$ to be the graph obtained after removing all possible bars of length $p$ from $\langle\lambda\rangle$, then we shall prove that irreducible spin representations of $\Gamma_{n}$ belong to the same $p$-block if and only if their corresponding graphs have the same $p$-core.

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