

SYMMETRY BREAKING FOR GROUND-STATE SOLUTIONS OF HÉNON SYSTEMS IN A BALL

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Abstract. We consider in this paper the problem

$$\begin{cases} -\Delta u = |x|^\alpha v^p, & x \in \Omega, \\ -\Delta v = |x|^\beta u^{q_\varepsilon}, & x \in \Omega, \\ u > 0, \quad v > 0, & x \in \Omega, \\ u = v = 0, & x \in \partial\Omega, \end{cases} \quad (1)$$

where Ω is the unit ball in \mathbb{R}^N centred at the origin, $0 \leq \alpha < pN$, $\beta > 0$, $N \geq 3$. Suppose $q_\varepsilon \rightarrow q$ as $\varepsilon \rightarrow 0^+$ and q_ε, q satisfy, respectively,

$$\frac{N}{p+1} + \frac{N}{q_\varepsilon+1} > N-2, \quad \frac{N}{p+1} + \frac{N}{q+1} = N-2;$$

we investigate the asymptotic estimates of the ground-state solutions $(u_\varepsilon, v_\varepsilon)$ of (1) as $\beta \rightarrow +\infty$ with p, q_ε fixed. We also show the symmetry-breaking phenomenon with α, β fixed and $q_\varepsilon \rightarrow q$ as $\varepsilon \rightarrow 0^+$. In addition, the ground-state solution is non-radial provided that $\varepsilon > 0$ is small or β is large enough.

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1. Introduction. In this paper, we investigate the limiting behaviour of the ground-state solutions of the problem

$$\begin{cases} -\Delta u = |x|^\alpha v^p, & x \in \Omega, \\ -\Delta v = |x|^\beta u^{q_\varepsilon}, & x \in \Omega, \\ u = v = 0, & x \in \partial\Omega, \end{cases} \quad (2)$$

where Ω is the unit ball in \mathbb{R}^N centred at the origin, $0 \leq \alpha < pN$, $\beta > 0$, $N \geq 3$. We assume in this paper that $q_\varepsilon \rightarrow q$ as $\varepsilon \rightarrow 0^+$ and q_ε, q satisfy, respectively,

$$\frac{N}{p+1} + \frac{N}{q_\varepsilon+1} > N-2, \quad \frac{N}{p+1} + \frac{N}{q+1} = N-2.$$

Problem (2) has two features. First, it is a Hénon-type system. The Hénon equation with Dirichlet boundary conditions

$$\begin{cases} -\Delta u = |x|^\alpha u^p, & x \in \Omega, \\ u = 0, & x \in \partial\Omega \end{cases} \quad (3)$$

was found in [10], which stems from rotating stellar structures. A standard compactness argument show that the infimum

$$\inf_{u \in H_0^1(\Omega) \setminus \{0\}} \frac{\int_{\Omega} |\nabla u|^2 \, dx}{\left(\int_{\Omega} |x|^\alpha |u|^{p+1} \, dx\right)^{\frac{2}{p+1}}} \tag{4}$$

is achieved for any $1 < p < 2^* - 1$, $\alpha > 0$. In 1982, Ni [14] proved that the infimum

$$\inf_{u \in H_{0,\text{rad}}^1(\Omega) \setminus \{0\}} \frac{\int_{\Omega} |\nabla u|^2 \, dx}{\left(\int_{\Omega} |x|^\alpha |u|^{p+1} \, dx\right)^{\frac{2}{p+1}}} \tag{5}$$

is achieved for any $p \in (1, \frac{N+2+2\alpha}{N-2})$ by a function in $H_{0,\text{rad}}^1(\Omega)$, the space of radial $H_0^1(\Omega)$ functions. Thus, radial solutions of (3) exist also for (Sobolev) supercritical exponents p . A natural question is whether any minimizer of (4) must be radially symmetric in the range $1 < p < \frac{N+2}{N-2}$ and $\alpha > 0$. Since the weight $|\cdot|^\alpha$ is an increasing function, neither rearrangement arguments nor the moving plane techniques of [7] can be applied.

For $\alpha > 0$, Smets et al. proved in [15] some symmetry-breaking results for (3). They proved that the minimizers of (4) (the so-called ground-state solutions, or least energy solutions) cannot be radial for α large enough. As a consequence, (3) has at least two solutions when α is sufficiently large (see also [16]).

Quite recently, Cao and Peng [3] proved that for $p + 1$ sufficiently close to 2^* , the ground-state solutions of (3) possess a unique maximum point whose distance from $\partial\Omega$ tends to zero as $p \rightarrow \frac{N+2}{N-2}$.

For more results about symmetry breaking phenomena for solutions of problem (3) either α is large enough or $p \rightarrow \frac{N+2}{N-2}$, see for instance [2, 1] and references therein.

Second, the system in (2) is a Hamiltonian-type system, which is strongly indefinite. The existence of solutions of the Hamiltonian elliptic system

$$\begin{cases} -\Delta u = v^p, & x \in \Omega, \\ -\Delta v = u^q, & x \in \Omega, \\ u = v = 0, & x \in \partial\Omega \end{cases} \tag{6}$$

was first considered in [5] and [11] with $\frac{1}{p+1} + \frac{1}{q+1} > \frac{N-2}{N}$; the curve of $(p, q) \in \mathbb{R}^2$ satisfying $\frac{1}{p+1} + \frac{1}{q+1} = \frac{N-2}{N}$ is called critical hyperbola. Afterwards, various results were obtained in the literature. Extensions of problem (6) can be found in [6] and [13]. In [13], existence problems for Hardy-type systems and Hénon-type systems were established. Particularly, for Hénon-type systems

$$\begin{cases} -\Delta u = |x|^\alpha v^p, & x \in \Omega, \\ -\Delta v = |x|^\beta u^q, & x \in \Omega, \\ u = v = 0, & x \in \partial\Omega, \end{cases} \tag{7}$$

the critical hyperbola is $\frac{1}{p+1}(1 + \frac{\alpha}{N}) + \frac{1}{q+1}(1 + \frac{\beta}{N}) = \frac{N-2}{N}$.

In recent years, a study of the limiting behaviour of ground-state solutions of elliptic problems has attracted considerable attention. For the system (6), the limiting behaviour of solutions of (6) as $\frac{1}{p+1} + \frac{1}{q+1} \rightarrow \frac{N-2}{N}$ was discussed in [8]. For the system (7), Yang and He [9] proved that for $\frac{1}{p+1} + \frac{1}{q+1} \rightarrow \frac{N-2}{N}$, the ground-state solutions of (7) possess a unique maximum point whose distance from $\partial\Omega$ tends to zero as

$\frac{1}{p+1} + \frac{1}{q+1} \rightarrow \frac{N-2}{N}$. Such problems are closely related to solutions of the following problem:

$$\begin{cases} -\Delta U = V^p, & y \in \mathbb{R}^N, \\ -\Delta V = U^q, & y \in \mathbb{R}^N, \\ U(y) > 0, \quad V(y) > 0, & y \in \mathbb{R}^N, \\ U(0) = 1, \quad U \rightarrow 0, \quad V \rightarrow 0 \text{ as } |y| \rightarrow \infty, \end{cases} \tag{8}$$

where $\frac{1}{p+1} + \frac{1}{q+1} = \frac{N-2}{N}$. It was proved in [12] that $U \in D^{2, \frac{p+1}{p}}(\mathbb{R}^N)$, $V \in D^{2, \frac{q+1}{q}}(\mathbb{R}^N)$, where $D^{2,r}(\mathbb{R}^N)$ is the completion of $C_0^\infty(\mathbb{R}^N)$ with respect to the norm $\|\Delta \cdot\|_r$. Actually, U and V are radially symmetric for $p \geq 1$ as showed in [4]. Moreover, U and V are unique and decreasing in r . In the discussion of one equation problem, one uses the instanton for the best Sobolev constant. However, no explicit form of (U, V) was found for $p > \frac{2}{N-2}$ up to now. Instead, the asymptotic behaviour of (U, V) as $r \rightarrow \infty$ is sufficient for this purpose. It was found in [12] that

$$\lim_{r \rightarrow \infty} r^{N-2} V(r) = a, \quad \begin{cases} \lim_{r \rightarrow \infty} r^{N-2} U(r) = b & \text{if } p > \frac{N}{N-2}, \\ \lim_{r \rightarrow \infty} \frac{r^{N-2}}{\log r} U(r) = b & \text{if } p = \frac{N}{N-2}, \\ \lim_{r \rightarrow \infty} r^{p(N-2)-2} U(r) = b & \text{if } \frac{2}{N-2} < p < \frac{N}{N-2}, \end{cases} \tag{9}$$

and

$$\lim_{r \rightarrow \infty} \frac{rV'(r)}{V(r)} = 2 - N, \quad \begin{cases} \lim_{r \rightarrow \infty} \frac{rU'(r)}{U(r)} = 2 - N & \text{if } p \geq \frac{N}{N-2}, \\ \lim_{r \rightarrow \infty} \frac{rU'(r)}{U(r)} = 2 - p(N - 2) & \text{if } p \leq \frac{N}{N-2}. \end{cases} \tag{10}$$

In this paper, we are interested in the symmetry of ground-states solutions of (2). Now, we denote

$$E_\alpha(\Omega) = \left\{ u \in W^{2, \frac{p+1}{p}} \cap W_0^{1, \frac{p+1}{p}}(\Omega) : \int_\Omega |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx < \infty \right\}$$

and

$$E_\alpha^{\text{rad}}(\Omega) = \{u \in E_\alpha(\Omega) : u(x) = u(|x|)\}.$$

Our main results are as follows.

THEOREM 1.1. *Suppose $N \geq 3$, $0 \leq \alpha < pN$, $\beta > 0$, $p > \frac{2}{N-2}$, $q_\varepsilon > \frac{N+p}{Np-2p-1}$; then there exists $\beta^* > 0$ such that the ground-state solutions $u_{\alpha,\beta,\varepsilon}$ are non-radial provided $\beta > \beta^*$.*

THEOREM 1.2. *Suppose $N \geq 3$, $0 \leq \alpha < pN$, $\beta > 0$, $p > \frac{2}{N-2}$, $pq_\varepsilon > 1$; then there exists $\varepsilon^* > 0$ such that the ground-state solutions $u_{\alpha,\beta,\varepsilon}$ are non-radial provided $\varepsilon < \varepsilon^*$ or $q - q_\varepsilon < \varepsilon^*$.*

This paper is organized as follows. In section 2, we give some preliminaries which turn out to be essential. In section 3, we present some estimates for radial ground-state solutions of (2) with α, p, q_ε fixed and $\beta \rightarrow \infty$. This will lead us to get the first

symmetry-breaking result, stating that for β sufficiently large, the ground-state solution of problem (2) is non-radial. In section 4, another symmetry-breaking result is proved, with α, β, p fixed and $q_\varepsilon \rightarrow q$ as $\varepsilon \rightarrow 0$.

2. Preliminaries. Before proving our main results, we want to introduce some simple calculus lemma which turns out to be essential:

LEMMA 2.1. *Let u be a radially symmetric function of Ω (unit ball in \mathbb{R}^N) with $u(1) = 0, u'(0)$ exists. Then*

$$(i) |u(x)| \leq \frac{1}{w_{N-1}^{\frac{p}{p+1}}(p(N-1)-1)^{\frac{1}{p+1}} \left(|x|^{p(N-1)-1}\right)^{\frac{1}{p+1}}} \left(\int_{\Omega} |\nabla u|^{\frac{p+1}{p}} dy\right)^{\frac{p}{p+1}}$$

$$(ii) \left|\frac{\partial u}{\partial r}\right| \leq r^{1-N} \left(\frac{r^{N+\alpha}}{N+\alpha}\right)^{\frac{1}{p+1}} \left(\frac{1}{w_{N-1}} \int_{B_r(0)} |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx\right)^{\frac{p}{p+1}},$$

where w_{N-1} is the surface area of the unit ball in \mathbb{R}^N .

Proof. (i) For

$$u(1) - u(x) = \int_{|x|}^1 |u'(t)| dt,$$

$$|u(x)| \leq \int_{|x|}^1 |u'(t)| dt$$

$$\leq \left(\int_{|x|}^1 |u'(t)|^{\frac{p+1}{p}} t^{N-1} dt\right)^{\frac{p}{p+1}} \left(\int_{|x|}^1 t^{-(N-1)p} dt\right)^{\frac{1}{p+1}}.$$

Since

$$\int_{|x|}^1 t^{-(N-1)p} dt = \frac{1}{p(N-1)-1} \left(\frac{1}{|x|^{p(N-1)-1}} - 1\right) \leq \frac{1}{p(N-1)-1} \frac{1}{|x|^{p(N-1)-1}}$$

and

$$\int_{|x|}^1 |u'(t)|^{\frac{p+1}{p}} t^{N-1} dt = \frac{1}{w_{N-1}} \int \left(\int_{|x|}^1 |u'(t)|^{\frac{p+1}{p}} t^{N-1} dt\right) w(\theta) d\theta$$

$$= \frac{1}{w_{N-1}} \int_{|x| \leq |y| \leq 1} |\nabla u|^{\frac{p+1}{p}} dy$$

$$\leq \frac{1}{w_{N-1}} \int_{\Omega} |\nabla u|^{\frac{p+1}{p}} dy,$$

we find that

$$|u(x)| \leq \frac{1}{w_{N-1}^{\frac{p}{p+1}}(p(N-1)-1)^{\frac{1}{p+1}} \left(|x|^{p(N-1)-1}\right)^{\frac{1}{p+1}}} \left(\int_{\Omega} |\nabla u|^{\frac{p+1}{p}} dy\right)^{\frac{p}{p+1}}.$$

(ii) For

$$\begin{aligned} |r^{N-1} \frac{\partial u}{\partial r}| &\leq \int_0^r t^{N-1} |\Delta u| dt \\ &\leq \left(\int_0^r t^{\frac{(N-1+\alpha)(p+1)}{p+1}} dt \right)^{\frac{1}{p+1}} \left(\int_0^r (t^{\frac{p(N-1)-\alpha}{p+1}} |\Delta u|)^{\frac{p+1}{p}} dt \right)^{\frac{p}{p+1}} \\ &= \left(\frac{r^{N+\alpha}}{N+\alpha} \right)^{\frac{1}{p+1}} \left(\frac{1}{w_{N-1}} \int_{B_r(0)} |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx \right)^{\frac{p}{p+1}}, \end{aligned}$$

we have

$$\left| \frac{\partial u}{\partial r} \right| \leq r^{1-N} \left(\frac{r^{N+\alpha}}{N+\alpha} \right)^{\frac{1}{p+1}} \left(\frac{1}{w_{N-1}} \int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx \right)^{\frac{p}{p+1}}.$$

□

As a result of Lemma 2.1, we obtain the following corollary.

COROLLARY 2.1. *Under hypothesis of Lemma 2.1, and if $\frac{N+\beta}{q+1} + \frac{N+\alpha}{p+1} > N - 2$, then*

$$\int_{\Omega} |x|^{\beta} |u(x)|^{q+1} dx \leq C \left(\int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx \right)^{\frac{p(q+1)}{p+1}},$$

where

$$C = \left(\frac{1}{w_{N-1}} \right)^{\frac{p(q+1)}{p+1}} \frac{1}{N+\beta - \left(N-2 - \frac{N+\alpha}{p+1} \right) (q+1)} \left(\frac{1}{N-2 - \frac{N+\alpha}{p+1}} \right)^{q+1} \left(\frac{1}{N+\alpha} \right)^{\frac{q+1}{p+1}}.$$

Proof. From the above Lemma 2.1, we have

$$|u(|x|)| \leq \int_{|x|}^1 |u'(r)| dr \leq \int_{|x|}^1 \left(\frac{r^{N+\alpha}}{N+\alpha} \right)^{\frac{1}{p+1}} r^{1-N} dr \left(\frac{1}{w_{N-1}} \int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx \right)^{\frac{p}{p+1}}.$$

Since

$$\begin{aligned} \int_{|x|}^1 \left(\frac{r^{N+\alpha}}{N+\alpha} \right)^{\frac{1}{p+1}} r^{1-N} dr &= \left(\frac{1}{N+\alpha} \right)^{\frac{1}{p+1}} \int_{|x|}^1 r^{\frac{N+\alpha}{p+1}} r^{1-N} dr \\ &\leq \frac{1}{N-2 - \frac{N+\alpha}{p+1}} \left(\frac{1}{N+\alpha} \right)^{\frac{1}{p+1}} \frac{1}{|x|^{N-2 - \frac{N+\alpha}{p+1}}}, \end{aligned}$$

we have

$$\begin{aligned} &r^{\beta} r^{N-1} |u(r)|^{q+1} \\ &\leq r^{\beta} r^{N-1} \left(\frac{1}{N-2 - \frac{N+\alpha}{p+1}} \left(\frac{1}{N+\alpha} \right)^{\frac{1}{p+1}} \frac{1}{r^{N-2 - \frac{N+\alpha}{p+1}}} \right)^{q+1} \left(\frac{1}{w_{N-1}} \int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx \right)^{\frac{p(q+1)}{p+1}}, \end{aligned}$$

which implies

$$\int_0^1 r^\beta r^{N-1} |u(r)|^{q+1} dr \leq \left(\frac{1}{N-2-\frac{N+\alpha}{p+1}} \left(\frac{1}{N+\alpha} \right)^{\frac{1}{p+1}} \right)^{q+1} \times \int_0^1 \left(\frac{1}{r^{N-2-\frac{N+\alpha}{p+1}}} \right)^{q+1} r^\beta r^{N-1} dr \left(\frac{1}{w_{N-1}} \int_\Omega |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx \right)^{\frac{p(q+1)}{p+1}}.$$

Thus, the conclusion holds. □

3. Asymptotic estimates. Consider the minimization problem

$$S_{\alpha, \beta, \varepsilon}^{\text{rad}} = \inf_{u \in E_\alpha^{\text{rad}}(\Omega) \setminus \{0\}} R_{\alpha, \beta, \varepsilon}(u), \tag{11}$$

where

$$R_{\alpha, \beta, \varepsilon}(u) = \frac{\int_\Omega |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx}{\left(\int_\Omega |x|^\beta |u|^{q_\varepsilon+1} dx \right)^{\frac{p+1}{p(q_\varepsilon+1)}}, \quad u \in E_\alpha(\Omega) \setminus \{0\}, \tag{12}$$

is the Rayleigh quotient associated with (2). Similar to [14], we can also prove that

$$S_{\alpha, \beta, \varepsilon}^{\text{rad}}(\Omega) = \inf_{u \in E_\alpha^{\text{rad}}(\Omega) \setminus \{0\}} \frac{\int_\Omega |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx}{\left(\int_\Omega |x|^\beta |u|^{q_\varepsilon+1} dx \right)^{\frac{p+1}{p(q_\varepsilon+1)}}$$

is attained by some positive function $u_{\alpha, \beta, \varepsilon}^{\text{rad}}$. After scaling, $u_{\alpha, \beta, \varepsilon}^{\text{rad}}$ is also a solution of (2).

Now, we provide an estimate of the energy $S_{\alpha, \beta, \varepsilon}^{\text{rad}}$ as $\beta \rightarrow \infty$.

LEMMA 3.1. *If $N \geq 3$, there exists $C > 0$ depending on N, p such that*

$$S_{\alpha, \beta, \varepsilon}^{\text{rad}} \geq C\beta^{\frac{p+2+q_\varepsilon}{p(q_\varepsilon+1)}} \text{ as } \beta \rightarrow \infty.$$

Proof. Let $u \in E_\alpha^{\text{rad}}(\Omega)$ and define the rescaled function $v(|x|) = u(|x|^s)$, where $s = \frac{N}{\beta+N}$. Then

$$\int_\Omega |x|^\beta |u|^{q_\varepsilon+1} dx = w_{N-1} \int_0^1 r^{\beta+N-1} |u(r)|^{q_\varepsilon+1} dr = s \int_\Omega |x|^{s(\beta+N)-N} |v(x)|^{q_\varepsilon+1} dx,$$

and from Lemma 2.1, we have

$$\int_\Omega |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx \geq w_{N-1}^{\frac{p+1}{p}} (N+\alpha)^{\frac{1}{p}} r^{\frac{(p+1)(N-1)-(N+\alpha)}{p}} \left| \frac{\partial u}{\partial r} \right|^{\frac{p+1}{p}},$$

which implies that

$$\begin{aligned} \int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx &\geq w_{N-1}^{\frac{p+1}{p}} (N + \alpha)^{\frac{1}{p}} \int_0^1 r^{\frac{(p+1)(N-1)-(N+\alpha)}{p}} \left| \frac{\partial u}{\partial r} \right|^{\frac{p+1}{p}} dr \\ &= w_{N-1}^{\frac{p+1}{p}} (N + \alpha)^{\frac{1}{p}} \int_0^1 r^{\frac{(p+1)(N-1)-(N+\alpha)}{p}} s^{-\frac{p+1}{p}} r^{\frac{1-s}{s}} \frac{p+1}{p} \left| \frac{\partial v}{\partial t} \right|^{\frac{p+1}{p}} dr \\ &= w_{N-1}^{\frac{1}{p}} (N + \alpha)^{\frac{1}{p}} s^{-\frac{1}{p}} \int_{\Omega} |x|^{\frac{p(N-1)(s-1)-s(1+\alpha)+(1-s)}{p}} |\nabla v|^{\frac{p+1}{p}} dx. \end{aligned}$$

Thus, we obtain

$$\frac{\int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx}{\left(\int_{\Omega} |x|^{\beta} |u|^{q_{\varepsilon}+1} dx\right)^{\frac{p+1}{p(q_{\varepsilon}+1)}}} \geq \frac{w_{N-1}^{\frac{1}{p}} (N + \alpha)^{\frac{1}{p}} s^{-\frac{1}{p}} \int_{\Omega} |x|^{\frac{p(N-1)(s-1)-s(1+\alpha)+(1-s)}{p}} |\nabla v|^{\frac{p+1}{p}} dx}{\left(s \int_{\Omega} |v(x)|^{q_{\varepsilon}+1} dx\right)^{\frac{p+1}{p(q_{\varepsilon}+1)}}}.$$

It follows that

$$S_{\alpha, \beta, \varepsilon}^{\text{rad}} \geq w_{N-1}^{\frac{1}{p}} (N + \alpha)^{\frac{1}{p}} s^{-\frac{1}{p}} s^{-\frac{p+1}{p(q_{\varepsilon}+1)}} \inf_{v \in W_{0, \text{rad}}^{1, \frac{p+1}{p}}(\Omega)} \frac{\int_{\Omega} |x|^{\frac{p(N-1)(s-1)-s(1+\alpha)+(1-s)}{p}} |\nabla v|^{\frac{p+1}{p}} dx}{\left(\int_{\Omega} |v(x)|^{q_{\varepsilon}+1} dx\right)^{\frac{p+1}{p(q_{\varepsilon}+1)}}}.$$

Now, we claim that for every $0 \leq s \leq 1$, we have $p(N - 1)(s - 1) - s(1 + \alpha) + (1 - s) < 0$. Indeed, for

$$\begin{aligned} &p(N - 1)(s - 1) - s(1 + \alpha) + (1 - s) \\ &= (p(N - 1) - 2 - \alpha)s - (p(N - 1) - 1), \end{aligned}$$

if $p(N - 1) - 2 \leq \alpha < pN$, we have $p(N - 1)(s - 1) - s(1 + \alpha) + (1 - s) < 0$; if $\alpha < p(N - 1) - 2 < p(N - 1) - 1$, then for every $0 \leq s \leq 1$, we also have $p(N - 1)(s - 1) - s(1 + \alpha) + (1 - s) < 0$. Therefore,

$$\int_{\Omega} |\nabla v|^{\frac{p+1}{p}} dx \leq \int_{\Omega} |x|^{\frac{p(N-1)(s-1)-s(1+\alpha)+(1-s)}{p}} |\nabla v|^{\frac{p+1}{p}} dx,$$

which implies that

$$c_s = \inf_{v \in W_{0, \text{rad}}^{1, \frac{p+1}{p}}(\Omega)} \frac{\int_{\Omega} |x|^{\frac{p(N-1)(s-1)-s(1+\alpha)+(1-s)}{p}} |\nabla v|^{\frac{p+1}{p}} dx}{\left(\int_{\Omega} |v(x)|^{q_{\varepsilon}+1} dx\right)^{\frac{p+1}{p(q_{\varepsilon}+1)}}}$$

is achieved by standard arguments. Since $|x| \leq 1$, if $p(N - 1) - 2 \leq \alpha < pN$, $p(N - 1)(s - 1) - s(1 + \alpha) + (1 - s) \leq -(p(N - 1) - 1)$, which implies that $c_s \geq c_0$; if $\alpha < p(N - 1) - 2$, $p(N - 1)(s - 1) - s(1 + \alpha) + (1 - s) = (p(N - 1) - 2 - \alpha)s - (p(N - 1) - 1)$, then for every $0 \leq s \leq 1$, $|x|^{\frac{p(N-1)(s-1)-s(1+\alpha)+(1-s)}{p}}$ is non-increasing, which implies that c_s is non-increasing on $[0, 1]$; then $c_s \geq c_1$. Thus,

$$S_{\alpha, \beta, \varepsilon}^{\text{rad}} \geq C(N + \alpha)^{\frac{1}{p}} s^{-\frac{p+2+q_{\varepsilon}}{p(q_{\varepsilon}+1)}}, \quad \beta \rightarrow \infty.$$

□

By the assumptions on p, q_ε , the inclusion $W^{2, \frac{p+1}{p}}(\Omega) \hookrightarrow L^{q_\varepsilon+1}(\Omega)$ is compact. It implies that

$$S_{\alpha, \beta, \varepsilon}(\Omega) = \inf_{u \in E_\alpha(\Omega) \setminus \{0\}} \frac{\int_\Omega |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx}{\left(\int_\Omega |x|^\beta |u|^{q_\varepsilon+1} dx\right)^{\frac{p+1}{p(q_\varepsilon+1)}}$$

is attained by some positive function $u_{\alpha, \beta, \varepsilon}$. After scaling, $\bar{u}_{\alpha, \beta, \varepsilon}$ is a solution of (2).

LEMMA 3.2. Assume $N \geq 3$, for any p, q_ε satisfy $\frac{N}{p+1} + \frac{N}{q_\varepsilon+1} > N - 2$ with $p > \frac{2}{N-2}$, $q_\varepsilon > \frac{N+p}{Np-2p-1}$; there exists $\beta^* \geq 0$ such that $S_{\alpha, \beta, \varepsilon} < S_{\alpha, \beta, \varepsilon}^{\text{rad}}$ provided $\beta > \beta^*$.

Proof. For any fixed $u \in C_0^\infty(\Omega)$, define $u_\beta(x) = u(\beta(x - x_\beta))$, where $x_\beta = (1 - \frac{1}{\beta}, 0, \dots, 0)$. For $|\beta(x - x_\beta)| \leq 1$, that is, $|x - x_\beta| \leq \frac{1}{\beta}$, then $|x| \geq |x_\beta| - \frac{1}{\beta} = 1 - \frac{2}{\beta}$. One has

$$\int_\Omega |x|^{-\frac{\alpha}{p}} |\Delta u_\beta|^{\frac{p+1}{p}} dx \leq \left(1 - \frac{2}{\beta}\right)^{-\frac{\alpha}{p}} \beta^{\frac{2(p+1)}{p} - N} \int_\Omega |\Delta u|^{\frac{p+1}{p}} dx,$$

and

$$\int_\Omega |x|^\beta |u_\beta|^{q_\varepsilon+1} dx \geq \left(1 - \frac{2}{\beta}\right)^\beta \beta^{-N} \int_\Omega |u|^{q_\varepsilon+1} dx.$$

Hence by definition, one obtains

$$\begin{aligned} S_{\alpha, \beta, \varepsilon} &\leq \frac{\int_\Omega |x|^{-\frac{\alpha}{p}} |\Delta u_\beta|^{\frac{p+1}{p}} dx}{\left(\int_\Omega |x|^\beta |u_\beta|^{q_\varepsilon+1} dx\right)^{\frac{p+1}{p(q_\varepsilon+1)}}} \\ &\leq \frac{\left(1 - \frac{2}{\beta}\right)^{-\frac{\alpha}{p}} \beta^{\frac{2(p+1)}{p} - N} \int_\Omega |\Delta u|^{\frac{p+1}{p}} dx}{\left(\left(1 - \frac{2}{\beta}\right)^\beta \beta^{-N} \int_\Omega |u|^{q_\varepsilon+1} dx\right)^{\frac{p+1}{p(q_\varepsilon+1)}}} \\ &\leq C_1 \beta^{\frac{2(p+1)}{p} - N + \frac{N(p+1)}{p(q_\varepsilon+1)}} \frac{\int_\Omega |\Delta u|^{\frac{p+1}{p}} dx}{\left(\int_\Omega |u|^{q_\varepsilon+1} dx\right)^{\frac{p+1}{p(q_\varepsilon+1)}}}. \end{aligned}$$

Since u is fixed and $\frac{\int_\Omega |\Delta u|^{\frac{p+1}{p}} dx}{\left(\int_\Omega |u|^{q_\varepsilon+1} dx\right)^{\frac{p+1}{p(q_\varepsilon+1)}}}$ is independent of β , we have

$$S_{\alpha, \beta, \varepsilon} \leq C \beta^{\frac{2(p+1)}{p} - N + \frac{N(p+1)}{p(q_\varepsilon+1)}}.$$

From Lemma 3.1 $S_{\alpha, \beta, \varepsilon}^{\text{rad}} \geq C \beta^{\frac{p+2+q_\varepsilon}{p(q_\varepsilon+1)}}$ as $\beta \rightarrow \infty$, and $\frac{p+2+q_\varepsilon}{p(q_\varepsilon+1)} > \frac{2(p+1)}{p} - N + \frac{N(p+1)}{p(q_\varepsilon+1)}$, that is, $q_\varepsilon > \frac{N+p}{Np-2p-1}$. Hence $S_{\alpha, \beta, \varepsilon} < S_{\alpha, \beta, \varepsilon}^{\text{rad}}$ as $\beta \rightarrow \infty$. □

4. Analysis for ε close to 0. In this section, we analyse the case where ε is close to 0, that is, q_ε is close to q . We will show that for any fixed $0 \leq \alpha < pN$, $\beta > 0$, the minimizer of $R_{\alpha, \beta, \varepsilon}$ is non-radial provided that ε is sufficiently small.

LEMMA 4.1. *If $N \geq 3$, there exists $c_0 > 0$, such that for every q_ε and for every $0 \leq \alpha < pN$, $\beta > 0$,*

$$c_0 \beta^{\frac{p+1}{p(q_\varepsilon+1)}} \leq S_{\alpha,\beta,\varepsilon}^{rad}.$$

Proof. From Lemma 2.1, we obtain

$$\int_{\Omega} |x|^\beta |u(x)|^{q_\varepsilon+1} dx \leq C \left(\int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta u|^{\frac{p+1}{p}} dx \right)^{\frac{p(q_\varepsilon+1)}{p+1}},$$

where

$$C = \left(\frac{1}{w_{N-1}} \right)^{\frac{p(q_\varepsilon+1)}{p+1}} \frac{1}{N + \beta - (N - 2 - \frac{N+\alpha}{p+1})(q_\varepsilon + 1)} \left(\frac{1}{N - 2 - \frac{N+\alpha}{p+1}} \right)^{q_\varepsilon+1} \left(\frac{1}{N + \alpha} \right)^{\frac{q_\varepsilon+1}{p+1}}.$$

Since $u \in E_\alpha^{rad}(\Omega)$ is arbitrary,

$$c_0 \left(N + \beta - (N - 2 - \frac{N + \alpha}{p + 1})(q_\varepsilon + 1) \right)^{\frac{(p+1)}{p(q_\varepsilon+1)}} \leq S_{\alpha,\beta,\varepsilon}^{rad},$$

which ends the proof. □

Let us denote by S the classical Sobolev constant

$$S = \inf_{u \in W^{2, \frac{p+1}{p}} \cap W_0^{1, \frac{p+1}{p}}(\Omega)} \frac{\int_{\Omega} |\Delta u|^{\frac{p+1}{p}} dx}{\left(\int_{\Omega} |u|^{q+1} dx \right)^{\frac{p+1}{p(q+1)}}}.$$

It is standard that this Rayleigh quotient is invariant under translations and dilations.

LEMMA 4.2. *If $N \geq 3$ and $0 \leq \alpha < pN$, $\beta > 0$, then*

$$S = S_{\alpha,\beta,0} < S_{\alpha,\beta,0}^{rad}.$$

Proof. Using Corollary 2.1, it is easy to verify that $S_{\alpha,\beta,0}^{rad}$ is achieved, so that $S < S_{\alpha,\beta,0}^{rad}$. Now, we claim that $S = S_{\alpha,\beta,0}$. From the definition of $S_{\alpha,\beta,0}$, we know that $S \leq S_{\alpha,\beta,0}$. Thus, we will prove that $S_{\alpha,\beta,0} \leq S$. Indeed, for $\delta > 0$, we can choose $x_\delta = (1 - \frac{1}{|\ln \delta|}, 0, \dots, 0)$, $U_\delta(x) = U(\frac{x-x_\delta}{\delta})$ and $V_\delta(x) = V(\frac{x-x_\delta}{\delta})$, where (U, V) is the solution of (8). Let $\varphi_\delta \in C_0^\infty(\mathbb{R}^N)$ be a cut-off function satisfying

$$\varphi_\delta(x) = \begin{cases} 1, & x \in B(x_\delta, \frac{1}{2|\ln \delta|}), \\ 0, & x \in \mathbb{R}^N \setminus B(x_\delta, \frac{1}{|\ln \delta|}), \end{cases}$$

$0 \leq \varphi_\delta(x) \leq 1$, $|\nabla \varphi_\delta(x)| \leq C |\ln \delta|$, $|\Delta \varphi_\delta| \leq C |\ln \delta|^2$ in \mathbb{R}^N , where $C > 0$, is independent of δ . Set $w_\delta = \varphi_\delta U_\delta$, similar to [9], we have

$$\lim_{\delta \rightarrow 0} \lim_{\varepsilon \rightarrow 0} \frac{\int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta w_\delta|^{\frac{p+1}{p}} dx}{\left(\int_{\Omega} |x|^\beta |w_\delta|^{q_\varepsilon+1} dx \right)^{\frac{1}{q_\varepsilon+1} \frac{p+1}{p}}} \leq S,$$

and

$$\begin{aligned} & \lim_{\varepsilon \rightarrow 0} \frac{\int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta w_{\delta}|^{\frac{p+1}{p}} dx}{\left(\int_{\Omega} |x|^{\beta} |w_{\delta}|^{q_{\varepsilon}+1} dx\right)^{\frac{1}{q_{\varepsilon}+1}} \frac{p+1}{p}} \\ & \geq \lim_{\varepsilon \rightarrow 0} \frac{\int_{\Omega} |x|^{-\frac{\alpha}{p}} |\Delta u_{\alpha, \beta, \varepsilon}|^{\frac{p+1}{p}} dx}{\left(\int_{\Omega} |x|^{\beta} |u_{\alpha, \beta, \varepsilon}|^{q_{\varepsilon}+1} dx\right)^{\frac{1}{q_{\varepsilon}+1}} \frac{p+1}{p}} \\ & \geq \lim_{\varepsilon \rightarrow 0} \frac{\int_{\Omega} |\Delta u_{\alpha, \beta, \varepsilon}|^{\frac{p+1}{p}} dx}{\left(\int_{\Omega} |u_{\alpha, \beta, \varepsilon}|^{q+1} dx\right)^{\frac{1}{q+1}} \frac{p+1}{p}}, \end{aligned}$$

where $u_{\alpha, \beta, \varepsilon}$ is a minimizer of $S_{\alpha, \beta, \varepsilon}$; thus $S_{\alpha, \beta, 0} \leq S$. □

LEMMA 4.3. *Assume that $N \geq 3$. For any $n \in \mathbb{N}$ there exists $\delta_n > 0$ such that $S_{\alpha, \beta, \varepsilon} < S_{\alpha, \beta, \varepsilon}^{\text{rad}}$ provided $\beta \geq \frac{1}{n}$ and $\varepsilon < \delta_n$.*

Proof. By contradiction, assume that there exists $n \in \mathbb{N}$ and sequences $\beta_k \geq \frac{1}{n}$ and $\delta_k \rightarrow 0$ such that

$$S_{\alpha, \beta_k, \delta_k} = S_{\alpha, \beta_k, \delta_k}^{\text{rad}}. \tag{13}$$

From Lemma 3.2, there exists c_1 independent of q_{δ_k} , such that

$$S_{\alpha, \beta, \delta_k} \leq c_1 \beta_k^{\frac{2(p+1)}{p} - N + \frac{N(p+1)}{p(q_{\delta_k}+1)}}.$$

Lemma 4.1 implies that

$$c_0 \beta_k^{\frac{p+1}{p(q_{\delta_k}+1)}} \leq S_{\alpha, \beta, \delta_k}^{\text{rad}}.$$

Since

$$\frac{p+1}{p(q_{\delta_k}+1)} - \frac{2(p+1)}{p} + N - \frac{N(p+1)}{p(q_{\delta_k}+1)} = \frac{(pN - 2p - 2)(q_{\delta_k} + 1) - (p+1)(N-1)}{p(q_{\delta_k} + 1)},$$

we have

$$\beta_k^{\frac{(pN - 2p - 2)(q_{\delta_k} + 1) - (p+1)(N-1)}{p(q_{\delta_k} + 1)}} \leq \frac{c_1}{c_0}.$$

From $q+1 = \frac{N(p+1)}{Np-2p-2} > \frac{(N-1)(p+1)}{Np-2p-2}$ and $q_{\delta_k} \rightarrow q$ as $k \rightarrow +\infty$, it is implies that $(pN - 2p - 2)(q_{\delta_k} + 1) - (p+1)(N-1) > 0$ as $k \rightarrow +\infty$. Thus, β_k is bounded. We can assume that $\beta_k \rightarrow \beta \geq \frac{1}{n}$ as $k \rightarrow +\infty$.

Claim that

$$S_{\alpha, \beta, 0}^{\text{rad}} = \lim_{k \rightarrow +\infty} S_{\alpha, \beta_k, \delta_k}^{\text{rad}}. \tag{14}$$

Indeed, by upper semi-continuity, it follows that

$$S_{\alpha, \beta, 0}^{\text{rad}} \geq \limsup_{k \rightarrow +\infty} S_{\alpha, \beta_k, \delta_k}^{\text{rad}},$$

On the other hand, from

$$\int_{\Omega} |x|^{\beta_k} |u_k|^{q_{\delta_k}+1} dx \leq \left(\int_{\Omega} |x|^{\beta_k} |u_k|^{q+1} dx \right)^{\frac{q_{\delta_k}+1}{q+1}} \left(\int_{\Omega} |x|^{\beta_k} dx \right)^{\frac{q-q_{\delta_k}}{q+1}},$$

we have

$$S_{\alpha,\beta,0}^{\text{rad}} \leq \liminf_{k \rightarrow +\infty} S_{\alpha,\beta_k,\delta_k}^{\text{rad}}.$$

Similarly, by upper continuity,

$$S_{\alpha,\beta,0} \geq \limsup_{k \rightarrow +\infty} S_{\alpha,\beta_k,\delta_k}. \tag{15}$$

We obtain from (13)–(15), $S_{\alpha,\beta,0} \geq S_{\alpha,\beta,0}^{\text{rad}}$, which contradicts Lemma 4.2. □

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