

X-RAY EMISSION FROM BL Lac OBJECTS: COMPARISON TO THE SYNCHROTRON SELF-COMPTON MODELS

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As one part of our joint study of the X-ray properties of BL Lac objects, we compare the measured X-ray flux densities with those predicted using the synchrotron self-Compton (SSC) formalism (Jones et al. 1974). Naive application of the formalism predicts X-ray fluxes from  $10^{-3}$  to  $10^5$  those observed. We therefore ask what we can learn by simply assuming the SSC mechanism, and looking for ways to reconcile the observed and measured X-ray fluxes. This paper reports our investigation of beaming factors due to relativistic ejection of a radiation source which is isotropic in its own rest frame. We conclude that large Lorentz factors,  $\Gamma \geq 10$ , do not apply to BL Lac objects as a class.

Our present study (Table 1) is based on 16 sources for which a VLBI size has been measured by Weiler and Johnston (1980) and which have been observed (15 detected) in X-rays with the Einstein Imaging Proportional Counter. The X-ray data have been reduced in the manner described by Zamorani et al. (1981), except that we assume a steeper energy spectrum, of index 1.5. We use equation 4 of Jones et al. (1974) to predict the X-ray flux density  $S_x^{\prime}$ (Jy) at 1 keV.

TABLE 1 Distribution of Beaming Factors

Object	$S_x^{\prime}$ (Jy)	$S_x$ (Jy)	$\delta$	Object	$S_x^{\prime}$ (Jy)	$S_x$ (Jy)	$\delta$
0048-097	1.3E-3	2.2E-7	8.8	1101+38	3.2E-7	1.4E-5	0.39
0219+438	8.8E-9	1.5E-7 <sup>1</sup>	0.49	1219+28	3.2E-6	4.8E-7	1.6
0235+164	1.0E-1	1.6E-7	28.0	1308+326	2.1E-1	2.7E-7	30.0
0735+178	4.6E-5	2.4E-7	3.7	1400+162	1.9E-10	1.6E-7 <sup>1</sup>	0.18
0754+100	3.6E-5	2.0E-7	3.6	1538+149	1.9E-5	1.3E-7	3.5
0818-128	3.9E-6	<1.4E-7	>2.3	1652+398	2.1E-7	1.2E-5	0.36
0829+046	1.8E-7	1.9E-7	0.98	2201+04	8.9E-8	2.6E-7	0.76
0851+202	3.4E-4	2.3E-6	3.5	2254+074	8.1E-9	1.2E-7 <sup>1</sup>	0.51

<sup>1</sup>From Maccagni and Tarengi.

We consider a model where the radio source is ejected with a Lorentz factor  $\Gamma$ , at an angle  $\theta$  away from a line toward us. If we define the beaming factor  $\delta = 1/\Gamma(1-\beta\cos\theta)$ , then a Lorentz transformation gives  $S_x = S'_x\delta^{-2(\alpha+2)}$ . In Table 1 we have assumed  $\nu_m = 1$  GHz and  $\alpha = 0$ . Calculating with  $\nu_m = 0.3$  or 5 and  $\alpha = 0.2$  or 0.5 changes  $S'_x$  by a large factor; however, since  $S_x$  depends on  $\delta$  to a power similar to the dependence on the radio size, synchrotron absorption frequency  $\nu_m$  and radio flux density  $S_m$ ,  $\delta$  is determined within a factor of a few.

If the predicted  $S'_x$  were all within a factor of 10 or 100 from the measured  $S_x$ , it could be evidence for a single component SSC model. Instead, we must find a way to reconcile at least the cases where the predicted X-ray flux density greatly exceeds that observed. The assumption of relativistic beaming is sufficient and furthermore is reasonable in the sense that a quasi-isotropic distribution of  $\theta$ , and relatively small  $\Gamma$  are required.

We may predict the intrinsic distribution of  $\delta$  for a set of objects if the Lorentz factors have a probability density function  $\rho(\Gamma)$ :

$$\rho(\delta) = \frac{1}{2\delta^2} \int_{\Gamma_{\min}}^{\infty} \frac{\rho(\Gamma) d\Gamma}{\sqrt{\Gamma^2 - 1}} \quad (1)$$

where the minimum  $\Gamma_{\min} = (1+\delta^2)/2\delta$ . From this equation we can immediately see qualitatively that for a  $\delta$ -function distribution of  $\Gamma$ , values of  $\Gamma_0 = 2$  to 5 would suffice to span most of the range of observed  $\delta$  factors, from 0.1 to 10. A power law  $\rho(\Gamma) \propto \Gamma^{-2}$  would span this same range. On the other hand, for  $\Gamma \geq 10$ , the overwhelming majority of sources would have  $\delta \ll 1$ , which is not what we observe.

Suppose we only recognize an object as a "BL Lac" if it is beamed toward us within some angle  $\theta_m$ . In this case we can only integrate equation (1) up to some  $\Gamma_{\max}$  at which  $\delta \geq 1/\Gamma_{\max}(1-\cos\theta_m\sqrt{1-1/\Gamma_{\max}^2})$ , and the normalization factor 2 changes to  $1-\cos\theta_m$ . For the delta function distribution,  $\theta_m = 1/\Gamma_0$ , and it is easy to see that for  $\Gamma_0 \geq 10$  we do not expect any  $\delta \leq 10$ . We conclude that although some extreme values of  $\Gamma$  may occur for specific sources, the general distribution is dominated by  $\Gamma \lesssim 5$ .

#### REFERENCES

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