

Averaged changes in the orbital elements of meteoroids due to Yarkovsky-Radzievskij force

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Abstract. Yarkovsky-Radzievskij effect exceeds the Poynting-Robertson effect in the perturbing action on particles larger than 100 μm . We obtained formulae for averaged changes in a meteoroid's Keplerian orbital elements and used them to estimate dispersion in the Geminid meteoroid stream. It was found that dispersion in semi-major axis of the model shower increased nearly three times on condition that meteoroids rotation is fast, and the rotation axis is stable.

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To obtain the averaged changes in a meteoroid's orbit we use equations for the derivatives of the orbital elements in the form suggested by Burns (1976), namely $db/dt = F(a, \dots, \omega, v, E, m, r, F_r, F_t, F_n)$, where b is one of the standard Keplerian orbital elements a, e, i, Ω, ω ; r is the heliocentric distance, m is the meteoroid mass, v is the true anomaly, E is the eccentric anomaly, and F_r, F_t, F_n are the radial, transverse, and normal components of the perturbing force. As in (Burns *et al.* 1979), we average the time variation of the elements over the orbital period under the assumption that a and e are essentially constant over this time interval and that the angular momentum is conserved. The average change in any element b is then:

$$\left\langle \frac{da}{dt} \right\rangle = \frac{1}{PH} \oint \frac{da}{dt} r^2 dv, \quad (1)$$

where P is the orbital period and $H = [a\mu(1 - e^2)]^{1/2}$ is the orbital angular momentum of the meteoroid per unit mass, μ is the gravitational constant of the sun.

The perturbing Yarkovsky-Radzievskij force for fast rotators, according to Burns *et al.* (1979) is

$$F_Y = k_Y r^{-7/2}, \quad k_Y = 2.962 r_m^2 c^{-1} \sigma^{1/4} (1 - \alpha)^{1/4} \gamma \omega_m^{-1/2} S_0^{7/4} r_0^{7/2} \cos \xi, \quad (2)$$

where $(1/\gamma) = 300 \text{ J}/(\text{m}^2 \cdot \text{s}^{1/2} \cdot \text{K})$ is the thermal inertia, r_m is the radius of the meteoroid, $\alpha = 0.1$ is the albedo, $\omega_m = 10^4 \text{ rad/s}$ is the angular velocity of the meteoroid ($r_m = 1 \text{ mm}$), S_0 is the solar constant at heliocentric distance $r_0 = 1 \text{ au}$, ξ is the angle between the rotational axis of the particle and its orbital plane ($\xi = 45^\circ$ for prograde, or $\xi = 135^\circ$ for retrograde rotation), c is the speed of light, and σ is the Stefan-Boltzmann constant.

The values of the parameters in (2) correspond to the model accepted by Olsson-Steel (1987). As in (Olsson-Steel 1987), we assume that $F_r = F_t = F_n = F_Y/\sqrt{3}$ due to the precession of the rotation axis and the frequent changes in the rotation speed. Even now, when computers are much faster, the numerical integration of (1) is still very expensive, so analytic expressions for the averaged changes in the orbital elements were found:

$$\begin{aligned} \langle \dot{a} \rangle &= k_1 I_{52}, & \langle \dot{e} \rangle &= k_2 (I_{32c} + e I_{12} + I_{12c}), & \langle \dot{i} \rangle &= k_2 I_{12c} \cos \omega, \\ \langle \dot{\Omega} \rangle &= k_2 I_{12c} \sin \omega / \sin i, & \langle \dot{\omega} \rangle &= -k_2 I_{32c} / e - \cos i \langle \dot{\Omega} \rangle, \end{aligned}$$

Table 1. Coefficients of expansion (3)

Integrals	p_0	p_1	p_2	p_3	p_4	p_5	p_6	p_7	p_8
I_{52}	2	0	$\frac{15}{8}$	0	$-\frac{15}{2 \cdot 16^2}$	0	$-\frac{25}{2 \cdot 16^3}$	0	$-\frac{1575}{2 \cdot 16^5}$
I_{32c}	0	$\frac{3}{2}$	0	$-\frac{3}{4 \cdot 16}$	0	$-\frac{15}{2 \cdot 16^2}$	0	$-\frac{315}{2 \cdot 16^4}$	0
I_{12}	2	0	$-\frac{1}{8}$	0	$-\frac{15}{2 \cdot 16^2}$	0	$-\frac{105}{2 \cdot 16^3}$	0	$-\frac{15015}{2 \cdot 16^5}$
I_{12c}	0	$\frac{1}{2}$	0	$\frac{3}{4 \cdot 16}$	0	$\frac{35}{8 \cdot 16^2}$	0	$-\frac{1155}{2 \cdot 16^4}$	0

where

$$\begin{aligned}
 k_1 &= k_Y / [\pi \sqrt{3\mu} ma^2 (1 - e^2)^{5/2}], & k_2 &= k_Y / [2\pi \sqrt{3\mu} ma^3 (1 - e^2)^{3/2}], \\
 I_{52} &= \oint (1 + e \cos v)^{5/2} dv, & I_{32c} &= \oint (1 + e \cos v)^{3/2} \cos v dv, \\
 I_{12} &= \oint (1 + e \cos v)^{1/2} dv, & I_{12c} &= \oint (1 + e \cos v)^{1/2} \cos v dv.
 \end{aligned}$$

The integrals I are easily found as series by expanding the integrands using the binomial formula

$$I = \pi \sum_{j=0}^{\infty} p_j e^j. \tag{3}$$

Accuracy to 1% is achieved by keeping terms to eighth order in the eccentricity. The expansion coefficients p_j in (3) are given in Table 1.

Using the formulae we made some estimations for the Geminid meteoroid stream. The model of the stream was like in (Ryabova 2007) and the meteoroid mass was taken as 3×10^{-3} g. Direction of rotation (prograde or retrograde) for each meteoroid was chosen using a pseudorandom number generator. The dispersion of the Geminids is anisotropic, so the total dispersion and the dispersion observed at the Earth differ. The width of the model shower is about $1.5 - 2.5^\circ$ in solar longitude (Ryabova 2007). With YR-addition the width increases by 0.3° , remaining less than the observed width 5° (Fox *et al.* 1983). As to the semi-major axis, YR-effect increases Δa for the model *shower* nearly 3 times. We found that YR-effect, being the mass-dependant, increases the mass separation in the stream about twice.

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† Part of this research was firstly fulfilled about 25 years ago. It remained unnoticed to meteor astronomers community, because the English translation (Ryabova 1990) of the paper was published in a hard-to-reach journal and, until recently, was absent from SAO/NASA ADS database. The method was carefully revised and some corrections were introduced.