



Spectral theory of the invariant Laplacian on the disk and the sphere – a complex analysis approach

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Abstract. The central theme of this paper is the holomorphic spectral theory of the canonical Laplace operator of the complement $\Omega := \{(z, w) \in \widehat{\mathbb{C}}^2 : z \cdot w \neq 1\}$ of the “complexified unit circle” $\{(z, w) \in \widehat{\mathbb{C}}^2 : z \cdot w = 1\}$. We start by singling out a distinguished set of holomorphic eigenfunctions on the bidisk in terms of hypergeometric ${}_2F_1$ functions and prove that they provide a spectral decomposition of every holomorphic eigenfunction on the bidisk. As a second step, we identify the maximal domains of definition of these eigenfunctions and show that these maximal domains naturally determine the fine structure of the eigenspaces. Our main result gives an intrinsic classification of all closed Möbius invariant subspaces of eigenspaces of the canonical Laplacian of Ω . Generalizing foundational prior work of Helgason and Rudin, this provides a unifying complex analytic framework for the real-analytic eigenvalue theories of both the hyperbolic and spherical Laplace operators on the open unit disk resp. the Riemann sphere and, in particular, shows how they are interrelated with one another.

1 Introduction

Let $\widehat{\mathbb{C}} := \mathbb{C} \cup \{\infty\}$ denote the Riemann sphere. The purpose of this paper is to explore the spectral theory of the complex *invariant Laplace operator*

$$\Delta_{zw} = 4(1 - zw)^2 \partial_z \partial_w$$

of the *complement of the complexified unit circle*,

$$(1.1) \quad \Omega := \widehat{\mathbb{C}}^2 \setminus \{(z, w) \in \widehat{\mathbb{C}}^2 : z \cdot w = 1\}^1$$

by function-theoretic methods. This approach allows a unified study of the *real-analytic* spectral theories of the hyperbolic Laplacian

$$\Delta_{\mathbb{D}} := 4(1 - |z|^2)^2 \partial_z \partial_{\bar{z}}$$

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¹Here, we extend the arithmetic in \mathbb{C} in the usual manner by $z \cdot \infty = \infty = \infty \cdot z$ for $z \in \widehat{\mathbb{C}} \setminus \{0\}$ and $0 \cdot \infty = 1 = \infty \cdot 0$. We think of the complexified unit circle as the set $\{(z, w) \in \widehat{\mathbb{C}}^2 : z \cdot w = 1\}$.



on the open unit disk $\mathbb{D} := \{z \in \mathbb{C} : |z| < 1\}$ and the spherical Laplacian

$$\Delta_{\widehat{\mathbb{C}}} := -4(1 + |z|^2)^2 \partial_z \partial_{\bar{z}}$$

on the Riemann sphere $\widehat{\mathbb{C}}$ from a complex analytic point of view and, in addition, it also shows how they are interrelated to one another. Beyond that, the complex point of view taken in this paper offers several other useful advantages. In particular, it connects in a natural way the fine structure of the eigenspaces of the hyperbolic and spherical Laplacians as described by Helgason [21] and Rudin [37] with the *maximal* domain of existence of the corresponding holomorphic eigenfunctions of the invariant Laplacian Δ_{zw} .

As one instance, we analyze from a complex analysis point of view the building blocks of each λ -eigenspace of the hyperbolic Laplacian $\Delta_{\mathbb{D}}$, which have previously been identified, e.g., by Helgason in [21] using real-variable methods. It turns out that these so-called Poisson–Fourier modes naturally extend to holomorphic eigenfunctions of Δ_{zw} which are (maximally) defined either on Ω or on one of three distinguished subdomains of Ω depending on the choice of the eigenvalue λ . This then allows a transparent proof that each holomorphic eigenfunction of Δ_{zw} defined on any rotationally invariant subdomain of Ω has a unique spectral decomposition in form of a locally uniformly and absolutely convergent infinite series composed of Poisson–Fourier modes. In the special case of the bidisk \mathbb{D}^2 and further restriction to the “diagonal” $\{(z, \bar{z}) : z \in \mathbb{D}\}$, we recover the spectral decomposition of the smooth eigenfunctions of $\Delta_{\mathbb{D}}$ on the unit disk \mathbb{D} as described, e.g., in [21] or [8].

As a second instance, we investigate the structure of the closed “Möbius invariant” subspaces of any fixed eigenspace $X_\lambda(\mathbb{D})$ of $\Delta_{\mathbb{D}}$ from a complex analysis point of view. This topic has been investigated in detail by Rudin in [37] using purely “real” methods. We shall see that the distinguished subdomains of Ω mentioned above naturally lead to the same distinction between exceptional and non-exceptional eigenvalues which has been found by Rudin. In Rudin’s work, the exceptional cases correspond to the eigenvalues $\lambda = 4m(m + 1)$, $m = 0, 1, 2, \dots$, and they are characterized by the existence of three nontrivial Möbius invariant closed subspaces of $X_\lambda(\mathbb{D})$, exactly one of which, $X_\lambda^0(\mathbb{D})$ say, is finite dimensional. It turns out that a complex number $\lambda \in \mathbb{C}$ is an exceptional eigenvalue in the sense of Rudin if and only if the invariant Laplacian Δ_{zw} has a globally defined (that is, defined on Ω) holomorphic λ -eigenfunction. Moreover, in this case, the unique finite-dimensional invariant subspace $X_\lambda^0(\mathbb{D})$ corresponds precisely to the full λ -eigenspace of all globally defined λ -eigenfunctions of the invariant Laplacian Δ_{zw} , which then, in fact, is invariant under the *full* group of all Möbius transformations. For the other two nontrivial invariant subspaces of the exceptional $\Delta_{\mathbb{D}}$ -eigenspace $X_\lambda(\mathbb{D})$ discovered by Rudin as well as the full eigenspace $X_\lambda(\mathbb{D})$ itself, we give a similar but more intricate description in form of Runge-type approximation results in terms of holomorphic λ -eigenfunctions defined precisely on one of the distinguished three subdomains of Ω , see Theorem 2.10.

In the next section, we give an account of the main results of this work and their ramifications for the spectral theory of the hyperbolic and spherical Laplacian as well as an outline of the structure of the remaining sections. The accompanying papers [18, 19, 27, 32] are related to other aspects of the function theory of the set Ω , the

complement of the complexified unit circle, and its applications. Our interest in this set and its inhabitants, the holomorphic functions on Ω , first arose in connection with previous work [7, 10, 12, 28, 41, 43] on canonical Wick-type star products in strict deformation quantization of the unit disk and the Riemann sphere, and from our desire to understand the somehow mysterious role played by Ω and, in particular, by its function-theoretic properties in this regard. A partial explanation was given in [19], where it was indicated that invariant differential operators of Pöschl–Minda type, on the one hand, effectively facilitate and unify the study of the star products on the disk and the sphere, and on the other hand, are perhaps best understood as operators acting on the spaces of holomorphic functions on Ω and its three distinguished subdomains. We started wondering whether and how the most basic differential operator acting on Ω , the invariant Laplacian $\Delta_{z,w}$, and its spectral theory possibly fit into this emerging picture. This paper describes what we have found. In the forthcoming paper [32] of the second-named author, these endeavours will come to full circle: it is shown that there are globally defined eigenfunctions of the Laplacian $\Delta_{z,w}$ which form a Schauder basis of the Fréchet space $\mathcal{H}(\Omega)$ of all holomorphic functions on Ω . The results of the present paper then imply that the algebra

$$\mathcal{A}(\mathbb{D}) = \{f : \mathbb{D} \rightarrow \mathbb{C} \mid f(z) = F(z, \bar{z}) \text{ for all } z \in \mathbb{D} \text{ for some } F \in \mathcal{H}(\Omega)\},$$

for which the Wick-star product on \mathbb{D} in [28] is constructed, admits a spectral decomposition precisely into the finite-dimensional invariant subspaces of the exceptional eigenspaces of the hyperbolic Laplace operator $\Delta_{\mathbb{D}}$ on \mathbb{D} discovered by Rudin [37] many years ago. This provides an intrinsic characterization of the algebra $\mathcal{A}(\mathbb{D})$ in terms of the natural hyperbolic geometry of the unit disk and its canonical invariant Laplacian $\Delta_{\mathbb{D}}$.

2 Overview and main results

In order to place the results of this paper into a broader context, we begin by recalling in greater detail, the striking distinction between *exceptional* and *non-exceptional* eigenvalues of the hyperbolic Laplace operator $\Delta_{\mathbb{D}}$ and its relevance for the study of the *invariant* subspaces of the $\Delta_{\mathbb{D}}$ -eigenspaces which has been discovered by Rudin [37]. The Fréchet space of all twice continuously (real) differentiable functions $f : \mathbb{D} \rightarrow \mathbb{C}$ equipped with the standard compact-open topology is denoted by $C^2(\mathbb{D})$. For each $\lambda \in \mathbb{C}$, we denote by $X_{\lambda}(\mathbb{D})$, the vector space of all λ -eigenfunctions $f \in C^2(\mathbb{D})$ of the hyperbolic Laplacian, that is,

$$X_{\lambda}(\mathbb{D}) = \{f \in C^2(\mathbb{D}) : \Delta_{\mathbb{D}}f = \lambda f \text{ on } \mathbb{D}\}.$$

It is known that each such $\Delta_{\mathbb{D}}$ -eigenspace $X_{\lambda}(\mathbb{D})$ is a closed infinite-dimensional subspace of $C^2(\mathbb{D})$. The hyperbolic Laplacian $\Delta_{\mathbb{D}}$ is invariant under the full group of all conformal automorphisms (biholomorphic maps) T of the unit disk \mathbb{D} in the sense that

$$\Delta_{\mathbb{D}}(f \circ T) = (\Delta_{\mathbb{D}}f) \circ T \quad \text{for all } f \in C^2(\mathbb{D}).$$

In order to emphasize that this group consists entirely of Möbius transformations, we call it the *Möbius group* of \mathbb{D} and denote it by $\mathcal{M}(\mathbb{D})$. A closed subspace Y of

$C^2(\mathbb{D})$ is called *Möbius invariant* if $f \circ \psi \in Y$ for all $f \in Y$ and all $\psi \in \mathcal{M}(\mathbb{D})$.² It is called nontrivial if $Y \neq \{0\}$ and $Y \neq X$.

Theorem 2.1 (Rudin [37]) *Let $\lambda \in \mathbb{C}$.*

(NE) *If $\lambda \neq 4m(m+1)$ for $m = 0, 1, 2, \dots$, then $X_\lambda(\mathbb{D})$ has no nontrivial Möbius invariant subspaces.*

(E) *If $\lambda = 4m(m+1)$ for some $m = 0, 1, 2, \dots$, then $X_\lambda(\mathbb{D})$ has precisely three distinct nontrivial Möbius invariant subspaces. There is exactly one nontrivial Möbius invariant subspace of $X_\lambda(\mathbb{D})$ which is finite dimensional; its dimension is $2m+1$.*

The alternative (E) in Theorem 2.1 will be called the *exceptional* case and the unique finite-dimensional Möbius invariant subspace of $X_\lambda(\mathbb{D})$ will be denoted by $X_\lambda^0(\mathbb{D})$. The alternative (NE) will be referred to as the *non-exceptional* case.

One of the main results of the present paper is a complete analog of Theorem 2.1 with $\Delta_{\mathbb{D}}$ replaced by the differential operator Δ_{zw} (see Theorem 2.6). Apart from being potentially interesting in its own right, it provides a concrete function-theoretic description of the exceptional eigenspaces in Theorem 2.1. This also adds a conceptual component to Rudin's handling of the invariant eigenspaces of the hyperbolic Laplace operator $\Delta_{\mathbb{D}}$.

The characteristic feature of our approach is to look for *holomorphic* solutions F of the eigenvalue equation

$$(2.1) \quad \Delta_{zw}F = \lambda F$$

defined on a subdomain D of Ω which we wish to choose as large as possible depending on the eigenvalue λ . These *maximal domains of existence* (see Definition 2.5) of the λ -eigenfunctions turn out to be the only essential ingredients which are needed to give a complete description of the (invariant) λ -eigenspaces of the operator Δ_{zw} and its offsprings $\Delta_{\mathbb{D}}$ and $\Delta_{\widehat{\mathbb{C}}}$.

In order to state our main results, we have to adapt the notation which we have introduced above for the hyperbolic Laplacian to the case of the differential operator $\Delta_{zw} = 4(1-zw)^2 \partial_z \partial_w$. Instead of working in the Fréchet space $C^2(\mathbb{D})$, we now fix a subdomain D of the set $\Omega = \{(z, w) \in \widehat{\mathbb{C}} : z \cdot w \neq 1\}$, and work in the Fréchet space $\mathcal{H}(D)$ of all complex-valued holomorphic functions defined on D (again equipped with the topology of locally uniform convergence, this time on D). Our goal is to determine the holomorphic solutions $F : D \rightarrow \mathbb{C}$ of the eigenvalue equation (2.1), i.e., we are interested in the Δ_{zw} -eigenspaces

$$X_\lambda(D) := \{F \in \mathcal{H}(D) : \Delta_{zw}F = \lambda F \text{ on } D\}, \quad \lambda \in \mathbb{C}.$$

With regard to Rudin's theorem (Theorem 2.1), a particularly natural choice for the domain D is the *bidisk* $\mathbb{D}^2 := \mathbb{D} \times \mathbb{D}$ since it is easy to see that for every $F \in X_\lambda(\mathbb{D}^2)$, the "restriction" of F to the "diagonal" $\{(z, \bar{z}) : z \in \mathbb{D}\}$, that is, $f(z) := F(z, \bar{z})$, $z \in \mathbb{D}$, yields an eigenfunction $f \in C^2(\mathbb{D})$ of the hyperbolic Laplacian $\Delta_{\mathbb{D}}$ for the eigenvalue λ . In fact, the following result shows that every $f \in X_\lambda(\mathbb{D})$ arises in this fashion, that

²Möbius invariant spaces are called \mathcal{M} -spaces in [37]. We will reserve the symbol \mathcal{M} for some other purpose.

is, it has an “extension” to an eigenfunction F of $\Delta_{z\bar{w}}$ which is holomorphic on the bidisk \mathbb{D}^2 .

Theorem 2.2 (Smooth eigenfunctions of $\Delta_{\mathbb{D}}$ on \mathbb{D} vs. holomorphic eigenfunctions of $\Delta_{z\bar{w}}$ on \mathbb{D}^2) *Let $\lambda \in \mathbb{C}$ and $f \in C^2(\mathbb{D})$ such that $\Delta_{\mathbb{D}}f = \lambda f$ on \mathbb{D} . Then there is a uniquely determined function $F \in \mathcal{H}(\mathbb{D}^2)$ such that $\Delta_{z\bar{w}}F = \lambda F$ on \mathbb{D}^2 and $f(z) = F(z, \bar{z})$ for all $z \in \mathbb{D}$. Moreover, the induced bijective linear map*

$$X_{\lambda}(\mathbb{D}) \rightarrow X_{\lambda}(\mathbb{D}^2)$$

is continuous.

Remark 2.3 If $f \in C^2(\mathbb{D})$ solves $\Delta_{\mathbb{D}}f = \lambda f$ on \mathbb{D} , then f is real-analytic (see [38, Theorem 4.2.5]), so there is trivially a holomorphic function \tilde{F} defined on some open neighborhood $U \subseteq \mathbb{D}^2$ of the diagonal $((z, \bar{z}) \in \mathbb{D}^2 : z \in \mathbb{D})$ such that $\tilde{F}(z, \bar{z}) = f(z)$ for all $z \in \mathbb{D}$ satisfying $(z, \bar{z}) \in U$. In view of a well-known variant of the identity principle (see Lemma 3.1) there is only one such holomorphic extension $\tilde{F} : U \rightarrow \mathbb{C}$ of $f \in C^2(\mathbb{D})$ to U . The point of Theorem 2.2 is that \tilde{F} has a holomorphic extension (at least) to the bidisk \mathbb{D}^2 .

Theorem 2.2 gives rise to the following definition.

Definition 2.4 Let $\lambda \in \mathbb{D}$. Then the continuous bijective linear mapping

$$\mathcal{R}_h : X_{\lambda}(\mathbb{D}^2) \rightarrow X_{\lambda}(\mathbb{D}), \quad \mathcal{R}_h(F)(z) := F(z, \bar{z}) \quad (z \in \mathbb{D})$$

is called the *hyperbolic restriction map*. Its continuous inverse

$$\mathcal{E}_h := (\mathcal{R}_h)^{-1} : X_{\lambda}(\mathbb{D}) \rightarrow X_{\lambda}(\mathbb{D}^2)$$

is called the *hyperbolic extension map* from $X_{\lambda}(\mathbb{D})$ to $X_{\lambda}(\mathbb{D}^2)$.

The hyperbolic restriction and extension mappings provide the bridge between the holomorphic spectral theory of the invariant Laplacian $\Delta_{z\bar{w}}$ and the smooth spectral theory of the hyperbolic Laplacian $\Delta_{\mathbb{D}}$. In particular, one can study the spectral properties of $\Delta_{\mathbb{D}}$ on \mathbb{D} from the viewpoint of complex analysis on the bidisk \mathbb{D}^2 . Based on Theorem 2.2, we can now proceed to associate the fine structure of the eigenspaces $X_{\lambda}(\mathbb{D})$ with the *maximal domains of existence* of holomorphic eigenfunctions in $X_{\lambda}(\mathbb{D}^2)$, a concept which is defined as follows (cf. [14, p. 97]).

Definition 2.5 Let F be a holomorphic function on the bidisk \mathbb{D}^2 . A subdomain $D \subseteq \Omega$ that contains the bidisk \mathbb{D}^2 is called a *maximal domain of existence* of F if the function F has a holomorphic extension to D but to no strictly larger subdomain of Ω .

For our purposes, this definition is natural in several respects. First, the condition $D \supseteq \mathbb{D}^2$ obviously comes from Theorem 2.2. Second, the condition $D \subseteq \Omega$ is natural in view of the fact that every function which is holomorphic on a subdomain of $\widehat{\mathbb{C}}^2$ that is strictly larger than Ω is necessarily constant (see Theorem 5.3 in [18]). In particular, the largest possible maximal domain of existence of any eigenfunction of $\Delta_{z\bar{w}}$ in $\mathcal{H}(\mathbb{D}^2)$ is Ω . Note, this also means that the maximal domain of existence for the *constant* eigenfunctions is Ω .

With this concept at hand, we can now give a function-theoretic characterization of the exceptional eigenvalues of the hyperbolic Laplacian $\Delta_{\mathbb{D}}$ and the corresponding unique finite-dimensional nontrivial Möbius invariant subspaces $X_{\lambda}^0(\mathbb{D})$ of the $\Delta_{\mathbb{D}}$ -eigenspaces $X_{\lambda}(\mathbb{D})$. In addition, the following theorem provides an equivalent condition in terms of existence of globally defined eigenfunctions of the spherical Laplacian $\Delta_{\widehat{\mathbb{C}}}$.

Theorem 2.6 (Exceptional smooth eigenfunctions $\Delta_{\mathbb{D}}$ on \mathbb{D} vs. holomorphic eigenfunctions of $\Delta_{z,w}$ on Ω vs. smooth eigenfunctions of $\Delta_{\widehat{\mathbb{C}}}$ on $\widehat{\mathbb{C}}$) *Let $\lambda \in \mathbb{C}$. Then the following conditions are pairwise equivalent:*

- (i) *There exists a function $F \in X_{\lambda}(\mathbb{D}^2)$ with Ω as maximal domain of existence.*
- (ii) *There exists a function $g \in C^2(\widehat{\mathbb{C}})$ such that $\Delta_{\widehat{\mathbb{C}}}g = \lambda g$ on $\widehat{\mathbb{C}}$.*
- (iii) *λ is an exceptional eigenvalue of $\Delta_{\mathbb{D}}$, i.e., $\lambda = 4m(m+1)$ for some nonnegative integer m .*

If one of these conditions is in place and $\lambda = 4m(m+1)$, then $\dim X_{\lambda}(\Omega) = 2m+1$ and the following statements hold:

- (a) *(Function-theoretic description of $X_{\lambda}^0(\mathbb{D})$)*

The domain Ω is the maximal domain of existence for every $F \in X_{\lambda}^0(\mathbb{D}^2)$, i.e.,

$$X_{\lambda}^0(\mathbb{D}) = \mathcal{R}_h(X_{\lambda}(\Omega)).$$

- (b) *(Function-theoretic description of the smooth eigenfunctions of $\Delta_{\widehat{\mathbb{C}}}$ on $\widehat{\mathbb{C}}$)*

$$(2.2) \quad \{g \in C^2(\widehat{\mathbb{C}}) : \Delta_{\widehat{\mathbb{C}}}g = \lambda g \text{ on } \widehat{\mathbb{C}}\} = \{\widehat{\mathbb{C}} \ni z \mapsto F(z, -\bar{z}) : F \in X_{\lambda}(\Omega)\}.$$

Remark 2.7 (a) Note that $X_0^0(\mathbb{D}^2)$ and $X_0(\Omega)$ consist precisely of the constant functions.

- (b) In view of Theorem 2.6, an eigenvalue λ of $\Delta_{\mathbb{D}}$ is exceptional if and only if there exists a holomorphic eigenfunction of $\Delta_{z,w}$ on \mathbb{D}^2 with the largest possible maximal domain of existence, the set Ω .
- (c) In Theorem 2.6(a), we think of $X_{\lambda}(\Omega)$ as a subspace of $X_{\lambda}(\mathbb{D}^2)$. It is a closed subspace of $X_{\lambda}(\mathbb{D}^2)$ (in the topology of $\mathcal{H}(\mathbb{D}^2)$) because it is finite dimensional.
- (d) (Spherical restriction map)

If $\lambda \in \mathbb{C}$ is an exceptional eigenvalue, then the *spherical restriction map*

$$\mathcal{R}_s : X_{\lambda}(\Omega) \rightarrow C^2(\widehat{\mathbb{C}}), \quad \mathcal{R}_s(F)(z) := F(z, -\bar{z}) \quad (z \in \widehat{\mathbb{C}})$$

is well-defined. By Theorem 2.6(b), \mathcal{R}_s is a bijection from $X_{\lambda}(\Omega)$ onto the λ -eigenspace $\{g \in C^2(\widehat{\mathbb{C}}) : \Delta_{\widehat{\mathbb{C}}}g = \lambda g \text{ on } \widehat{\mathbb{C}}\}$ of $\Delta_{\widehat{\mathbb{C}}}$, which is clearly continuous.

With Theorem 2.6, we have reached two of our goals, a conceptual characterization of exceptional eigenvalues and the finite-dimensional nontrivial Möbius invariant subspaces of the eigenspaces $X_{\lambda}(\mathbb{D})$ of the hyperbolic Laplacian $\Delta_{\mathbb{D}}$. Next, we address the infinite-dimensional nontrivial Möbius invariant subspaces of the $\Delta_{\mathbb{D}}$ -eigenspaces $X_{\lambda}(\mathbb{D})$. This turns out to be more difficult, and first requires clarification of the invariance properties of the Laplacian $\Delta_{z,w}$. Implicitly, the underlying difficulty is already present in Theorem 2.6, and can be seen as follows. Let $\lambda \in \mathbb{C}$ be an exceptional eigenvalue of $\Delta_{\mathbb{D}}$. Then the finite-dimensional Möbius invariant subspace $X_{\lambda}^0(\mathbb{D})$ is invariant under all automorphisms of \mathbb{D} . However, the corresponding eigenspace

$X_\lambda(\Omega)$ for $\Delta_{z,w}$, which consists of functions holomorphic on Ω is, loosely speaking, invariant under a much larger group of automorphisms. In fact, we first note that the invariant Laplacian $\Delta_{z,w}$ of Ω is *not* invariant under all biholomorphic automorphisms of Ω in the sense that the invariance condition

$$(2.3) \quad \Delta_{z,w}(F \circ T) = (\Delta_{z,w}F) \circ T \quad \text{for all } F \in \mathcal{H}(\Omega)$$

holds for all $T \in \text{Aut}(\Omega)$. The reason is simply that the automorphism group $\text{Aut}(\Omega)$ is much too large (see [18]). However, $\Delta_{z,w}$ is invariant under the subgroup \mathcal{M} of $\text{Aut}(\Omega)$ defined by

$$(2.4) \quad \mathcal{M} := \bigcup_{\psi \in \mathcal{M}(\widehat{\mathbb{C}})} \left\{ (z, w) \mapsto \left(\psi(z), \frac{1}{\psi(1/w)} \right), (z, w) \mapsto \left(\psi(w), \frac{1}{\psi(1/z)} \right) \right\},$$

where we write $\mathcal{M}(\widehat{\mathbb{C}})$ for the group of all Möbius transformations $\psi : \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$. This is easy to prove by direct verification. (Conversely, one can show that every $T \in \text{Aut}(\Omega)$ for which the invariance property (2.3) holds does necessarily belong to the subgroup \mathcal{M} , see [18, Theorem 5.2], but we do not need such a result in this paper.) Since \mathcal{M} is “induced” by the set $\mathcal{M}(\widehat{\mathbb{C}})$ of all Möbius transformations, we call \mathcal{M} the *Möbius group of Ω* . Note that $\mathcal{M}(\widehat{\mathbb{C}})$ is strictly bigger than $\mathcal{M}(\mathbb{D})$,³ so the Möbius group \mathcal{M} of Ω is strictly larger than the Möbius group $\mathcal{M}(\mathbb{D})$ of \mathbb{D} . Now, while $X_\lambda^0(\mathbb{D})$ is invariant under each element of $\mathcal{M}(\mathbb{D})$, the set $X_\lambda(\Omega)$ is invariant under each element of \mathcal{M} , simply because $X_\lambda(\Omega)$ is the entire λ -eigenspace of the Laplacian $\Delta_{z,w}$ on Ω and $\Delta_{z,w}$ is invariant with respect to the Möbius group \mathcal{M} .

In view of this discussion, it is now clear that a suitable concept of invariance for the eigenspaces $X_\lambda(\mathbb{D}^2)$ of $\Delta_{z,w}$ on the bidisk \mathbb{D}^2 has to be based on the group

$$\mathcal{M}(\mathbb{D}^2) := \text{Aut}(\mathbb{D}^2) \cap \mathcal{M},$$

which we call the *Möbius group of the bidisk \mathbb{D}^2* . It consists precisely of all automorphisms of the bidisk \mathbb{D}^2 which have the invariance property (2.3). In fact, it is not difficult to show that

$$(2.5) \quad \mathcal{M}(\mathbb{D}^2) = \bigcup_{\psi \in \mathcal{M}(\mathbb{D})} \left\{ (z, w) \mapsto \left(\psi(z), \frac{1}{\psi(1/w)} \right), (z, w) \mapsto \left(\psi(w), \frac{1}{\psi(1/z)} \right) \right\}.$$

Clearly, each $\Delta_{z,w}$ -eigenspace $X_\lambda(\mathbb{D}^2)$ is invariant with respect to $\mathcal{M}(\mathbb{D}^2)$, that is, whenever $F \in X_\lambda(\mathbb{D}^2)$ and $T \in \mathcal{M}(\mathbb{D}^2)$, then $F \circ T \in X_\lambda(\mathbb{D}^2)$. Each closed subspace of $\mathcal{H}(\mathbb{D}^2)$ which is invariant under the Möbius group $\mathcal{M}(\mathbb{D}^2)$ will be called a *Möbius invariant subspace*. Theorem 2.2 implies immediately the following result.

Corollary 2.8 *Let $\lambda \in \mathbb{C}$ and Y be a subspace of $X_\lambda(\mathbb{D})$. Then Y is Möbius invariant if and only if $\mathcal{E}_h(Y)$ is a Möbius invariant subspace of $X_\lambda(\mathbb{D}^2)$.*

We can now give a function-theoretic characterization of the Möbius invariant subspaces of $X_\lambda(\mathbb{D}^2)$ and thereby, in view of Corollary 2.8, the Möbius invariant

³ $\mathcal{M}(\widehat{\mathbb{C}})$ is real six dimensional, while $\mathcal{M}(\mathbb{D})$ is real three dimensional.

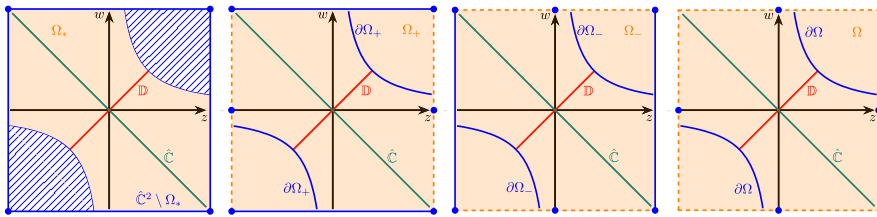


Figure 1: Schematic picture of the sets Ω_* , Ω_+ , Ω_- , and Ω (from left to right) with points at infinity. Here, \mathbb{D} is identified with the diagonal $\{(z, \bar{z}) : z \in \mathbb{D}\}$ and $\widehat{\mathbb{C}}$ with the rotated diagonal $\{(z, -\bar{z}) : z \in \widehat{\mathbb{C}}\}$.

subspaces of $X_\lambda(\mathbb{D})$. The following subdomains of Ω play the essential role for this purpose:

$$\begin{aligned} \Omega_+ &:= \Omega \setminus \{(z, \infty) : z \in \widehat{\mathbb{C}}\}, \\ \Omega_- &:= \Omega \setminus \{(\infty, w) : w \in \widehat{\mathbb{C}}\}, \\ \Omega_* &:= \{(z, w) \in \mathbb{C}^2 : zw \in \mathbb{C} \setminus [1, \infty)\}. \end{aligned}$$

Note that each of these three subdomains of Ω contains the bidisk \mathbb{D}^2 . Moreover,

$$\Omega = \Omega_+ \cup \Omega_- \cup \{(\infty, \infty)\}$$

and

$$\mathcal{H}(\Omega) = \mathcal{H}(\Omega_+) \cap \mathcal{H}(\Omega_-).$$

Figure 1 provides a schematic view of Ω and its distinguished subsets.

We note in passing that the subdomains Ω_+ and Ω_- arise naturally in the study of the Fréchet space structure of $\mathcal{H}(\Omega)$ (see [18]) and also for studying invariant differential operators of Peschl–Minda type acting on $\mathcal{H}(\Omega)$ (see [19]). The following result shows that they are also useful for describing the Möbius invariant subspaces of the eigenspaces of the invariant Laplacian Δ_{zw} . We use the following terminology.

Definition 2.9 Let $U \subseteq V$ be subdomains of some complex manifold, and let $Y \subseteq X$ be subsets of $\mathcal{H}(U)$. We say that “ $X \cap \mathcal{H}(V)$ is dense in Y ” if $X \cap \mathcal{H}(V) \subseteq Y$ and if every function in Y can be approximated locally uniformly on U by functions in X which have a holomorphic extension to V .

We are now, finally, in a position to formulate the main result of this paper.

Theorem 2.10 Let $\lambda \in \mathbb{C}$, and let Y be a nontrivial Möbius invariant subspace of $X_\lambda(\mathbb{D}^2)$. Then one and only one of the following four alternatives holds.

- (E₀) $Y = X_\lambda(\mathbb{D}^2) \cap \mathcal{H}(\Omega)$ and $\dim Y < \infty$.
- (E₊) $X_\lambda(\mathbb{D}^2) \cap \mathcal{H}(\Omega_+)$ is dense in Y .
- (E₋) $X_\lambda(\mathbb{D}^2) \cap \mathcal{H}(\Omega_-)$ is dense in Y .
- (NE) $X_\lambda(\mathbb{D}^2) \cap \mathcal{H}(\Omega_*)$ is dense in Y . In this case $Y = X_\lambda(\mathbb{D}^2)$.

In addition,

- (i) (E_0) , (E_+) , and (E_-) cannot occur if $\lambda \neq 4m(m+1)$ for every nonnegative integer m .
- (ii) (E_0) holds if and only if $\dim Y < \infty$. In this case, $\dim Y = 2m+1$ and $\lambda = 4m(m+1)$ for some nonnegative integer m .
- (iii) None of the density statements (NE) , (E_+) , and (E_-) may be improved to equalities.

Remark 2.11 Recall that two domains $U \subseteq V$ in \mathbb{C}^n form a Runge pair (U, V) if every function in $\mathcal{H}(U)$ can be approximated locally uniformly in U by functions in $\mathcal{H}(V)$. Identifying domains $U \subseteq V$ as Runge pairs is a fundamental and, in many cases, challenging problem in complex analysis. Note that in our terminology, (U, V) is a Runge pair if and only if $\mathcal{H}(U) \cap \mathcal{H}(V)$ is dense in $\mathcal{H}(U)$. For subspaces $Y \subseteq X$ of $\mathcal{H}(U)$, it is tempting to call $(Y, X, \mathcal{H}(V))$ a *Runge triple*, if $X \cap \mathcal{H}(V)$ is dense in Y . Then for every nontrivial infinite-dimensional Möbius invariant subspace Y of $X_\lambda(\mathbb{D}^2)$ exactly one of the triples

$$(Y, X_\lambda(\mathbb{D}^2), \mathcal{H}(\Omega_+)), \quad (Y, X_\lambda(\mathbb{D}^2), \mathcal{H}(\Omega_-)), \quad (Y, X_\lambda(\mathbb{D}^2), \mathcal{H}(\Omega_*))$$

is a Runge triple. Theorem 2.10 can therefore be regarded as a Runge-type approximation theorem for the Möbius invariant spaces of eigenfunctions of the invariant Laplacian Δ_{z_w} .

The plan of the paper is as follows. We introduce some basic concepts and notation in a preliminary Section 3. In Sections 4 and 5, we develop the general spectral theory of the invariant Laplacian Δ_{z_w} on Ω in analogy to the well-established spectral theory of the hyperbolic Laplacian $\Delta_{\mathbb{D}}$ on \mathbb{D} . In a sense, we rather closely follow the presentation Berenstein and Gay [8, Section 1.6] have given for the spectral theory of the hyperbolic Laplacian, but we have made an effort either to provide even more rigorous proofs or to give precise references to the literature for all auxiliary results which are needed. In contrast to [8], we completely work in the holomorphic setting. On the one hand, this makes it possible to take advantage of many efficient tools from complex analysis which are not available otherwise. On the other hand, we need to incorporate from the beginning the maximal domain of existence of eigenfunctions; an issue which does not even show up when working “only” on the unit disk. Here, our approach requires some finer analysis of the building blocks of the eigenfunctions, namely, certain hypergeometric functions and their integral representations in terms of Poisson–Fourier modes.

In Section 6, we prove Theorem 2.2 and show that it is in some sense best possible by providing an explicit example. This implies that the smooth spectral theory of the hyperbolic Laplacian $\Delta_{\mathbb{D}}$ on the disk \mathbb{D} and the holomorphic spectral theory of the Laplacian Δ_{z_w} on the bidisk \mathbb{D}^2 are essentially equivalent. In the same spirit, we relate the smooth spectral theory of the spherical Laplacian $\Delta_{\widehat{\mathbb{C}}}$ with the holomorphic spectral theory of the Laplacian Δ_{z_w} on Ω as well as the exceptional eigenvalues of the hyperbolic Laplacian $\Delta_{\mathbb{D}}$ (see Theorem 6.2). Section 7 is devoted to a study of the transformation behavior of the Poisson–Fourier modes under precompositions with elements of the Möbius group \mathcal{M} . These results are needed for Section 8, where we prove our main results, Theorems 2.6 and 2.10. By and large, we follow Rudin’s [37] treatment of invariant subspaces of eigenfunctions of the hyperbolic Laplacian,

but again completely working in the holomorphic setting; we briefly comment on the similarities and differences between our and Rudin's approach in Remark 8.5. We close the paper with Section 9, which connects the Poisson–Fourier modes to the invariant differential operators of Peschl–Minda type studied in [19].

Four final preliminary remarks are in order. First, treating the hyperbolic eigenvalue equation $\Delta_{\mathbb{D}}f = \lambda f$ and the spherical eigenvalue equation $\Delta_{\widehat{\mathbb{C}}}f = \lambda f$ as special cases of the more general complex eigenvalue equation $\Delta_{zw}F = \lambda F$ has been a recurrent theme in the literature for a long time. To mention but a few of the many references, we refer for instance to the papers [4–6] and their bibliographies. What seems to be new is the systematic study of the maximal domains of existence of the holomorphic solutions of $\Delta_{zw}F = \lambda F$ and its ramifications for the study of the invariant subspaces of the $\Delta_{\mathbb{D}}$ -eigenspaces. As a second remark, we should point out that Rudin's work [37] is in fact concerned with the invariant Laplace operator on the unit ball of \mathbb{C}^n , while our focus is exclusively on the complex one-dimensional case $n = 1$. Third, even though we are superficially dealing with holomorphic functions of two complex variables, we only need very few and only elementary facts from the theory of several complex variables. Finally, Helgason [20–22] has systematically studied invariant differential operators and their eigenvalue problem in the setting of homogeneous spaces. In contrast to our holomorphic approach, he used entirely real methods. This Lie theoretic approach has since been generalized significantly. While providing a comprehensive list of references would go beyond what we can achieve here, we would like to mention [23, 26], who generalized the theory to higher-dimensional symmetric spaces, and Maass [29], who initiated the vast and fruitful research of Maass wave forms.

3 Notation and preliminaries

We denote the open unit disk in \mathbb{C} by \mathbb{D} , the bidisk $\mathbb{D} \times \mathbb{D}$ by \mathbb{D}^2 and the Riemann sphere by $\widehat{\mathbb{C}}$. The open disk of radius $r > 0$ centered at the origin is denoted by \mathbb{D}_r . Moreover, we write $\mathbb{N} := \{1, 2, \dots\}$ for the set of positive integers, $\mathbb{N}_0 := \mathbb{N} \cup \{0\}$ and \mathbb{Z} for the set of all integers. For an open subset U of $\widehat{\mathbb{C}}$ or $\widehat{\mathbb{C}}^2$, we write ∂U for its boundary and \bar{U} for its closure. The set of all twice resp. infinitely (real) differentiable functions $f : U \rightarrow \mathbb{C}$ is denoted $C^2(U)$ resp. $C^\infty(U)$, and we write $\mathcal{H}(U)$ for set of all holomorphic functions $f : U \rightarrow \mathbb{C}$.

The set $\Omega = \widehat{\mathbb{C}}^2 \setminus \{(z, w) \in \widehat{\mathbb{C}}^2 : z \cdot w \neq 1\}$ is a complex manifold of complex dimension 2 and an open submanifold of $\widehat{\mathbb{C}}^2$. In order to describe the complex structure of Ω only two charts are necessary, the *standard chart*

$$\Omega \cap (\mathbb{C} \times \mathbb{C}) \rightarrow \mathbb{C}^2, \quad (u, v) \mapsto (u, v)$$

and the *flip chart*

$$\Omega \cap (\widehat{\mathbb{C}} \setminus \{0\} \times \widehat{\mathbb{C}} \setminus \{0\}) \rightarrow \mathbb{C}^2, \quad (u, v) \mapsto (1/v, 1/u).$$

In these local coordinates, the *invariant Laplace operator* Δ_{zw} of the complex manifold Ω is then given by

$$\Delta_{zw}f(z, w) = 4(1 - zw)^2 \partial_z \partial_w f(z, w),$$

and it is easily seen that Δ_{zw} is a well-defined object. Here, ∂_z and ∂_w denote the Wirtinger derivatives with respect to z and w , respectively.

One of the few elementary results from the theory of functions of several complex variables we need in this paper is the following simple lemma.

Lemma 3.1 (“Two variable identity principle” (p. 18 in [35])) *Let U be a subdomain of \mathbb{C}^2 which contains a point of the form (z, \bar{z}) resp. $(z, -\bar{z})$. Let $f : U \rightarrow \mathbb{C}$ be a holomorphic function such that*

$$\{f(z, \bar{z}) : (z, \bar{z}) \in U\} = \{0\} \quad \text{resp.} \quad \{f(z, -\bar{z}) : (z, -\bar{z}) \in U\} = \{0\}.$$

Then $f \equiv 0$.

Finally, we briefly recall some standard terminology from linear algebra. We denote by $\text{span } M$ the collection of all *finite* linear combinations of elements of a subset M of a given vector space X . The vector spaces that occur in this paper are spaces of smooth or holomorphic functions defined on some open subset U of Ω , which we equip with the standard topology of uniform convergence on compact subsets of U . If M denotes a set of smooth or holomorphic functions on U , we denote by $\text{clos}_U M$ the closure of M with respect to locally uniform convergence on U .

4 Homogeneous eigenfunctions and Poisson–Fourier modes

By making a separation of variables approach, Rudin [37] showed that every $\lambda \in \mathbb{C}$ is an eigenvalue of $\Delta_{\mathbb{D}}$. The analogous result is true for Δ_{zw} : let $n \in \mathbb{Z}$ and suppose $f_n : D \rightarrow \mathbb{C}$ is a holomorphic function defined on a domain $D \subseteq \mathbb{C}^2$ containing the origin $(0, 0)$. Further, assume that f_n is *n-homogeneous*, i.e.,

$$(4.1) \quad f_n(\eta z, w/\eta) = \eta^n f_n(z, w) \quad \text{for all } \eta \in \partial\mathbb{D}$$

and for all $(z, w) \in \mathbb{C}^2$ belonging to the bidisk $\mathbb{D}_r^2 = \mathbb{D}_r \times \mathbb{D}_r$ for some (and hence all) $r > 0$ such that $\mathbb{D}_r^2 \subseteq D$. A consideration of the power series expansion of f_n at $(0, 0)$ in \mathbb{D}_r^2 implies in a straightforward way that there is a holomorphic function $y_n : \mathbb{D}_{r^2} \rightarrow \mathbb{C}$ such that

$$(4.2) \quad f_n(z, w) = y_n(zw)z^n \quad \text{if } n \geq 0 \quad \text{and} \quad f_n(z, w) = y_n(zw)w^{|n|} \quad \text{if } n \leq 0$$

for all $(z, w) \in \mathbb{D}_r^2$. It is, then, easy to see that $\Delta_{zw}f_n = \lambda f_n$ on U if and only if

$$(4.3) \quad 4(1-t)^2 [ty_n''(t) + (|n|+1)y_n'(t)] = \lambda y_n(t)$$

for all $t \in \mathbb{D}_{r^2}$. A power series ansatz shows that there is at most one solution y_n of (4.3) which is holomorphic in a neighborhood of $t = 0$ and normalized such that $y_n(0) = 1$. In order to find this solution, we convert (4.3) into a hypergeometric differential equation as follows. We choose $\mu \in \mathbb{C}$ such that $\lambda = 4\mu(\mu - 1)$, and let $\hat{y}_n(t) := (1-t)^{-\mu} y_n(t)$. Then (4.3) is equivalent to

$$(4.4) \quad t(1-t)\hat{y}_n'' + [c - (a+b+1)t]\hat{y}_n' - ab\hat{y}_n = 0$$

with

$$a = \mu, \quad b = \mu + |n|, \quad c = |n| + 1.$$

It is well-known (see [33, Section 15.10]) that the only solution \hat{y}_n of (4.4) which is holomorphic at $t = 0$ and normalized at $t = 0$ by $\hat{y}_n(0) = 1$ is the hypergeometric series

$${}_2F_1(a, b; c; t) := \sum_{k=0}^{\infty} \frac{(a)_k (b)_k}{(c)_k} \frac{t^k}{k!}, \quad |t| < 1,$$

where

$$(4.5) \quad (\alpha)_k := \prod_{j=0}^{k-1} (\alpha + j), \quad \alpha \in \mathbb{C}, k \in \mathbb{N}_0$$

denotes the (rising) Pochhammer symbol. This procedure leads to the conclusion that

$$(4.6) \quad y_n(t) = (1 - t)^\mu {}_2F_1(\mu, \mu + |n|; |n| + 1; t)$$

is the unique solution of (4.3) which is holomorphic in $t = 0$ and normalized by $y_n(0) = 1$. Note that there are, in fact, two complex numbers $\mu \in \mathbb{C}$ such that $\lambda = 4\mu(\mu - 1)$. However, as we have seen, both necessarily lead to the same holomorphic solution y_n of (4.3) with $y_n(0) = 1$.

Remark 4.1 Note that if $\mu \in \mathbb{C}$ is one solution to $\lambda = 4\mu(\mu - 1)$, then $1 - \mu$ is the other one. As we have seen, both numbers induce the same function (4.6). This also follows from the transformation formula (see [1, Equation (15.3.3)])

$${}_2F_1(a, b; c; t) = (1 - t)^{c-a-b} {}_2F_1(c - a, c - b; c; t).$$

In the following, it will often be convenient to choose μ such that $\text{Re } \mu \geq 1/2$.

By a standard fact about hypergeometric functions (see [1, 15.3.1]), the function y_n has a holomorphic extension at least to the slit plane

$$\mathbb{C} \setminus [1, \infty).$$

Returning to (4.2) with this choice of y_n , we therefore see that

$$(4.7) \quad F_n^\mu(z, w) := (1 - zw)^\mu {}_2F_1(\mu, \mu + |n|; |n| + 1; zw) \cdot \begin{cases} z^n, & \text{if } n \geq 0, \\ |w|^{|n|}, & \text{if } n \leq 0, \end{cases}$$

is holomorphic at least on the domain

$$\Omega_* := \{(z, w) \in \mathbb{C}^2 : zw \in \mathbb{C} \setminus [1, \infty)\}$$

and provides the unique n -homogeneous solution of $\Delta_{zw} f = \lambda f$ on Ω_* up to a multiplicative constant.

Next, we relate the n -homogeneous eigenfunction F_n^μ of Δ_{zw} to a complexified version of the Poisson kernel of the unit disk. This slight change of perspective will turn out to be important in the sequel. In fact, a well-known integral representation formula for the hypergeometric function ${}_2F_1(\mu, \mu + |n|; |n| + 1; zw)$ (see [11, Section 2.5.1, Formula (10), p. 81]) shows that for every $n \in \mathbb{Z}$,

$$(4.8) \quad \binom{\mu + |n| - 1}{|n|} F_n^\mu(z, w) = \frac{(1 - zw)^\mu}{2\pi} \int_0^{2\pi} \frac{e^{int}}{(1 - ze^{-it})^\mu (1 - we^{it})^\mu} dt,$$

with the *generalized binomial coefficients* $\binom{\mu}{n} := \frac{(\mu-n+1)_n}{n!}$ for $\mu \in \mathbb{C}$, $n \in \mathbb{N}_0$. On the right-hand side of this identity, one can recognize the $(-n)$ th Fourier coefficient of (a suitably defined power of order μ of) the function

$$(4.9) \quad P: \mathbb{D}^2 \times \partial\mathbb{D} \rightarrow \mathbb{C}, \quad P(z, w; \xi) := \frac{1 - zw}{(1 - z/\xi)(1 - w\xi)}.$$

Note that if $w = \bar{z} \in \mathbb{D}$, then $P(z, \bar{z}; \xi)$ is the standard Poisson kernel of the unit disk.

Definition 4.2 (Generalized Poisson kernel, Poisson–Fourier mode, PFM)

- (a) The function $P: \mathbb{D}^2 \times \partial\mathbb{D} \rightarrow \mathbb{C}$ defined by (4.9) is called *generalized Poisson kernel*.
 (b) Let $\mu \in \mathbb{C}$ and $n \in \mathbb{Z}$. Then the function $P_n^\mu: \mathbb{D}^2 \rightarrow \mathbb{C}$,

$$(4.10) \quad P_n^\mu(z, w) := \frac{(1 - zw)^\mu}{2\pi} \int_0^{2\pi} \frac{e^{-int}}{(1 - ze^{-it})^\mu (1 - we^{it})^\mu} dt$$

is called the *n*th *Poisson–Fourier mode (PFM) of order μ* .

Here, a^μ is defined for $a \in \mathbb{C} \setminus (-\infty, 0]$ as $\exp(\mu \log a)$, where \log denotes the principal branch of the logarithm.

Remark 4.3 If $w = \bar{z} \in \mathbb{D}$, then $P_n^\mu(z, \bar{z})$ is the *n*th Fourier coefficient of the μ -power of the (real-valued and, in fact, nonnegative) Poisson kernel of the unit disk \mathbb{D} . Further, if $\mu = m \in \mathbb{Z}$, then $P_n^m(z, w)$ also is the *n*th Fourier coefficient of the *m*-power of the generalized Poisson kernel from (4.9).

We can now reformulate (4.8) in terms of PFMs as follows:

$$(4.11) \quad P_n^\mu = (-1)^n \binom{-\mu}{|n|} F_{-n}^\mu = \binom{\mu + |n| - 1}{|n|} F_{-n}^\mu.$$

In particular, P_n^μ has a holomorphic extension from \mathbb{D}^2 to Ω_* , which we continue to denote by P_n^μ .

In analogy with (4.1), we call a domain $D \subseteq \Omega$ *rotationally invariant* if $(\eta z, w/\eta) \in D$ for all $(z, w) \in D$ and all $\eta \in \partial\mathbb{D}$. Summarizing our considerations leads to the following complete description of the *n*-homogeneous eigenfunctions of Δ_{zw} .

Theorem 4.4 Let D be a rotationally invariant subdomain of Ω containing the origin and $n \in \mathbb{Z}$. Suppose that $f \in \mathcal{H}(D)$ is an *n*-homogeneous solution to $\Delta_{zw} f = \lambda f$ for some $\lambda \in \mathbb{C}$ of the form $\lambda = 4\mu(\mu - 1)$ with $\operatorname{Re} \mu \geq 1/2$. Then there is a constant $c \in \mathbb{C}$ such that

$$f(z, w) = c P_{-n}^\mu(z, w), \quad (z, w) \in D \cap \Omega_*.$$

In particular, f has a holomorphic extension to $D \cup \Omega_*$. Moreover, the following dichotomy holds:

(NE) (Non-exceptional case)

If $\lambda \neq 4m(m + 1)$ for all $m \in \mathbb{N}_0$, then Ω_* is the maximal domain of existence of f .

(E) (Exceptional case)

If $\lambda = 4m(m + 1)$ for some $m \in \mathbb{N}_0$, then the maximal domain of existence of f is

$$\begin{aligned} \Omega_+ & \quad n < -m; \\ \Omega_- & \quad \text{if } n > m; \\ \Omega & \quad |n| \leq m. \end{aligned}$$

Proof It remains to prove that for any nonconstant n -homogeneous eigenfunction F_n^μ of Δ_{zw} in $\mathcal{H}(\Omega_*)$ the dichotomy “(NE) vs. (E)” holds.

(i) Let $\mu \notin \mathbb{N}$ and assume that Ω_* is not the maximal domain of existence of F_n^μ which is contained in Ω (see Definition 2.5). Then F_n^μ has a holomorphic extension to some point $(z_0, w_0) \in \Omega$ such that $z_0 w_0 \in \mathbb{R} \cup \{\infty\}$ and $z_0 w_0 > 1$. We first consider the case $n \geq 0$. By definition of F_n^μ , see (4.7), and in view of [1, Eq. 15.3.4], we have

$$\begin{aligned} F_n^\mu(z, w) &= (1 - zw)^\mu {}_2F_1(\mu, \mu + n; |n| + 1; zw) z^n \\ (4.12) \quad &= {}_2F_1\left(\mu, 1 - \mu; |n| + 1; \frac{zw}{zw - 1}\right) z^n, \end{aligned}$$

from which we infer that

$$t \mapsto {}_2F_1(\mu, 1 - \mu; |n| + 1; t)$$

is holomorphic in a neighborhood of the point $x_0 := z_0 w_0 / (z_0 w_0 - 1) \geq 1$. However, see [33, 15.2.3],

$$\begin{aligned} & \lim_{y \rightarrow 0^+} [{}_2F_1(\mu, 1 - \mu; |n| + 1; x + iy) - {}_2F_1(\mu, 1 - \mu; |n| + 1; x - iy)] \\ (4.13) \quad &= \frac{2\pi i}{\Gamma(\mu)\Gamma(1 - \mu)} (x - 1)^{|n|} {}_2F_1(|n| + 1 - \mu, |n| + \mu; |n| + 1; 1 - x) \end{aligned}$$

for all $x > 1$. By our assumption, the left-hand side of (4.13) has to vanish for all $x \in \mathbb{R}$ in some open interval $(x_0, x_0 + \varepsilon)$ with $\varepsilon > 0$ and hence the same is true for the right-hand side. Since the right-hand side is a holomorphic function of x on the domain $\mathbb{C} \setminus (-\infty, 0]$ this is clearly only possible if $\mu \in \mathbb{N}$. The remaining case $n \leq 0$ follows from the case $n \geq 0$ and $F_{-n}^\mu(z, w) = F_n^\mu(w, z)$.

(ii) Let $\mu \in \mathbb{N}$ and $|n| \leq \mu - 1$. If $n \geq 0$, then

$$F_n^\mu(z, w) = \frac{G_n^\mu(zw)}{(1 - zw)^{\mu-1}} z^n,$$

where

$$\begin{aligned} G_n^\mu(zw) &:= (1 - zw)^{2\mu-1} {}_2F_1(\mu, \mu + |n|; |n| + 1; zw) \\ (4.14) \quad &= {}_2F_1(1 - \mu, |n| + 1 - \mu; |n| + 1; zw) \end{aligned}$$

is a polynomial in zw of degree $\mu - n - 1 \geq 0$ (see [1, 15.3.3 and 15.1.1]). Hence, F_n^μ is the product of z^n and a rational function in zw of numerator degree $\mu - n - 1$ and of denominator degree $\mu - 1$ with pole only at the point 1. Therefore, $F_n^\mu \in \mathcal{H}(\Omega)$. Since $F_n^\mu(z, w) = F_{-n}^\mu(w, z)$, this implies $F_n^\mu \in \mathcal{H}(\Omega)$ also for $1 - \mu \leq n \leq 0$.

(iii) Let $\mu \in \mathbb{N}$ and $|n| \geq \mu$. Since $F_n^\mu(z, w) = F_{-n}^\mu(w, z)$ and $\mathcal{H}(\Omega_+) \cap \mathcal{H}(\Omega_-) = \mathcal{H}(\Omega)$ we may assume $n \geq \mu$. In this case, the function G_n^μ in (4.14) is a polynomial

in zw of degree $\mu - 1$, and thus F_n^μ is the product of z^n times a rational function in zw of numerator degree $\mu - 1$ and of denominator degree $\mu - 1$, and is therefore holomorphic on Ω_- . Moreover, F_n^μ has no holomorphic extension to a point (∞, w_0) with $w_0 \in \mathbb{C} \setminus \{0\}$ since in view of (4.14) and (4.12)

$$\begin{aligned} \lim_{z \rightarrow \infty} \frac{G_n^\mu(zw_0)}{(1 - zw_0)^{\mu-1}} &= \lim_{z \rightarrow \infty} {}_2F_1\left(\mu, 1 - \mu; |n| + 1; \frac{zw_0}{zw_0 - 1}\right) = {}_2F_1(\mu, 1 - \mu; |n| + 1; 1) \\ &= \frac{\Gamma(n + 1)\Gamma(n)}{\Gamma(n + 1 - \mu)\Gamma(n + \mu)} \in (0, \infty), \end{aligned}$$

by [1, Eq. 15.1.20], so $|F_n^\mu(z, w_0)| \rightarrow \infty$ as $z \rightarrow \infty$. This implies that Ω_- is the maximal domain of existence of F_n^μ . ■

Corollary 4.5 For every $\lambda \in \mathbb{C}$ and each $n \in \mathbb{Z}$, there is an n -homogeneous holomorphic solution of $\Delta_{zw}f = \lambda f$ on the domain Ω_* . More precisely, if $\lambda = 4\mu(\mu - 1) \in \mathbb{C}$ with $\text{Re } \mu \geq 1/2$, then every such solution has the form cP_{-n}^μ for some $c \in \mathbb{C}$.

We close this section by collecting some elementary properties of the PFM P_n^μ which will be needed in the sequel.

Remark 4.6 (Elementary properties of Poisson–Fourier modes) Let $\mu \in \mathbb{C}$, $n \in \mathbb{N}_0$ and $z, w \in \mathbb{D}$.

(a) The Poisson–Fourier modes are related to the hypergeometric function ${}_2F_1$ via

$$\begin{aligned} P_n^\mu(z, w) &= (-1)^n \binom{-\mu}{n} (1 - zw)^\mu w^n {}_2F_1(\mu, \mu + n; n + 1; zw) \\ (4.15) \quad &= (-1)^n \binom{-\mu}{n} w^n {}_2F_1\left(\mu, 1 - \mu; n + 1; \frac{zw}{zw - 1}\right). \end{aligned}$$

This is (4.7) together with (4.11) resp. (4.12).

(b) Using the series representation of ${}_2F_1$ functions, we see that

$$(4.16) \quad P_n^\mu(z, w) = (-1)^n (1 - zw)^\mu \sum_{k=0}^\infty \binom{-\mu}{k+n} \binom{-\mu}{k} z^k w^{k+n}.$$

If $-\mu = m \in \mathbb{N}_0$, then $P_{\pm n}^{-m} = 0$, whenever $n > m$. Otherwise,

$$(4.17) \quad P_n^{-m}(z, w) = (-1)^n \sum_{k=0}^{m-n} \binom{m}{k+n} \binom{m}{k} \frac{z^k w^{k+n}}{(1 - zw)^m}.$$

(c) The PFM are symmetric in the sense that $P_n^\mu(z, w) = P_{-n}^\mu(w, z)$. This allows us to simplify our proofs in the following: we will often prove identities for P_n^μ only which then implies the corresponding result for P_{-n}^μ .

(d) Remark 4.1 implies

$$(4.18) \quad \binom{\mu - 1}{n} P_{\pm n}^\mu = \binom{-\mu}{n} P_{\pm n}^{1-\mu}.$$

Remark 4.7 (Invariant representative functions and the finite-dimensional invariant eigenspaces) We consider \mathbb{D} as a symmetric space $\mathcal{M}(\mathbb{D})/\partial\mathbb{D}$ over its automorphism group $\mathcal{M}(\mathbb{D})$ with $\partial\mathbb{D} \cong U(1)$ acting by rotations. This yields an alternative way of deriving the restrictions to \mathbb{D} of those Poisson–Fourier modes from (4.17)

which are defined on all of Ω , using finite-dimensional representation theory. Such considerations are very much in the spirit of [21]. The group $\mathcal{M}(\mathbb{D})$ is isomorphic to the projective split unitary group

$$\text{PSU}(1,1) = \left\{ \begin{pmatrix} a & b \\ \bar{b} & \bar{a} \end{pmatrix} : |a|^2 - |b|^2 = 1 \right\} / \left\{ \pm \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \right\}$$

as a real Lie group. By Schur’s lemma, every irreducible representation $\pi: \mathcal{M}(\mathbb{D}) \rightarrow \text{GL}_n(\mathbb{C})$ as invertible $(n \times n)$ -matrices induces eigenfunctions of the Laplacian on $\text{PSU}(1,1)$ via

$$(4.19) \quad \pi_{k,\ell}: \text{PSU}(1,1) \rightarrow \mathbb{C}, \quad \pi_{k,\ell}(g) := \pi(g)_{k,\ell}, \quad 1 \leq k, \ell \leq n,$$

which are known as *representative functions* or *matrix elements*. For a systematic discussion, we refer to the textbook [9, Chapter III]. The corresponding eigenvalues may be computed from representation-theoretic data (see, e.g., [16, Proposition 10.6]). Note that in the literature, many results are formulated for general invariant differential operators or the corresponding Casimir elements and joint eigenfunctions thereof instead of the Laplacian, which is always invariant and thus constitutes a special case. However, the disk \mathbb{D} is a two-point homogeneous space, so all $\mathcal{M}(\mathbb{D})$ -invariant differential operators are polynomials in the Laplacian (see [20, Theorem 11]). Adapting [16, Example 4.10], one may parameterize the irreducible representations of $\text{PSU}(1,1)$ as follows: let $m \in \mathbb{N}_0$ and

$$V_m := \text{span} \{ z^k w^{m-k} \in \mathbb{C}[z, w] \mid 0 \leq k \leq m \}$$

the vector space of polynomials of total degree m with $\dim(V_m) = m + 1$. Then

$$\pi_m: \text{SU}(1,1) \rightarrow \text{GL}(V_m), \quad \pi_m \begin{pmatrix} a & b \\ \bar{b} & \bar{a} \end{pmatrix} p(z, w) := p(\bar{a}z - bw, -\bar{b}z + aw)$$

for $m \in \mathbb{N}$ constitutes a complete list of the irreducible finite-dimensional representations of $\text{SU}(1,1)$ (see also [16, Proposition 4.11 and Section 4.6]). Moreover, π_m descends to the quotient $\text{PSU}(1,1)$ if and only if m is even, providing a description of all irreducible finite-dimensional representations of $\text{PSU}(1,1)$. A computation then yields explicit formulas for (4.19), which completes the eigenvalue theory of the Laplacian on $\text{PSU}(1,1)$. Finally, parameterizing a copy of the rotation group $U(1) \subseteq \text{PSU}(1,1)$ by $\begin{pmatrix} i\eta & 0 \\ 0 & -i\bar{\eta} \end{pmatrix}$ with $\eta \in \partial\mathbb{D}$ yields that exactly the representative functions $(\pi_{2m})_{m,d}$ with $1 \leq d \leq 2m + 1$ are invariant under the action of $U(1)$ and thus pass to the quotient $\mathbb{D} \cong \text{PSU}(1,1)/U(1)$. By a computation, $(\pi_{2m})_{m,d} = P_{m-d}^{-m}$ with the exceptional Poisson–Fourier modes from (4.17).

Note that these considerations only recover the finite-dimensional invariant eigenspace, i.e., the case (E_0) in Theorem 2.10. It would be interesting to study whether this approach generalizes to the other invariant subspaces by incorporating suitable representations on infinite-dimensional spaces.

5 Spectral decomposition of eigenspaces

In this section, we show that for every rotationally invariant domain $D \subseteq \Omega$ containing the origin each holomorphic eigenfunction of the invariant Laplacian Δ_{zw} on D

has a unique representation as a *Poisson–Fourier series*, a doubly infinite series with Poisson–Fourier modes as building blocks. We shall also see that if $D = \Omega$ this series corresponds to a finite sum, and when D is one of the distinguished domains Ω_+ or Ω_- , then the series is one-sided infinite.

Theorem 5.1 *Let D be a rotationally invariant subdomain of Ω containing the origin, and let $f \in \mathcal{H}(D)$ be such that $\Delta_{z,w} f = \lambda f$ for some $\lambda \in \mathbb{C}$ of the form $\lambda = 4\mu(\mu - 1)$ with $\operatorname{Re} \mu \geq 1/2$. Then there are uniquely determined coefficients $c_n \in \mathbb{C}$ such that*

$$(5.1) \quad f = \sum_{n=-\infty}^{\infty} c_n P_n^\mu := \sum_{n=0}^{\infty} c_n P_n^\mu + \sum_{n=1}^{\infty} c_{-n} P_{-n}^\mu.$$

Here, both series converge absolutely and locally uniformly in D .

Proof (i) For each $n \in \mathbb{Z}$ and all $(z, w) \in D$, we consider

$$f_n(z, w) := \frac{1}{2\pi i} \int_{\partial \mathbb{D}} f\left(\eta z, \frac{w}{\eta}\right) \eta^{-(n+1)} d\eta.$$

Since D is rotationally invariant, f_n is well-defined. Clearly, f_n is holomorphic on D and n -homogeneous. By Theorem 4.4, there are complex numbers $c_n \in \mathbb{C}$ such that

$$f_n(z, w) = c_n P_{-n}^\mu(z, w).$$

(ii) We fix $(z, w) \in D$ and positive constants $r < 1 < R$ such that $(\eta z, w/\eta) \in D$ for all $r < |\eta| < R$. This is possible as D is a rotationally invariant domain. Then $\eta \mapsto f(\eta z, w/\eta)$ is holomorphic in the annulus $r < |\eta| < R$ and therefore has a representation as the Laurent series

$$f\left(\eta z, \frac{w}{\eta}\right) = \sum_{n=-\infty}^{\infty} f_n(z, w) \eta^n,$$

which converges locally uniformly in $r < |\eta| < R$. In particular,

$$f(z, w) = \sum_{n=-\infty}^{\infty} f_n(z, w), \quad (z, w) \in D.$$

This series converges, in fact, uniformly on each compact set $K \subseteq D$. In order to see this, let K be such a compact subset of D . We can then choose positive constants $r_1 < 1 < R_1$ such that $K_1 := \{(\eta z, w/\eta) : (z, w) \in K, r_1 \leq |\eta| \leq R_1\}$ is a compact subset of D , and we let $M_1 := \max\{|f(z, w)| : (z, w) \in K_1\}$. Note that

$$f_n(z, w) = \frac{1}{2\pi i} \int_{\partial \mathbb{D}_\rho} f\left(\eta z, \frac{w}{\eta}\right) \eta^{-(n+1)} d\eta$$

for all $n \in \mathbb{Z}$, all $(z, w) \in K$ and every $\rho \in [r, R]$. We therefore get

$$|f_n(z, w)| \leq M_1 \cdot \rho^{-n} \quad \text{for all } (z, w) \in K \text{ and all } \rho \in [r_1, R_1].$$

In particular, $|f_n(z, w)| \leq M_1 r_1^{-n}$ for all $n < 0$ and $|f_n(z, w)| \leq M_1 R_1^{-n}$ for all $n \geq 0$, and this ensures the absolute and uniform convergence of the two series

$$\sum_{n=-\infty}^{-1} f_n(z, w) \quad \text{and} \quad \sum_{n=0}^{\infty} f_n(z, w)$$

on the compact set K .

(iii) In view of Step (ii), the coefficients $f_n(z, w)$ are exactly the Laurent coefficients of $\eta \mapsto f(\eta z, w/\eta)$ in an annulus containing the unit circle and are thus uniquely determined by f . Hence, $f_n(z, w) = c_n P_{-n}^\mu(z, w)$ shows that the coefficients c_n are uniquely determined by f . ■

In fact, the previous proof provides the following more precise information.

Corollary 5.2 *Let D be a subdomain of Ω containing the origin, and let $f \in \mathcal{H}(D)$ be such that $\Delta_{zw} f = \lambda f$ for some $\lambda \in \mathbb{C}$.*

(i) *If $D = \Omega$, then $\lambda = 4m(m + 1)$ for some $m \in \mathbb{N}_0$ and there are uniquely determined coefficients $c_n \in \mathbb{C}$ such that*

$$f = \sum_{n=-m}^m c_n P_n^{m+1}.$$

(ii) *If $D = \Omega_+$, then $\lambda = 4m(m + 1)$ for some $m \in \mathbb{N}_0$ and there are uniquely determined coefficients $c_n \in \mathbb{C}$ such that*

$$f = \sum_{n=-\infty}^m c_n P_n^{m+1}.$$

(iii) *If $D = \Omega_-$, then $\lambda = 4m(m + 1)$ for some $m \in \mathbb{N}_0$ and there are uniquely determined coefficients $c_n \in \mathbb{C}$ such that*

$$f = \sum_{n=-m}^{\infty} c_n P_n^{m+1}.$$

Proof (i) If $f \in \mathcal{H}(\Omega)$, then the functions

$$f_n(z, w) := \frac{1}{2\pi i} \int_{\partial\mathbb{D}} f\left(\eta z, \frac{w}{\eta}\right) \eta^{-(n+1)} d\eta$$

are holomorphic and n -homogeneous on Ω and

$$f_n(z, w) = c_n P_{-n}^\mu(z, w)$$

with $\mu \in \mathbb{C}$ such that $\lambda = 4\mu(\mu - 1)$ and $\text{Re } \mu \geq 1/2$. Hence, $\lambda = 4m(m + 1)$ for some $m \in \mathbb{N}_0$ since otherwise P_{-n}^μ is not holomorphic on Ω by Theorem 4.4. If $\lambda = 4m(m + 1)$, then for each $|n| > m$, the function P_{-n}^{m+1} is not holomorphic on Ω again by Theorem 4.4, which forces $c_n = 0$ for those n . Parts (ii) and (iii) follow in the same way. ■

Since all PFM P_n^μ are holomorphic on Ω_* , it is natural to inquire whether the series (5.1) in Theorem 5.1 converges on some bigger domain than D . In Section 6, we shall see that, in general, this is not the case.

6 Comparison with the eigenvalue equation of the Laplacian on the unit disk and the Riemann sphere

In this section, we relate the spectral theory of the Laplacian on the unit disk \mathbb{D} resp. the Riemann sphere $\widehat{\mathbb{C}}$ with the spectral theory of the invariant Laplacian on Ω which we have developed so far. In particular, we show that all eigenfunctions of $\Delta_{\mathbb{D}}$ on \mathbb{D} resp. $\Delta_{\widehat{\mathbb{C}}}$ on $\widehat{\mathbb{C}}$ do have holomorphic extensions to eigenfunctions of Δ_{z^w} on the bidisk \mathbb{D}^2 resp. Ω . Our approach is similar to the one employed in [8, Section 1.6], which deals exclusively with the hyperbolic Laplacian $\Delta_{\mathbb{D}}$ on the unit disk \mathbb{D} . However, we need some fine properties of hypergeometric functions, in addition to those which have been employed in [8]. We begin with the following lemma which is crucial for our approach.

Lemma 6.1 *Let $\mu \in \mathbb{C}$ with $\operatorname{Re} \mu > 0$. Then*

$$(6.1) \quad \lim_{n \rightarrow \infty} {}_2F_1(\mu, 1 - \mu; n + 1; \omega) = 1$$

locally uniformly for $\omega \in \mathbb{C} \setminus [1, \infty)$.

Proof Fix $m \in \mathbb{N}_0$ such that $m > \operatorname{Re} \mu - 2$, and let $n \in \mathbb{N}_0$. Then by [11, Formula (11), p. 76] or [30, p. 84] there are complex numbers A_k , $k = 1, \dots, m$, such that

$${}_2F_1(\mu, 1 - \mu; n + 1; \omega) = 1 + \sum_{k=1}^m \frac{A_k}{(n+1) \dots (n+k)} \frac{\omega^k}{k!} + \rho_{m+1}(n, \omega),$$

where

(6.2)

$$\rho_{m+1}(n, \omega) := \tilde{\gamma}_n \int_0^1 \int_0^1 t^{1-\mu+m} (1-t)^{n+\mu-1} (1-st\omega)^{-\mu-1-m} (1-s)^m ds dt \cdot \omega^{m+1}$$

and

$$\tilde{\gamma}_n := \frac{\Gamma(n+1)\Gamma(\mu+m)}{\Gamma(1-\mu)\Gamma(\mu+n)\Gamma(\mu) m!}.$$

Note that our choice of the nonnegative integer m guarantees that the integral in (6.2) converges. To prove (6.1), it therefore suffices to show that $\rho_{m+1}(n, \omega) \rightarrow 0$ uniformly on every compact subset K of $\mathbb{C} \setminus [1, \infty)$ as $n \rightarrow \infty$. Fix such a compact set K . Obviously,

$$M_K := \min_{\omega \in K, 0 \leq s, t \leq 1} |1 - st\omega| > 0,$$

and since $-\operatorname{Re} \mu - 1 - m < 0$, we have

$$(6.3) \quad \begin{aligned} |\rho_{m+1}(n, \omega)| &\leq |\tilde{\gamma}_n| \int_0^1 t^{1-\operatorname{Re} \mu+m} (1-t)^{n+\operatorname{Re} \mu-1} dt \frac{|\omega|^{m+1}}{M_K^{m+1+\operatorname{Re} \mu}} \\ &= |\tilde{\gamma}_n| \frac{\Gamma(2+m-\operatorname{Re} \mu) \Gamma(n+\operatorname{Re} \mu)}{\Gamma(2+m+n)} \frac{|\omega|^{m+1}}{M_K^{m+1+\operatorname{Re} \mu}} \\ &= \gamma^* \frac{\Gamma(n+\operatorname{Re} \mu)}{\Gamma(n+1)} \frac{\Gamma(n+1)}{|\Gamma(n+\mu)|} \frac{\Gamma(n+1)}{\Gamma(2+m+n)} \frac{|\omega|^{m+1}}{M_K^{m+1+\operatorname{Re} \mu}}, \end{aligned}$$

where

$$\gamma^* := \frac{|\Gamma(\mu + m)| \Gamma(2 + m - \operatorname{Re} \mu)}{|\Gamma(1 - \mu)| |\Gamma(\mu)| m!}.$$

An application of Stirling’s formula

$$(6.4) \quad \lim_{n \rightarrow \infty} \frac{\Gamma(\alpha + n)}{\Gamma(n + 1)} e^{-(\alpha-1) \log(n+1)} \rightarrow 1, \quad \alpha \in \mathbb{C},$$

to each of the first three quotients in (6.3) shows that there is a constant $\gamma > 0$ depending only on μ, m and K such that

$$|\rho_{m+1}(n, \omega)| \leq \gamma \cdot (n + 1)^{\operatorname{Re} \mu - 1} (n + 1)^{1 - \operatorname{Re} \mu} (n + 1)^{-m-1} |\omega|^{m+1} = \gamma \cdot \left(\frac{|\omega|}{n + 1} \right)^{m+1}.$$

In particular, $\rho_{m+1}(n, \omega) \rightarrow 0$ as $n \rightarrow \infty$ uniformly for $\omega \in K$. This completes the proof of (6.1). ■

We are now in a position to prove Theorem 2.2.

Proof of Theorem 2.2 Let $f \in X_\lambda(\mathbb{D})$, and write $\lambda = 4\mu(\mu - 1)$ for some complex number $\mu \in \mathbb{C}$ with $\operatorname{Re} \mu \geq 1/2$. By [8, Theorem 1.6.18], there are uniquely determined coefficients $c_n \in \mathbb{C}$ such that

$$f(z) = \sum_{n=-\infty}^{\infty} c_n P_n^\mu(z, \bar{z});$$

the series converges absolutely and pointwise for each $z \in \mathbb{D}$. As it is shown in the proof of [8, Theorem 1.6.18], the coefficients c_n do have the additional property that

$$(6.5) \quad \sum_{n=-\infty}^{\infty} |c_n| r^{|n|} < \infty \quad \text{for all } 0 < r < 1.$$

We proceed to show that (6.5) and Lemma 6.1 together guarantee that the series

$$(6.6) \quad F(z, w) := \sum_{n=-\infty}^{\infty} c_n P_n^\mu(z, w)$$

converges locally uniformly for $(z, w) \in \mathbb{D}^2$, and hence defines a function $F \in \mathcal{H}(\mathbb{D}^2)$ with the required properties. The identity principle shows further that F is then uniquely determined.

It remains to prove the local uniform convergence of the series (6.6) in \mathbb{D}^2 . Fix $r \in (0, 1)$. We begin by noting that (4.11) and (4.7) lead to

$$(6.7) \quad \begin{aligned} |P_n^\mu(z, w)| &= \left| \binom{\mu}{|n|} \right| |F_{-n}^\mu(z, w)| \\ &\leq \frac{|\Gamma(\mu + |n|)|}{|\Gamma(\mu)| \Gamma(|n| + 1)} \left| {}_2F_1 \left(\mu, 1 - \mu; |n| + 1; \frac{zw}{zw - 1} \right) \right| r^{|n|} \end{aligned}$$

for all $|z|, |w| \leq r$. Since the Möbius transformation

$$\xi \mapsto \frac{\xi}{\xi - 1}$$

maps the unit disk \mathbb{D} conformally onto the half-plane $\{\zeta \in \mathbb{C} : \operatorname{Re} \zeta < 1/2\}$, the set

$$K := \left\{ \frac{zw}{zw-1} : |z| \leq r, |w| \leq r \right\}$$

is a compact subset of $\mathbb{C} \setminus [1, \infty)$. Thus, Lemma 6.1 implies

$${}_2F_1 \left(\mu, 1 - \mu; |n| + 1; \frac{zw}{zw-1} \right) \rightarrow 1 \quad \text{uniformly for } |z| \leq r, |w| \leq r$$

as $|n| \rightarrow \infty$. Combining this with Stirling's formula (6.4) we see from inequality (6.7) that there is a constant $\gamma > 0$ such that

$$(6.8) \quad |P_n^\mu(z, w)| \leq \gamma \cdot (|n| + 1)^{\operatorname{Re} \mu - 1} r^{|n|}$$

for all $|z| \leq r, |w| \leq r$ and every $n \in \mathbb{Z}$. This estimate together with (6.5) implies that the series (6.6) converges uniformly for $|z|, |w| \leq r$, as required. In particular, we have shown that the restriction map

$$\mathcal{R}_h : X_\lambda(\mathbb{D}^2) \rightarrow X_\lambda(\mathbb{D}), \quad \mathcal{R}_h(F)(z) := F(z, \bar{z}) \quad (z \in \mathbb{D})$$

is bijective. Since $X_\lambda(\mathbb{D}^2)$ and $X_\lambda(\mathbb{D})$ are both Fréchet spaces with respect to the topology of locally uniform convergence on \mathbb{D}^2 resp. \mathbb{D} (see [38, Corollary 1 to Theorem 4.2.4] for the fact that $X_\lambda(\mathbb{D})$ is a Fréchet space) and the restriction map \mathcal{R}_h is obviously continuous, its inverse is continuous as well by the Open Mapping Theorem. This completes the proof of Theorem 2.2. ■

Theorem 6.2 (C^2 -eigenfunctions of $\Delta_{\widehat{\mathbb{C}}}$ on $\widehat{\mathbb{C}}$ vs. holomorphic eigenfunctions of Δ_{zw} on Ω) *Let $\lambda \in \mathbb{C}$. Then the following are equivalent:*

- (i) *There is a function $f \in C^2(\widehat{\mathbb{C}})$ such that $\Delta_{\widehat{\mathbb{C}}} f = \lambda f$ on $\widehat{\mathbb{C}}$.*
- (ii) *There is a function $F \in \mathcal{H}(\Omega)$ such that $\Delta_{zw} F = \lambda F$ on Ω .*
- (iii) $\lambda = 4m(m+1)$ for some $m \in \mathbb{N}_0$.

Theorem 6.2 is a special case of Theorem 2.6.

Proof The implication (ii) \Rightarrow (iii) is Corollary 5.2, and (iii) \Rightarrow (ii) is Theorem 4.4. Clearly, (ii) implies (i), so we only need to prove that (i) implies (ii). Accordingly, we write $\lambda = 4\mu(\mu-1)$ with $\mu \in \mathbb{C}$ and $\operatorname{Re} \mu \geq 1/2$. For $n \in \mathbb{Z}$, consider

$$f_n(z) := \frac{1}{2\pi i} \int_{\partial \mathbb{D}} f(\eta z) \eta^{-(n+1)} d\eta = \frac{1}{2\pi} \int_0^{2\pi} f(e^{it} z) e^{-int} dt.$$

Then $f_n \in C^2(\widehat{\mathbb{C}})$ is n -homogeneous and $\Delta_{\widehat{\mathbb{C}}} f_n = 4\mu(\mu-1)f_n$ on $\widehat{\mathbb{C}}$. Arguing in a similar way as in the proof of Theorem 4.4, we see that there is a constant $c_n \in \mathbb{C}$ such that

$$f_n(z) = c_n \binom{-\mu}{|n|} {}_2F_1 \left(\mu, 1 - \mu; |n| + 1; \frac{|z|^2}{1 + |z|^2} \right) \cdot \begin{cases} \frac{z^{|n|}}{|z|^{|n|}}, & \text{if } n \geq 0, \\ (-z)^{|n|}, & \text{if } n \leq 0. \end{cases}$$

Therefore, the behavior of f_n as $z \rightarrow \infty$ depends essentially on the value ${}_2F_1(\mu; 1 - \mu; |n| + 1; 1)$. It is well-known, see [3, Theorem 2.1.3], that

$$\lim_{x \rightarrow 1^-} \frac{{}_2F_1(\mu, 1 - \mu; |n| + 1; 1)}{-\log(1 - x)} = \frac{1}{\Gamma(\mu)\Gamma(1 - \mu)} \quad \text{if } n = 0,$$

and, see [3, Theorem 2.2.2],

$${}_2F_1(\mu, 1 - \mu; |n| + 1; 1) = \frac{\Gamma(|n| + 1)\Gamma(|n|)}{\Gamma(|n| + 1 - \mu)\Gamma(|n| + \mu)} \quad \text{if } n \neq 0.$$

Hence, if $n = 0$, then f_n is defined at the point ∞ only if $c_0 = 0$ or $\mu \in \mathbb{N}$. If $n \neq 0$, then ${}_2F_1(\mu, 1 - \mu; |n| + 1; 1) = 0$ if and only if $\mu = |n| + 1 + k$ for some $k \in \mathbb{N}_0$. Thus, $f_n : \widehat{\mathbb{C}} \rightarrow \mathbb{C}$ is defined at the point ∞ only if $c_n = 0$ or $\mu \in \mathbb{N}$, $\mu > |n|$. We conclude that either $c_n = 0$ for every $n \in \mathbb{Z}$ and then $f \equiv 0$, or $c_n \neq 0$ for at least one $n \in \mathbb{Z}$ and then $\lambda = 4m(m + 1)$ for some $m \in \mathbb{N}_0$. In the latter case, we see that $c_n \neq 0$ forces $m + 1 = \mu = |n| + 1 + k$ for some $k \in \mathbb{N}_0$, so $m = |n| + k$. In particular, $c_n = 0$ for all $|n| > m$, so $f(z) = F(z, -\bar{z})$ for

$$F(z, w) = \sum_{n=-m}^m c_n P_n^\mu(z, w) \in \mathcal{H}(\Omega). \quad \blacksquare$$

Remark 6.3 We see that the spectrum of the hyperbolic Laplacian $\Delta_{\mathbb{D}}$ is \mathbb{C} whereas the spectrum of the spherical Laplacian $\Delta_{\widehat{\mathbb{C}}}$ is notably smaller as it only consists of the scalars $4m(m + 1)$ for $m \in \mathbb{N}_0$. Furthermore, by Theorem 6.2, every eigenfunction of the spherical Laplacian can be extended to the whole domain Ω . This is different from the hyperbolic case where the extension to \mathbb{D}^2 provided by Theorem 2.2 is maximal at least for the category of rotationally invariant domains as the following example shows.

Example 6.4 Fix $\mu \in \mathbb{C}$ with $\text{Re } \mu \geq \frac{1}{2}$, let $(z_0, w_0) \in \Omega \setminus \mathbb{D}^2$ with $z_0 w_0 \notin [1, \infty)$, and define

$$b_n := \left(\frac{\Gamma(\mu + |n|)}{\Gamma(\mu)|n|!} {}_2F_1(\mu; 1 - \mu; |n| + 1) \frac{z_0 w_0}{z_0 w_0 - 1} \right)^{-1}.$$

Note that asymptotically

$$|b_n| \sim (|n| + 1)^{1 - \text{Re } \mu} \quad (|n| \rightarrow \infty).$$

This follows from Lemma 6.1 and Stirling’s formula (6.4). In view of (6.8), this shows that the series

$$F(z, w) = \sum_{n=-\infty}^{\infty} b_n P_n^\mu(z, w)$$

converges locally uniformly in \mathbb{D}^2 . Hence, $F \in \mathcal{H}(\mathbb{D}^2)$ and, since every function P_n^μ is an eigenfunction of $\Delta_{z,w}$, we have $\Delta_{z,w} F = 4\mu(\mu - 1)F$ in \mathbb{D}^2 . In particular,

$$f(z) := F(z, \bar{z}) = \sum_{n=-\infty}^{\infty} b_n P_n^\mu(z, \bar{z})$$

is of class $C^2(\mathbb{D})$ and an eigenfunction of $\Delta_{\mathbb{D}}$. By Theorem 2.2, $F \in \mathcal{H}(\mathbb{D}^2)$ is the uniquely determined function in $\mathcal{H}(\mathbb{D}^2)$ such that $F(z, \bar{z}) = f(z)$ in \mathbb{D} . Now assume $F \in \mathcal{H}(D)$, where D is a rotationally invariant domain such that $\mathbb{D}^2 \not\subseteq D \subseteq \Omega$ and $(z_0, w_0) \in D \setminus \mathbb{D}^2$. By Theorem 5.1, there are coefficients $(\tilde{b}_n)_{n \in \mathbb{Z}} \subseteq \mathbb{C}$ such that

$$F(z, w) = \sum_{n=-\infty}^{\infty} \tilde{b}_n P_n^{\mu}(z, w)$$

for all $(z, w) \in D$. Since the coefficients are uniquely determined, we conclude $\tilde{b}_n = b_n$. This is a contradiction because

$$F(z_0, w_0) = \sum_{n=0}^{\infty} w_0^n + \sum_{n=1}^{\infty} z_0^n$$

is a divergent series since $|z_0| \geq 1$ or $|w_0| \geq 1$.

Corollary 6.5 *Let $\lambda \in \mathbb{C}$ and D be a rotationally invariant subdomain of Ω such that $\mathbb{D}^2 \not\subseteq D$. Then there exists a function $f \in X_{\lambda}(\mathbb{D}^2)$ that cannot be analytically continued to D .*

7 Poisson–Fourier modes and the Möbius group

Recall the Möbius group $\mathcal{M}(\mathbb{D}^2)$ of the bidisk from (2.5). In order to give a characterization of the Möbius invariant eigenspaces of $\Delta_{z,w}$ (e.g., proving Theorem 2.10), it will turn out that all we need to understand are precompositions of PFM with automorphisms in $\mathcal{M}(\mathbb{D}^2)$, that is,

$$P_n^{\mu} \circ T \text{ with } \operatorname{Re} \mu \geq 1/2 \text{ and } T \in \mathcal{M}(\mathbb{D}^2).$$

Moreover, since $\mathcal{M}(\mathbb{D}^2)$ is closely related to the Möbius group \mathcal{M} of Ω from (2.4), we can use the following description of \mathcal{M} .

Lemma 7.1 *The group \mathcal{M} is generated by the flip map $\mathcal{F}(u, v) := (1/v, 1/u)$ and the mappings*

$$(7.1) \quad T_{z,w}(u, v) := \left(\frac{z-u}{1-wu}, \frac{w-v}{1-zv} \right) \quad \text{and} \quad \rho_{\gamma}(u, v) := \left(\gamma u, \frac{1}{\gamma} v \right), \quad (u, v) \in \Omega$$

with $(z, w) \in \Omega \cap \mathbb{C}^2$ and $\gamma \in \mathbb{C}^*$. More precisely, for every $T \in \mathcal{M}$, there exist $z, w \in \mathbb{C}$ and $\gamma \in \mathbb{C}^*$ such that

$$T = \rho_{\gamma} \circ T_{z,w} \quad \text{or} \quad T = \rho_{\gamma} \circ T_{z,w} \circ \mathcal{F}.$$

Remark 7.2 Lemma 7.1 is exactly Lemma 2.2 in [19] where we have replaced the automorphisms $\Phi_{z,w}$ by the $T_{z,w} = \Phi_{z,w} \circ \rho_{-1}$ automorphisms in (7.1). The reason will become apparent in (7.3) and (7.4). Essentially, the automorphisms $T_{z,w}$ in (7.1) are precisely the automorphisms of Ω interchanging a given point $(z, w) \in \Omega \cap (\mathbb{C} \times \mathbb{C})$ with $(0, 0)$ – instead of only sending $(0, 0)$ to (z, w) (see [32, Section 2.3] for more details on this).

Therefore, understanding precompositions of a PFM P_n^{μ} with elements $T \in \mathcal{M}(\mathbb{D}^2)$ breaks down to understanding precompositions with the above generators. Note

that the mappings \mathcal{F} and $T_{z,w}$ for $w \neq \bar{z}$ are not elements of $\mathcal{M}(\mathbb{D}^2)$. However, the precompositions with these mappings still make sense for those PFM defined on all of Ω , that is, in view of Theorem 4.4, the precompositions

$$P_n^{-m} \circ T \text{ with } m \in \mathbb{N}_0 \text{ and } T \in \mathcal{M}$$

are well-defined. In the case that $T = \rho_\gamma$, we have $P_n^\mu \circ T = \gamma^{-n} P_n^\mu$ by the $(-n)$ -homogeneity of the PFM for all $\mu \in \mathbb{C}$. Next, if $T = \mathcal{F}$, it is easily seen by direct verification based on (4.17) that

$$(P_n^{-m} \circ \mathcal{F})(z, w) = (-1)^m P_n^{-m}(z, w)$$

for $m \in \mathbb{N}_0$ and $|n| \leq m$. In order to treat the case that $T = T_{z,w}$ resp. $T = T_{z,\bar{z}}$, some preliminary observations are useful. First, every automorphism $T_{z,w}$ is induced by a Möbius transformation $\psi_{z,w}$ of the form

$$(7.2) \quad \psi_{z,w}(u) := \frac{z - u}{1 - wu},$$

in the sense that $T_{z,w}(u, v) = (\psi_{z,w}(u), 1/\psi_{z,w}(1/v))$.

Remark 7.3 Choosing $w = \bar{z}$ in (7.2) yields all self-inverse automorphisms of \mathbb{D} except for the identity, and, similarly, the choice $w = -\bar{z} \in \mathbb{D}$ yields all self-inverse rigid motions of $\widehat{\mathbb{C}}$.

Further, we will employ the definition of the PFM via integrals from (4.10). By Remark 4.3, when considering points $(z, \bar{z}) \in \mathbb{D}^2$, this definition coincides with the n th Fourier mode of the μ -power of the Poisson kernel on the unit disk. The Poisson kernel on \mathbb{D} satisfies the well-known properties

$$(7.3) \quad P(z, \bar{z}; \xi) = \frac{\psi'_{z,\bar{z}}(\xi)}{\psi_{z,\bar{z}}(\xi)} \xi$$

and

$$(7.4) \quad P(\psi_{z,\bar{z}}(u), \overline{\psi_{z,\bar{z}}(u)}; \psi_{z,\bar{z}}(\xi)) = \frac{P(u, \bar{u}; \xi)}{P(z, \bar{z}; \xi)},$$

where $z, u \in \mathbb{D}$, $\xi \in \partial\mathbb{D}$, and $\psi_{z,\bar{z}} \in \mathcal{M}(\mathbb{D})$ (see [38, Theorem 3.3.5]).

Remark 7.4 One can define the generalized Poisson kernel P even on $\Omega \times \partial\mathbb{D}$. Then P is never zero, and for each $\xi \in \partial\mathbb{D}$, the function $P(\cdot; \xi)$ is meromorphic in the sense of [13, Chapter VI.2]. The identities (7.3) and (7.4) then take the form

$$P(\psi_{z,w}^{-1}(0), 1/\psi_{z,w}^{-1}(\infty); \xi) = \frac{\psi'_{z,w}(\xi)}{\psi_{z,w}(\xi)} \xi = P(z, w; \xi),$$

$$P\left(\psi_{z,w}(u), \frac{1}{\psi_{z,w}(1/v)}; \psi_{z,w}(\xi)\right) = \frac{P(u, v; \xi)}{P(\psi_{z,w}^{-1}(0), 1/\psi_{z,w}^{-1}(\infty); \xi)} = \frac{P(u, v; \xi)}{P(z, w; \xi)},$$

where $(z, w) \in \Omega \cap \mathbb{C}^2$, $(u, v) \in \Omega$, and $\xi \in \partial\mathbb{D}$.

With these preparations, we are now in a position to analyze the precompositions of P_n^μ with elements of $\mathcal{M}(\mathbb{D}^2)$ resp. \mathcal{M} . We start with the simplest case, $P_0^{m+1} \circ T_{u,v}$.

Proposition 7.5 *Let $m \in \mathbb{N}_0$, $(z, w) \in \Omega$, and $(u, v) \in \Omega \cap \mathbb{C}^2$. Then*

$$(7.5) \quad (P_0^{m+1} \circ T_{u,v})(z, w) = (P_0^{-m} \circ T_{u,v})(z, w) = \sum_{j=-m}^m P_{-j}^{m+1}(u, v) P_j^{-m}(z, w).$$

Proof The way of reasoning is as follows: first, fix $z, u \in \mathbb{D}$, and let $v := \bar{u}$. Then $\psi_{u, \bar{u}} \in \mathcal{M}(\mathbb{D})$ which means, in particular, that $T_{u, \bar{u}}(z, \bar{z}) \in \{(t, \bar{t}) : t \in \mathbb{D}\}$. By the two variable identity principle, Lemma 3.1, if we can show (7.5) in this special case, then, keeping $u \in \mathbb{D}$ fixed, (7.5) also holds for points $(z, w) \in \Omega$ since both sides of the equation are holomorphic functions in Ω . Next, fix $(z, w) \in \Omega$ and note that both sides of (7.5) are holomorphic as functions of $(u, v) \in \Omega \cap \mathbb{C}^2$. Assuming (7.5) for $v = \bar{u} \in \mathbb{D}$ then implies the claim, again by Lemma 3.1.

It remains to show (7.5) for $w = \bar{z}$ and $v = \bar{u}$ with $z, u \in \mathbb{D}$. For this purpose, we compute

$$\begin{aligned} (P_0^{m+1} \circ T_{u, \bar{u}})(z, \bar{z}) &\stackrel{(4.18)}{=} (P_0^{-m} \circ T_{u, \bar{u}})(z, \bar{z}) \\ &= \frac{1}{2\pi} \int_0^{2\pi} P\left(\psi_{u, \bar{u}}(z), \frac{1}{\psi_{u, \bar{u}}(1/\bar{z})}; e^{it}\right)^{-m} dt \\ &\stackrel{(7.4)}{=} \frac{1}{2\pi} \int_0^{2\pi} P(z, \bar{z}; \psi_{u, \bar{u}}^{-1}(e^{it}))^{-m} P(u, \bar{u}; \psi_{u, \bar{u}}^{-1}(e^{it}))^m dt \\ &\stackrel{\psi_{u, \bar{u}}(e^{is})=e^{it}}{=} \frac{1}{2\pi} \int_0^{2\pi} P(z, \bar{z}; e^{is})^{-m} P(u, \bar{u}; e^{is})^{m+1} ds. \end{aligned}$$

In the last step, we used (7.3). Now, we can interpret the resulting integral as the inner product $\langle \cdot, \cdot \rangle$ on $L^2([0, 2\pi], \mathbb{C})$, the space of square integrable functions $f : [0, 2\pi] \rightarrow \mathbb{C}$. Using Parseval's identity in (P), we obtain

$$\begin{aligned} (P_0^{-m} \circ T_{u, \bar{u}})(z, \bar{z}) &= \left\langle P(u, \bar{u}; e^{is})^{m+1}, \overline{P(z, \bar{z}; e^{is})^{-m}} \right\rangle \\ &\stackrel{(P)}{=} \sum_{j=-\infty}^{\infty} \left\langle P(u, \bar{u}; e^{it})^{m+1}, e^{-ijt} \right\rangle \left\langle e^{-ijs}, \overline{P(z, \bar{z}; e^{is})^{-m}} \right\rangle \\ &= \sum_{j=-\infty}^{\infty} \frac{1}{(2\pi)^2} \int_0^{2\pi} P(u, \bar{u}; e^{it})^{m+1} e^{ijt} dt \int_0^{2\pi} P(z, \bar{z}; e^{is})^{-m} e^{-ijs} ds \\ &= \sum_{j=-m}^m P_{-j}^{m+1}(u, \bar{u}) P_j^{-m}(z, \bar{z}), \end{aligned}$$

where the series reduces to a finite sum because P_j^{-m} vanishes for $|j| > m$ (see Remark 4.6(b)). ■

The previous proof can be modified to establish the following result.

Theorem 7.6 Let $\mu \in \mathbb{C}$, $\operatorname{Re} \mu \geq 1/2$, $n \in \mathbb{N}_0$, and $u \in \mathbb{D}$. Then

$$(7.6a) \quad \begin{aligned} & (P_n^\mu \circ T_{u, \bar{u}})(z, w) \\ &= \sum_{j=-\infty}^{\infty} \left(\sum_{k=0}^n \sum_{\ell=0}^{\infty} (-1)^{k-n-\ell} \binom{n}{k} \binom{-n}{\ell} u^\ell \bar{u}^k P_{j-k+n+\ell}^{1-\mu}(u, \bar{u}) \right) P_{-j}^\mu(z, w) \end{aligned}$$

and

$$(7.6b) \quad \begin{aligned} & (P_{-n}^\mu \circ T_{u, \bar{u}})(z, w) \\ &= \sum_{j=-\infty}^{\infty} \left(\sum_{k=0}^n \sum_{\ell=0}^{\infty} (-1)^{n-k+\ell} \binom{n}{k} \binom{-n}{\ell} u^k \bar{u}^\ell P_{j+k-n-\ell}^{1-\mu}(u, \bar{u}) \right) P_{-j}^\mu(z, w) \end{aligned}$$

for all $(z, w) \in \mathbb{D}^2$. Moreover, both series converge locally uniformly and absolutely w.r.t. (z, w) in \mathbb{D}^2 .

Proof Using (7.3) and (7.4), we compute

$$\begin{aligned} (P_n^\mu \circ T_{u, \bar{u}})(z, \bar{z}) &= \frac{1}{2\pi} \int_0^{2\pi} P(z, \bar{z}; e^{is})^\mu P(u, \bar{u}; e^{is})^{1-\mu} \psi_{u, \bar{u}}(e^{is})^{-n} ds \\ &= \frac{1}{2\pi} \int_0^{2\pi} P(z, \bar{z}; e^{is})^\mu P(u, \bar{u}; e^{is})^{1-\mu} (1 - \bar{u} e^{is})^n (u - e^{is})^{-n} ds. \end{aligned}$$

Note that taking complex powers is unproblematic since the appearing Poisson kernels are positive real quantities. Applying Lemma 3.1 and, additionally, the generalized binomial theorem leads to

$$\begin{aligned} & (P_n^\mu \circ T_{u, \bar{u}})(z, w) \\ &= \sum_{k=0}^n \sum_{\ell=0}^{\infty} \binom{n}{k} \binom{-n}{\ell} (-1)^{k-n+\ell} \bar{u}^k u^\ell \frac{1}{2\pi} \int_0^{2\pi} P(z, w; e^{is})^\mu P(u, \bar{u}; e^{is})^{1-\mu} e^{is(k-n-\ell)} ds. \end{aligned}$$

Interpreting the above integral as a $L^2([0, 2\pi], \mathbb{C})$ inner product as it was done in the proof of Proposition 7.5 shows (7.6a). A similar computation yields (7.6b).

Let $K \subseteq \mathbb{D}^2$ be compact. It remains to prove that

$$\sum_{k=0}^n \sum_{\ell=0}^{\infty} \left| \binom{n}{k} \binom{-n}{\ell} \bar{u}^k u^\ell \right| \sum_{j=-\infty}^{\infty} \left| P_{j-k+n+\ell}^{1-\mu}(u, \bar{u}) \right| \sup_{(z,w) \in K} \left| P_{-j}^\mu(z, w) \right| < \infty.$$

We have

$$\left| P_{j-k+n+\ell}^{1-\mu}(u, \bar{u}) \right| \leq \frac{1}{2\pi} \int_0^{2\pi} |P(u, \bar{u}; e^{it})^{1-\mu}| dt := M_u < \infty$$

as $u \in \mathbb{D}$ is fixed. Since K is compact, we find a nonnegative real number $r < 1$ such that $|z|, |w| \leq r$ for all $(z, w) \in K$. Hence, using (6.8), we can ensure that

$$\sum_{j=-\infty}^{\infty} \sup_{(z,w) \in K} \left| P_{-j}^\mu(z, w) \right| \leq \sum_{j=-\infty}^{\infty} \gamma_r \cdot (|j| + 1)^{\operatorname{Re} \mu - 1} r^{|j|} := N_{r, \mu} < \infty,$$

where $\gamma_r > 0$ depends on r (and μ). Thus,

$$\begin{aligned} & \sum_{k=0}^n \sum_{\ell=0}^{\infty} \left| \binom{n}{k} \binom{-n}{\ell} \bar{u}^k u^\ell \right| \sum_{j=-\infty}^{\infty} \left| P_{j-k+n+\ell}^{1-\mu}(u, \bar{u}) \right| \sup_{(z,w) \in K} \left| P_{-j}^{\mu}(z, w) \right| \\ & \leq \sum_{k=0}^n \sum_{\ell=0}^{\infty} \left| \binom{n}{k} \binom{-n}{\ell} \bar{u}^k u^\ell \right| \cdot M_u \cdot N_{r,\mu} < \infty \end{aligned}$$

since the binomial series over k and ℓ converge absolutely. ■

Let $m, n \in \mathbb{N}_0, n \leq m$. For $\mu = -m$, the j -series in (7.6) terminate. Moreover, in this case, when replacing (u, \bar{u}) by $(u, v) \in \mathbb{D}^2$ in (7.6), a similar argument as in the proof of Theorem 7.6 shows that the series in (7.6) are absolutely and locally uniformly convergent with respect to (u, v) in \mathbb{D}^2 , too. Further, we know that $P_{\pm n}^{-m} \circ T_{u,v} \in \mathcal{H}(\Omega)$. Thus, both sides of (7.6a) resp. (7.6b) (with (u, \bar{u}) replaced by (u, v)) define holomorphic functions w.r.t. $(u, v) \in \mathbb{D}^2$, which allows us to apply the two variable identity principle, Lemma 3.1. We obtain the following.

Corollary 7.7 *Let $m, n \in \mathbb{N}_0, n \leq m, (z, w) \in \Omega$ and $(u, v) \in \mathbb{D}^2$. Then*

$$\begin{aligned} & (P_n^{-m} \circ T_{u,v})(z, w) \\ & = \sum_{j=-m}^m \left(\sum_{k=0}^n \sum_{\ell=0}^{\infty} (-1)^{k-n-\ell} \binom{n}{k} \binom{-n}{\ell} u^\ell v^k P_{j-k+n+\ell}^{m+1}(u, v) \right) P_{-j}^{-m}(z, w) \end{aligned}$$

and

$$\begin{aligned} & (P_{-n}^{-m} \circ T_{u,v})(z, w) \\ & = \sum_{j=-m}^m \left(\sum_{k=0}^n \sum_{\ell=0}^{\infty} (-1)^{n-k+\ell} \binom{n}{k} \binom{-n}{\ell} u^k v^\ell P_{j+k-n-\ell}^{m+1}(u, \bar{u}) \right) P_{-j}^{-m}(z, w). \end{aligned}$$

In particular, $P_n^{-m} \circ T_{u,v} \in \text{span}\{P_j^{-m} : j = -m, \dots, m\} \subseteq \mathcal{H}(\Omega)$.

8 Proof of Theorems 2.6 and 2.10: invariant eigenspaces

Using the theory we have developed so far, we are now able to give an explicit description of the eigenspaces of the Laplace operator $\Delta_{z,w}$ on \mathbb{D}^2 resp. on Ω . This will allow us to conclude Theorems 2.6 and 2.10 afterwards.

Theorem 8.1 (Möbius invariant subspaces of $X_\lambda(\mathbb{D}^2)$) *Let $\lambda \in \mathbb{C}$.*

- (NE) *If $\lambda \neq 4m(m+1)$ for all $m \in \mathbb{N}_0$, then $X_\lambda(\mathbb{D}^2)$ has no nontrivial Möbius invariant subspaces.*
- (E) *If $\lambda = 4m(m+1)$ for some $m \in \mathbb{N}_0$, then $X_\lambda(\mathbb{D}^2)$ has precisely three nontrivial Möbius invariant subspaces $X_\lambda^+(\mathbb{D}^2), X_\lambda^-(\mathbb{D}^2), X_\lambda^0(\mathbb{D}^2)$ of which exactly one, say $X_\lambda^0(\mathbb{D}^2)$, is finite dimensional of dimension $2m+1$. Explicitly, these spaces are given by*

$$(8.1a) \quad X_\lambda^+(\mathbb{D}^2) = \text{clos}_{\mathbb{D}^2} \left(\text{span}\{P_n^{m+1} : -m \leq n < \infty\} \right),$$

$$(8.1b) \quad X_{\lambda}^{-}(\mathbb{D}^2) = \text{clos}_{\mathbb{D}^2} \left(\text{span}\{P_n^{m+1} : -\infty < n \leq m\} \right),$$

$$(8.1c) \quad X_{\lambda}^0(\mathbb{D}^2) = \text{span}\{P_n^{m+1} : -m \leq n \leq m\},$$

where $\text{clos}_{\mathbb{D}^2}$ denotes the closure with respect to the topology of locally uniform convergence on \mathbb{D}^2 .

Note that Rudin [37] gives an explicit characterization of the three nontrivial subspaces in the second case of his theorem (see Theorem 2.1) too. We will discuss the similarities and the differences between Rudin’s and our approach in Remark 8.5.

Three crucial ingredients that we will use for the proof of Theorem 8.1 are the representation of eigenfunctions in terms of PFM from Theorem 5.1, the pullback formula for PFM from Theorem 7.6 and the differential operators D^+ and D^- defined by

$$(8.2) \quad \begin{aligned} D^+ f(z, w) &:= \partial_z f(z, w) - w^2 \partial_w f(z, w), \\ D^- f(z, w) &:= \partial_w f(z, w) - z^2 \partial_z f(z, w), \end{aligned}$$

for $f \in \mathcal{H}(D)$ and any subdomain $D \subseteq \mathbb{C}^2$. We note an alternative description of these operators based on the automorphisms from (7.1).

Lemma 8.2 *Let $D \subseteq \mathbb{C}^2$ be a subdomain, $f \in \mathcal{H}(D)$. For all $(z, w) \in D$, it then holds that*

$$(8.3) \quad \begin{aligned} D^+ f(z, w) &= \partial_u (f \circ T_{u,0} \circ \rho_{-1}) \Big|_{u=0} (z, w), \\ D^- f(z, w) &= \partial_v (f \circ T_{0,v} \circ \rho_{-1}) \Big|_{v=0} (z, w). \end{aligned}$$

What makes these operators useful is that they act by shifts of the Fourier index on the PFM.

Lemma 8.3 *The operators D^+ and D^- commute with Δ_{zw} . Furthermore, for $\mu \in \mathbb{C}$ and $n \in \mathbb{Z}$, it holds that*

$$D^+ P_n^\mu = (\mu - n - 1) P_{n+1}^\mu \quad \text{and} \quad D^- P_n^\mu = (\mu + n + 1) P_{n-1}^\mu.$$

Proof The operator D^+ commutes with Δ_{zw} in view of the latter’s \mathcal{M} -invariance, (8.3) and commutativity of partial derivatives. Moreover, D^+ lowers homogeneity in the sense of (4.1) by one degree. Since P_n^μ is a $(-n)$ -homogeneous eigenfunction of Δ_{zw} to the eigenvalue $4\mu(\mu - 1)$, this implies that $D^+ P_n^\mu$ is a $-(n + 1)$ -homogeneous eigenfunction of Δ_{zw} to the same eigenvalue. Hence, Theorem 4.4 yields $D^+ P_n^\mu = c P_{n+1}^\mu$ for some $c \in \mathbb{C}$. Applying (8.2) and evaluating (4.16) in $(z, w) = (0, 1)$ yields $c = \mu - n - 1$. Note that one has to distinguish the cases $n \geq 0$ and $n < 0$. The second equality in (8.3) may be derived analogously. ■

Remark 8.4 One may thus regard the operators D^+ and D^- as ladder operators and together with their commutator

$$\frac{1}{2} (D^- \circ D^+ - D^+ \circ D^-) = z \partial_z - w \partial_w,$$

which is the Euler vector field corresponding to the homogeneity (4.1), they generate the Lie algebra of \mathcal{M} . By (8.3), all three operators may, moreover, be regarded as fundamental vector fields of the natural action of \mathcal{M} as automorphisms of Ω .

Proof (of Theorem 8.1) Write $\lambda = 4\mu(\mu - 1)$ for $\mu \in \mathbb{C}$, $\operatorname{Re} \mu \geq 1/2$, and let Y be a nontrivial Möbius invariant subspace of $X_\lambda(\mathbb{D}^2)$. Our approach is as follows: we first show that Y contains a PFM P_n^μ for some $n \in \mathbb{Z}$. If the assumption of (NE) holds, then this will imply that $P_n^\mu \in Y$ for all $n \in \mathbb{Z}$. Conversely, if the assumption of (E) holds for $m \in \mathbb{N}_0$, the existence of P_n^{-m} for some $n \in \mathbb{Z}$ will imply that certain other PFM need to be contained in Y , too. Here, we will need to distinguish three cases which will lead to the three specific spaces described on the right-hand side of (8.1). This shows that there are at most three possible choices for Y . Then, in a second step, we show that the three spaces found in the first step are in fact Möbius invariant, which then proves the existence of exactly three nontrivial proper Möbius invariant subspaces.

Step 1:

- (i) Using (8.2), it follows from Y being closed and Möbius invariant that the differential quotients $D^+f, D^-f \in Y$. Therefore, Y is invariant with respect to D^+ and D^- , and, by iteration, $(D^+)^k f, (D^-)^k f \in Y$ for every $k \in \mathbb{N}$.
- (ii) Since Y is Möbius invariant, Y is, in particular, rotationally invariant. Thus, for every $n \in \mathbb{Z}$, we have $f_n \in Y$, where f_n is defined by

$$f_n(z, w) := \frac{1}{2\pi i} \int_{\partial\mathbb{D}} f\left(\eta z, \frac{w}{\eta}\right) \eta^{-(n+1)} d\eta.$$

By Theorem 4.4, f_n is a multiple of the PFM P_n^μ . Moreover, by Theorem 5.1, there are $c_n \in \mathbb{C}$ for all $n \in \mathbb{Z}$ depending on f such that f has the representation

$$f = \sum_{n=-\infty}^{\infty} c_n P_n^\mu.$$

Since $Y \neq \{0\}$, we may assume $f \neq 0$, i.e. $c_n \neq 0$ for some $n \in \mathbb{Z}$. Thus, the above observations imply $P_n^\mu \in Y$ for some $n \in \mathbb{Z}$.

- (iii) By Part (ii), there exists $P_n^\mu \in Y$ for some $n \in \mathbb{Z}$, and by Part (i), we conclude $(D^+)^k P_n^\mu, (D^-)^k P_n^\mu \in Y$ for every $k \in \mathbb{N}$. Lemma 8.3 shows that these functions are multiples of PFM again, so $(\mu - n - k)_k P_{n+k}^\mu, (\mu + n - k)_k P_{n-k}^\mu \in Y$ for every $k \in \mathbb{N}$. Recall that the prefactors denote (rising) Pochhammer symbols (see (4.5)).

For $\mu \notin \mathbb{Z}$, the factors $(\mu \pm n - k)_k$ never vanish. In this case, $P_n^\mu \in Y$ for every $n \in \mathbb{Z}$. Using Theorem 5.1 again, this implies $Y = X_\lambda(\mathbb{D}^2)$ which proves (NE).

For $\mu \in \mathbb{N}$, the factors $(\mu \pm n - k)_k$ vanish for appropriate $k \in \mathbb{N}$. This leads to the dichotomy in case (E): assume $\mu = m + 1 \in \mathbb{N}$, and $P_n^{m+1} \in Y$. Then necessarily $P_k^{m+1} \in Y$ for $-m \leq k \leq m$. Since shifting P_k^{m+1} with D^\pm eventually produces the zero function, for $q < -m$ or $q > m$ the function P_q^{m+1} does not necessarily need to be contained in Y . However, if $P_q^{m+1} \in Y$ for some $q < -m$, then since P_q^{m+1} can be shifted to each of these functions without producing the zero function, it must hold that $P_p^{m+1} \in Y$ for $-\infty < p \leq m$. The same argument works for $q > m$. In summary, these considerations show that every nontrivial closed Möbius invariant subspace of $X_\lambda(\mathbb{D}^2)$ contains one of the three spaces:

$$\begin{aligned} Y_+ &:= \operatorname{clos}_{\mathbb{D}^2} \left(\operatorname{span}\{P_n^{m+1} : -m \leq n < \infty\} \right), \\ Y_- &:= \operatorname{clos}_{\mathbb{D}^2} \left(\operatorname{span}\{P_n^{m+1} : -\infty < n \leq m\} \right) \text{ and} \\ Y_0 &:= \operatorname{span}\{P_n^{m+1} : -m \leq n \leq m\}. \end{aligned}$$

Step 2: It remains to show that Y_+ , Y_- , and Y_0 are in fact Möbius invariant. Let $\mu - 1 = m \in \mathbb{N}_0$, $n \in \mathbb{Z}$, and $z, w, u \in \mathbb{D}$.

- (iv) Assume $n \geq 0$. Theorem 7.6 shows that $(P_n^{m+1} \circ T_{u,\bar{u}})(z, w)$ is a series of scalar multiples of $P_k^{m+1}(z, w)$ with $-m \leq k < \infty$. To see this recall that by Remark 4.6, we have $P_n^{-m} \equiv 0$ if $|n| > m$. Moreover, Theorem 7.6 also implies that this series converges locally uniformly with respect to (z, w) on \mathbb{D}^2 . Thus, $P_n^{m+1} \circ T_{u,\bar{u}} \in Y_+$. The same argument applies to $P_{-n}^{m+1} \circ T_{u,\bar{u}}$.
- (v) Part (iv) implies that every linear combination

$$\sum_{n=-m}^N c_n P_n^{m+1} \circ T_{u,\bar{u}}$$

belongs to Y_+ . Since, by definition, every element in Y_+ is a series of the form $\sum_{n=-m}^\infty c_n P_n^{m+1}$, Möbius invariance of Y_+ follows. The same argument applies to Y_- .

- (vi) Now, consider $-m \leq n \leq m$. In this case, P_n^{m+1} is a multiple of P_n^{-m} by (4.18), and Corollary 7.7 shows that $(P_n^{m+1} \circ T_{u,\bar{u}})(z, w)$ can be expressed as linear combination of $P_k^{-m}(z, w)$ with $-m \leq k \leq m$. This shows Möbius invariance of Y_0 . ■

Remark 8.5 (a) In his proof of Theorem 2.1, Rudin employs exclusively a specific differential operator A (see [37, Formula (5), p. 143]), which in our terminology can be defined by

$$A := D^+ + D^-.$$

The operator $A^\sharp := i(D^+ - D^-)$ is “conjugate” to A in the sense that

$$A - iA^\sharp = 2D^+ \quad \text{and} \quad A + iA^\sharp = 2D^-.$$

Note the analogy with the Wirtinger derivatives ∂_z and $\partial_{\bar{z}}$, which satisfy

$$\partial_x - i\partial_y = 2\partial_z \quad \text{and} \quad \partial_x + i\partial_y = 2\partial_{\bar{z}}.$$

We note that in a related, but different context, namely in [36], Rudin also considers the operators D^+ and D^- . He denotes them by Q and \bar{Q} (see [36, Section 4.1, Formula (1)]).

- (b) Step 1 of our proof closely follows the argumentation of Rudin’s proof (Steps 1–3 in his labeling). However, because of the usage of his A -operator (see Part (a)), which does not respect homogeneity when applied to a PFM, Rudin additionally needs to consider the projections on n -homogeneous components in his proof. In our approach based on the D^\pm -operators, this issue does not arise, since these operators lower resp. raise the degree of homogeneity by exactly one.
- (c) Step 2 of our proof, where we made heavy use of Proposition 7.5 resp. Theorem 7.6, is different from Rudin’s approach (Step 4 in his labeling). In short, Rudin makes a series expansion in the disk automorphism parameter and uses again that Möbius invariant subspaces are invariant under repeated application of A (see Part (a)).
- (d) The proof of Theorem 8.1 heavily relies on the PFM. In fact, the usage of n -homogeneous functions is also crucial in the original proof of Theorem 2.1

in [37], although Rudin does not explicitly define these functions but works with projections on the subspaces of n -homogeneous eigenfunctions.

We are now in a position to prove Theorems 2.6 and 2.10.

Proof (of Theorem 2.10) Let $\lambda \in \mathbb{C}$ and $Y \subseteq \mathcal{H}(\mathbb{D}^2)$ be a nontrivial Möbius invariant subspace of $X_\lambda(\mathbb{D}^2)$. Let, first, $Y = X_\lambda(\mathbb{D}^2)$ be the full eigenspace. Every PFM is holomorphic at least on Ω_* . By rotational invariance of \mathbb{D}^2 and Theorem 5.1, this implies that $Y \cap \mathcal{H}(\Omega_*)$ is dense in Y , i.e., (NE) holds. Assume now that Y is a proper subspace of $X_\lambda(\mathbb{D}^2)$. By Theorem 8.1, this is only possible if λ is an exceptional eigenvalue. Furthermore,

$$Y \in \{X_\lambda^+(\mathbb{D}^2), X_\lambda^-(\mathbb{D}^2), X_\lambda^0(\mathbb{D}^2)\}.$$

Combining (8.1a), (8.1b), resp. (8.1c) with Corollary 5.2 yields (E_+) , (E_-) resp. (E_0) . We have already discussed the additional statements in the latter case in Theorem 8.1.

Finally, Corollary 6.5 shows that none of the density statements (E_+) , (E_-) , and (NE) may be improved to holomorphic extensibility of every element. ■

Proof (of Theorem 2.6) The equivalence of (i), (ii), and (iii) was already established in Theorem 6.2. Let $\lambda = 4m(m + 1)$ for some $m \in \mathbb{N}_0$. By (8.1c), we then have $\dim(X_\lambda(\Omega)) = 2m + 1$. Investing moreover Theorem 4.4 yields (a). Remark 6.3 proves the inclusion “ \subseteq ” in (2.2). Conversely, restricting an eigenfunction $F \in X_\lambda(\Omega)$ to the rotated diagonal $\{(z, -\bar{z}) : z \in \widehat{\mathbb{C}}\}$ yields an eigenfunction of $\Delta_{\widehat{\mathbb{C}}}$. This completes the proof. ■

The finite-dimensional spaces $X_{4m(m+1)}(\Omega)$ can also be characterized in terms of the PFM P_0^{-m} only.

Corollary 8.6 *Let $m \in \mathbb{N}_0$. The function P_0^{-m} is cyclic in $X_{4m(m+1)}(\Omega)$ with respect to the natural action of \mathcal{M} by pullbacks. That is,*

$$(8.4) \quad X_{4m(m+1)}(\Omega) = \text{span} \{P_0^{-m} \circ T \mid T \in \mathcal{M}\}.$$

In addition,

$$(8.5) \quad X_{4m(m+1)}(\Omega) = \text{span} \{P_0^{-m} \circ T_{u,\bar{u}} \mid u \in \mathbb{D}\} = \text{span} \{P_0^{-m} \circ T_{u,-\bar{u}} \mid u \in \mathbb{C}\}.$$

Proof We begin by showing the first equality in (8.5), which implies (8.4). Recall that, by (8.1c), the set $\{P_j^{-m} : -m \leq j \leq m\}$ constitutes a basis of $X_{4m(m+1)}(\Omega)$. By (7.5),

$$P_0^{-m} \circ T_{u,\bar{u}} = \sum_{j=-m}^m P_{-j}^{m+1}(u, \bar{u}) P_j^{-m}$$

for every $u \in \mathbb{D}$. We interpret the claim as a change of basis from $\{P_j^{-m} \mid -m \leq j \leq m\}$ to $\{P_0^{-m} \circ T_{u_k, \bar{u}_k} \mid -m \leq k \leq m\}$. Therefore, we have to find suitable points $u_{-m}, \dots, u_m \in \mathbb{D}$ such that the coefficient matrix

$$(P_j^{m+1}(u_k, \bar{u}_k))_{-m \leq j, k \leq m}$$

is invertible. However, it is true in general that given linearly independent complex valued functions F_1, \dots, F_M for some $M \in \mathbb{N}$ on a set containing at least M elements, it is possible to find the same number of points x_1, \dots, x_M in the same set such that the

matrix $(F_j(x_k))_{1 \leq j, k \leq M}$ is invertible (see [15, Proof of Proposition 7.28]). The second equality in (8.5) follows similarly. ■

9 Poisson–Fourier modes and Peshl–Minda operators

In [19], the classical Peshl–Minda differential operators, which were introduced by Peshl [34] and studied, e.g., by [2, 17, 24, 25, 31, 39, 40, 42], have been extended to differential operators acting on holomorphic functions defined on subdomains of $\Omega \cap \mathbb{C}^2$. It is the purpose of this section to put these (generalized) Peshl–Minda operators into the context of the present paper. In particular, we relate the Poisson–Fourier modes with the Peshl–Minda operators.

We first briefly recall the definition from [19]. Let U be an open subset of $\Omega \cap \mathbb{C}^2$ and $f \in C^\infty(U)$. The Peshl–Minda derivative $D^{m,n}f$ at the point $(z, w) \in U$ is defined by

$$D^{m,n}f(z, w) := \left. \frac{\partial^{m+n}}{\partial u^m \partial v^n} (f \circ T_{z,w} \circ \rho_{-1})(u, v) \right|_{(u,v)=(0,0)}.$$

We write $D_z^n := D^{n,0}$ and $D_w^n := D^{0,n}$ and refer to these operators as *pure* Peshl–Minda operators. Comparing with (8.3), we note that the roles of (z, w) and (u, v) have been swapped, which yields the operators D_z^1 and D_w^1 instead of D^+ and D^- , respectively.

Given a linear mapping $T: V \rightarrow W$ we write $\ker(T) := \{x \in V: Tx = 0\}$ for its kernel. We observe that the pure Peshl–Minda operators reproduce the generalized Poisson kernel from (4.9) in the following way.

Lemma 9.1 *Let $\mu \in \mathbb{C}$, $(z, w) \in \mathbb{D}^2$, $\xi \in \partial\mathbb{D}$, and $k \in \mathbb{N}_0$. Then*

$$(9.1a) \quad D_z^k P(z, w; \xi)^\mu = (\mu)_k P(z, w; \xi)^\mu (\psi_{z,w}(\xi))^{-k}$$

$$(9.1b) \quad D_w^k P(z, w; \xi)^\mu = (\mu)_k P(z, w; \xi)^\mu (\psi_{z,w}(\xi))^k$$

with $\psi_{z,w}$ from (7.2). In particular, if $\mu = -m \in (-\mathbb{N}_0)$, then

$$P^{-m} \in \ker(D_z^{m+1}) \cap \ker(D_w^{m+1}).$$

Proof If $w = \bar{z}$, this is most easily proved by an induction on k and the help of Proposition 3.8 in [19]. Then, the general result follows from Lemma 3.1. ■

We note in passing that Lemma 9.1 for $w = \bar{z}$ says that up to multiplication with a unimodular constant the (classical) Poisson kernel $z \mapsto P(z, \bar{z}; \xi)$ is a joint eigenfunction of the (classical) Peshl–Minda operators studied, e.g., in [24]. Remarkably, the Peshl–Minda operators also act as weighted shifts when applied to the zeroth PFM.

Proposition 9.2 *For $\mu \in \mathbb{C}$ and $n \in \mathbb{N}_0$, we have*

$$(-\mu + 1)_n \cdot P_n^\mu = D_z^n P_0^\mu \quad \text{and} \quad (-\mu + 1)_n \cdot P_{-n}^\mu = D_w^n P_0^\mu.$$

Proof Let $z \in \mathbb{D}$. We prove the claim for points (z, \bar{z}) first from which then follows the result by the identity principle (Lemma 3.1). We compute

$$\begin{aligned} D^{n,0}P_0^\mu(z, \bar{z}) &= \frac{1}{2\pi} \int_0^{2\pi} D^{n,0}P(z, \bar{z}; e^{it})^\mu dt \\ &\stackrel{(9.1a)}{=} \frac{1}{2\pi} \int_0^{2\pi} (\mu)_n P(z, \bar{z}; e^{it})^\mu (\psi_{z,\bar{z}}(e^{it}))^{-n} dt \\ &\stackrel{(7.4)}{=} \frac{(\mu)_n}{2\pi} \int_0^{2\pi} P(z, \bar{z}; \psi_{z,\bar{z}}(e^{it}))^{-\mu} (\psi_{z,\bar{z}}(e^{it}))^{-n} dt \\ &\stackrel{7.3}{=} \frac{(\mu)_n}{2\pi} \int_0^{2\pi} P(z, \bar{z}; e^{it})^{1-\mu} e^{-int} dt \\ &= (\mu)_n P_n^{1-\mu}(z, \bar{z}) \stackrel{(4.18)}{=} (-\mu + 1)_n P_n^\mu(z, \bar{z}). \end{aligned}$$

The computation for $D^{0,n}$ and P_{-n}^μ is analogous. ■

Remark 9.3 Since Poisson–Fourier modes are closely related to hypergeometric functions, see (4.15), an alternative proof can be given by induction using Gauss’ contiguous identities, more precisely with equation 15.2.24 and with a combination of equations 15.2.1 and 15.2.6 in [1].

In view of [19, Corollary 4.4], we find another connection between the PFM and the pure Peschl–Minda derivatives. Let $\lambda = 4m(m + 1)$, $m, n \in \mathbb{N}_0$. It is true that

$$P_n^{-m} \in \ker(D_w^k) \quad \text{and} \quad P_{-n}^{-m} \in \ker(D_z^k) \quad \text{for all } k > m.$$

Proposition 9.4 Let $M \in \mathbb{N}_0$. Then

$$(9.2) \quad \{P_n^{-m} : m \in \mathbb{N}_0, n \in \mathbb{Z}, m \leq M, |n| \leq m\} = \ker(D_z^{M+1}) \cap \ker(D_w^{M+1}).$$

Proof Denote the left-hand side of (9.2) by X_M . Note that $\ker(D_z^{M+1}) \cap \ker(D_w^{M+1})$ and X_M are finite-dimensional spaces. Moreover, X_M is the direct sum

$$X_M = \bigoplus_{m=0}^M X_{4m(m+1)}^0(\mathbb{D}^2)$$

and $\dim X_{4m(m+1)}^0(\mathbb{D}^2) = 2m + 1$ by Theorem 8.1. Therefore, we compute on the one hand

$$\dim X_M = \sum_{m=0}^M (2m + 1) = (M + 1)^2.$$

On the other hand, by [19, Corollary 4.4], we have

$$(9.3) \quad \ker(D_z^{M+1}) \cap \ker(D_w^{M+1}) = \text{span} \left\{ \frac{z^j w^k}{(1 - zw)^M} : 0 \leq j, k \leq M \right\}.$$

Consequently, $\dim(\ker(D_z^{M+1}) \cap \ker(D_w^{M+1})) = (M + 1)^2$. Moreover, combining (4.17) with (9.3) yields $X_M \subseteq (\ker(D_z^M) \cap \ker(D_w^M))$. This proves (9.2) by linear algebra. ■

Remark 9.5 Note that both the PFM P_n^{-m} and the generators from (9.3) may be understood as holomorphic functions on all of Ω . It thus makes sense to speak of

$$\bigoplus_{M \in \mathbb{N}_0} \left(\ker(D_z^{M+1}) \cap \ker(D_w^{M+1}) \right)$$

as the subspace of $\mathcal{H}(\Omega)$ consisting of all finite linear combinations of elements in the kernels of the Pöschl–Minda differential operators. We can even go one step further than Proposition 9.4 and conclude that as sets

$$\begin{aligned} & \text{clos}_\Omega \left(\text{span} \{ P_n^{-m} : m \in \mathbb{N}_0, n \in \mathbb{Z}, |n| \leq m \} \right) \\ &= \text{clos}_\Omega \left(\bigoplus_{M \in \mathbb{N}_0} \left(\ker(D_z^{M+1}) \cap \ker(D_w^{M+1}) \right) \right) \\ &= \mathcal{H}(\Omega), \end{aligned}$$

where the last equality follows from a combination of (9.3) with [18, Eq. (4.1) and Corollary 4.8]. This observation leads to the question whether the PFM form a Schauder basis of $\mathcal{H}(\Omega)$. The second author answers this question affirmatively in [32]. Perhaps similar results hold for $\mathcal{H}(\Omega_\pm)$ and the corresponding Poisson–Fourier modes from Theorem 4.4.

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