# A NEW PROXIMITY FUNCTION GENERATING THE BEST KNOWN ITERATION BOUNDS FOR BOTH LARGE-UPDATE AND SMALL-UPDATE INTERIOR-POINT METHODS 

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(Received 16 December, 2006; revised 6 August, 2007)


#### Abstract

Interior-Point Methods (IPMs) are not only very effective in practice for solving linear optimization problems but also have polynomial-time complexity. Despite the practical efficiency of large-update algorithms, from a theoretical point of view, these algorithms have a weaker iteration bound with respect to small-update algorithms. In fact, there is a significant gap between theory and practice for large-update algorithms. By introducing self-regular barrier functions, Peng, Roos and Terlaky improved this gap up to a factor of $\log n$. However, checking these self-regular functions is not simple and proofs of theorems involving these functions are very complicated. Roos et al. by presenting a new class of barrier functions which are not necessarily self-regular, achieved very good results through some much simpler theorems. In this paper we introduce a new kernel function in this class which yields the best known complexity bound, both for large-update and small-update methods.


2000 Mathematics subject classification: primary 90C05; secondary 90C51.
Keywords and phrases: linear optimization, interior-point method, primal-dual method, kernel function, large-update, small-update, polynomial complexity.

## 1. Introduction

We deal with the standard linear optimization problem

$$
\begin{equation*}
\min \left\{c^{\top} x: A x=b, x \geq 0\right\} \tag{1.1}
\end{equation*}
$$

where $A \in R^{m \times n}$ is a real $m \times n$ matrix with $\operatorname{rank}(A)=m, b \in R^{m}$ and $c \in R^{n}$. The dual problem of (1.1) is given by

$$
\begin{equation*}
\max \left\{b^{\top} y: A^{\top} y+s=c, s \geq 0\right\} \tag{1.2}
\end{equation*}
$$

[^0]Kernel functions play an important role in defining new search directions in primaldual interior point algorithms for solving linear optimization problems. A kernel function is a univariate strictly convex function which is defined for all positive reals $t$ and is minimal at $t=1$ where the minimal value equals zero. In other words $\psi(t): D \rightarrow R_{+}$with $R_{++} \subseteq D$, is a kernel function when it satisfies

$$
\psi^{\prime}(1)=\psi(1)=0, \quad \psi^{\prime \prime}(t)>0
$$

Moreover, $\psi(t)$ has the barrier property, that is, $\psi(t)$ goes to infinity if either $t \rightarrow 0$ or $t \rightarrow \infty$. Note that the above properties imply that $\psi(t)$ is completely determined by its second derivative

$$
\psi(t)=\int_{1}^{t} \int_{1}^{\xi} \psi^{\prime \prime}(\zeta) d \zeta d \xi
$$

This kernel function may be extended to a positive $n$-dimensional vector $\mu$ by

$$
\begin{equation*}
\Psi(\nu)=\sum_{i=1}^{n} \psi\left(v_{i}\right) \tag{1.3}
\end{equation*}
$$

yielding the scaled barrier function $\Psi(\nu)$. Note that the barrier function $\Psi(\nu)$ is nonnegative, and zero if and only if $v$ is a vector of ones. Therefore, the value of the barrier function can be considered as a measure for the closeness of $x$ and $(y, s)$ to the $\mu$-centres of (1.1) and (1.2). Hence some authors also call $\Psi(\nu)$ a proximity function. In the next section we briefly describe how any such barrier function defines a primal-dual interior-point method. The iteration bound for so-called large-update methods is obtained by showing that each iteration decreases $\Psi(v)$ by a sufficient amount. Table 1 gives some examples of kernel functions that have been analyzed so far, and the complexity results for the corresponding algorithms.

Note that all kernel functions in this table depend on a parameter and the mentioned iteration bounds in the table occur by choosing this parameter as indicated in the third column of the table. The first kernel function, the so-called self-regular function, was introduced and analyzed by Peng et al. Also the second kernel function is self regular and is the special case of the self regular function $\gamma_{p, q}(t)$, for $p=1$. The third function is not self regular and has been proposed and analyzed by Bai et al. A surprising feature of this kernel function is that it is finite for $t=0$, a property which separates it from self-regular functions, because self-regular functions become unbounded when $t$ approaches zero. Also Bai et al. introduced the following kernel function in [2]:

$$
\frac{t^{2}-1}{2}+\frac{e^{1 / t}-e}{e}
$$

TABLE 1. Kernel functions with the best known iteration bound for large-update methods. The iteration bound for all kernel functions is $O(\sqrt{n} \log n) \log (n / \epsilon)$.

| Kernel function | Parameter | References |
| :--- | :---: | :---: |
| $\frac{t^{2}-1}{2}+\frac{t^{1-q}-1}{q-1}, q>1$ | $q=\frac{1}{2} \log n$ | $[5,6]$ |
| $\frac{t^{2}-1}{2}+\frac{t^{1-q}-1}{q(q-1)}-\frac{q-1}{q}(t-1), q>1$ |  |  |
| $\frac{t^{2}-1}{2}+\frac{1}{\sigma}\left(e^{\sigma(1-1)}-1\right), \sigma>0$ | $\sigma=O(\log n)$ | $[1]$ |
| $\frac{t^{2}-1}{2}+\int_{1}^{t} e^{q(1 / \xi-1)} d \xi, q \geq 1$ | $q=O(\log (1+n))$ | $[3]$ |

They showed that the iteration bound for the corresponding algorithm is $O\left(\sqrt{n}(\log n)^{2} \log (n / \epsilon)\right)$. This bound is a factor $\log n$ worse than the bound in Table 1 . In this paper we introduce a new kernel function as follows:

$$
\begin{equation*}
\psi(t)=\frac{t^{2}-1}{2}+\frac{e^{q(1 / t-1)}-1}{q}, \quad q \geq 1 . \tag{1.4}
\end{equation*}
$$

We show that this kernel function yields the best-known iteration bound for largeupdate methods, that is, $O(\sqrt{n} \log n) \log (n / \epsilon)$. Figure 1 depicts the graph of $\psi(t)$.


Figure 1. The graph of $\psi$ for $q=4$.
The paper is organized as follows. In Section 2, we briefly recall how a given kernel function defines a primal-dual interior-point algorithm. In Section 3, we describe the
simple conditions on the kernel function that define a general class of kernel functions introduced by Bai et al. in [2]. Then we show that the new kernel function introduced in (1.4) satisfies these conditions. In what follows, we use the general scheme for analyzing the generic algorithm, as presented in [2]. We obtain the iteration bounds for both large-update and small-update methods based on a new kernel function in Section 4 and finally, Section 5 contains some concluding remarks.

## 2. The generic primal-dual interior-point algorithm

It is well known that finding an optimal solution of (1.1) and (1.2) is equivalent to solving the following system:

$$
\begin{cases}A x=b, & x \geq 0 \\ A^{\top} y+s=c, & s \geq 0 \\ x s=0 . & \end{cases}
$$

The basic idea of the primal-dual IPMs is to replace the third equation, the socalled complementarity condition, by the parameterized equation $x s=\mu e$, with $\mu>0$. Thus we consider the following system:

$$
\begin{cases}A x=b, & x \geq 0 \\ A^{\top} y+s=c, & s \geq 0 \\ x s=\mu e . & \end{cases}
$$

We assume that the primal and dual problem in (1.1) and (1.2) satisfies the interiorpoint condition (IPC), that is, there exists $\left(x^{0}, s^{0}, y^{0}\right)$ such that

$$
A x^{0}=b, \quad x^{0}>0, \quad A^{\top} y^{0}+s^{0}=c, \quad s^{0}>0 .
$$

It is well known that the IPC can be assumed without loss of generality. By using the self-dual embedding model, we will assume that $x^{0}=s^{0}=e[4,7]$. If $\operatorname{rank}(A)=m$ and the IPC holds, then for each $\mu>0$ the above parameterized system has a unique solution. We denote this solution as $(x(\mu), y(\mu), s(\mu))$ and call $x(\mu)$ the $\mu$-centre of (1.1) and $(y(\mu), s(\mu))$ the $\mu$-centre of (1.2). The set of $\mu$-centres (with $\mu$ running through all positive real numbers) gives a homotopy path, which is called the central path of (1.1) and (1.2). If $\mu \rightarrow 0$ then the limit of the central path exists and since the limit point satisfies the complementarity condition, the limit yields to optimal solutions for (1.1) and (1.2). If $(x(\mu), y(\mu), s(\mu))$ is known for some positive $\mu$, then we decrease $\mu$ to $\mu:=(1-\theta) \mu$, for some fixed $\theta \in(0,1)$ and solve the following
system:

$$
\left\{\begin{align*}
A \Delta x & =0  \tag{2.1}\\
A^{\top} \Delta y+\Delta s & =0 \\
s \Delta x+x \Delta s & =\mu e-x s
\end{align*}\right.
$$

This system uniquely defines a search direction ( $\Delta x, \Delta s, \Delta y$ ). This direction approximates the next $\mu$-centre. Hence all IPMs follow the central path approximately. Now we define

$$
v=\sqrt{\frac{x s}{\mu}} .
$$

One can easily check that System (2.1) which defines the search direction can be rewritten as follows:

$$
\left\{\begin{align*}
\bar{A} d_{x} & =0,  \tag{2.2}\\
\bar{A}^{\top} \Delta y+d_{s} & =0, \\
d_{x}+d_{s} & =v^{-1}-v,
\end{align*}\right.
$$

where $d_{x}=\nu \Delta x / x, d_{s}=\nu \Delta s / s$ and $\bar{A}=A V^{-1} X$, with $V=\operatorname{diag}(\nu)$ and $X=$ $\operatorname{diag}(x)$. The third equation in (2.2) is called the scaled centring equation. The right-hand side in the scaled centring equation equals minus the gradient of the basic logarithmic barrier function, that is,

$$
\Psi_{c}(v)=\sum_{i=1}^{n}\left(\frac{v_{i}^{2}-1}{2}-\log v_{i}\right)
$$

The basic idea in IPMs is to replace the scaled barrier function $\Psi_{c}(\nu)$ by an arbitrary strictly convex function $\Psi(\nu)$, such that $\Psi(\nu)$ is minimal at $\nu=e$ with $\Psi(e)=0$. Thus System (2.2) converts to

$$
\left\{\begin{aligned}
\bar{A} d_{x} & =0 \\
\bar{A}^{\top} \Delta y+d_{s} & =0 \\
d_{x}+d_{s} & =-\Delta \Psi(\nu)
\end{aligned}\right.
$$

Since the vectors $d_{x}$ and $d_{s}$ belong to the null and row spaces of the matrix $\bar{A}$, these vectors are orthogonal. By taking a suitable step size $\alpha \in(0,1)$, these search directions construct a new triple ( $x_{+}, y_{+}, s_{+}$) with

$$
x_{+}=x+\alpha \Delta x, \quad y_{+}=y+\alpha \Delta y, \quad s_{+}=s+\alpha \Delta s
$$

We repeat the procedure until we find an iterate in a certain neighbourhood of $(x(\mu), y(\mu), s(\mu))$. Then $\mu$ is again reduced by the factor $1-\theta$ and we apply Newton's method targeting the new $\mu$-centres, and so on. This process is repeated until $\mu$ is small enough and at this stage we have found an $\epsilon$-solution of the problems (1.1) and (1.2). The generic form of the algorithm is shown in Figure 2.

## Input:

A threshold parameter $\tau>0$;
An accuracy parameter $\epsilon>0$;
A fixed barrier update parameter $\theta, 0<\theta<1$;
begin

```
\(x:=e ; s:=e ; \mu:=1 ;\)
while \(n \mu>\epsilon\) do
begin
    \(\mu:=(1-\theta) \mu ;\)
    while \(\Psi(\nu)>\tau\) do
                begin
                    \(x_{+}:=x+\alpha \Delta x ;\)
                        \(y_{+}:=y+\alpha \Delta y\);
            \(\nu:=\sqrt{\frac{x s}{\mu}} ;\)
            end
        end
    end
```

Figure 2. A generic primal-dual algorithm for Lo

## 3. A general class of the kernel functions

In [2] Bai, El Ghami and Roos introduced a general class of kernel functions by using the following conditions:

$$
\begin{align*}
t \psi^{\prime \prime}(t)+\psi^{\prime}(t)>0, & t<1,  \tag{3.1}\\
t \psi^{\prime \prime}(t)-\psi^{\prime}(t)>0, & t>1,  \tag{3.2}\\
\psi^{\prime \prime \prime}(t)<0, & t>0,  \tag{3.3}\\
2 \psi^{\prime \prime}(t)^{2}-\psi^{\prime}(t) \psi^{\prime \prime \prime}(t)>0, & t<1,  \tag{3.4}\\
\psi^{\prime \prime}(t) \psi^{\prime}(\beta t)-\beta \psi^{\prime}(t) \psi^{\prime \prime \prime}(\beta t)>0, & t>1, \beta>1 . \tag{3.5}
\end{align*}
$$

The kernel function $\psi(t)$ is called eligible if it satisfies (3.1) and (3.3) to (3.5). It was shown in [2] that (3.2) and (3.3) imply (3.5). So a kernel function $\psi(t)$ is
eligible if it satisfies conditions (3.1) to (3.4). First we show that the kernel function introduced in (1.4) satisfies these conditions. The first three derivatives of $\psi(t)$ are shown in the following table:

Table 2. $\psi$ and its derivatives.

| $\psi(t)$ | $\frac{t^{2}-1}{2}+\frac{e^{q / t}-e^{q}}{q e^{q}}$ |
| :---: | :---: |
| $\psi^{\prime}(t)$ | $t-\frac{e^{q(1 / t-1)}}{t^{2}}$ |
| $\psi^{\prime \prime}(t)$ | $1+\frac{q+2 t}{t^{4}} e^{q(1 / t-1)}$ |
| $\psi^{\prime \prime \prime}(t)$ | $-\frac{q^{2}+6 q t+6 t^{2}}{t^{6}} e^{q(1 / t-1)}$ |

Since for $0<t<1$ and $q \geq 1$

$$
t \psi^{\prime \prime}(t)+\psi^{\prime}(t)=2 t+\left[\frac{q+2 t}{t^{3}}-\frac{1}{t^{2}}\right] e^{q(1 / t-1)}=2 t+\left(\frac{q+t}{t^{3}}\right) e^{q(1 / t-1)}>0
$$

Condition (3.1) is satisfied. Furthermore, for $t>1$

$$
t \psi^{\prime \prime}(t)-\psi^{\prime}(t)=\frac{q+3 t}{t^{3}} e^{q(1 / t-1)}>0
$$

So (3.2) is satisfied as well. Also from Table 2 it is obvious that (3.3) is satisfied. Finally for satisfying (3.4), it is seen that

$$
\begin{aligned}
2 \psi^{\prime \prime}(t)^{2} & -\psi^{\prime}(t) \psi^{\prime \prime \prime}(t) \\
& =\frac{2\left[t^{4}+(q+2 t) e^{q(1 / t-1)}\right]^{2}+e^{q(1 / t-1)}\left(q^{2}+6 q t+6 t^{2}\right)\left(t^{3}-e^{q(1 / t-1)}\right)}{t^{8}}
\end{aligned}
$$

If we set $A=e^{q(1 / t-1)}$, after some elementary reduction we have

$$
2 \psi^{\prime \prime}(t)^{2}-\psi^{\prime}(t) \psi^{\prime \prime \prime}(t)=\frac{A^{2}\left[2 t^{2}+2 q t+q^{2}\right]+A\left[14 t^{2}+10 q t+q^{2} t^{3}\right]+2 t^{8}}{t^{8}}
$$

On the other hand, because $0<t<1$, it follows that $1 / t-1>0$ and for $q \geq 1$ we have

$$
A=e^{q(1 / t-1)}>1
$$

So we have $2 \psi^{\prime \prime}(t)^{2}-\psi^{\prime}(t) \psi^{\prime \prime \prime}(t)>0$ and hence condition (3.4) holds. Therefore $\psi(t)$ lies in the general class introduced in [2]. Having such a kernel function, we can now construct the barrier function $\Psi(\nu)$ by (1.3). The function $\Psi(\nu)$ not only serves to define a search direction, but also as a measure of closeness of the current iterates to the $\mu$-centre. We use the norm-based proximity measure $\delta(\nu)$ defined by

$$
\delta(\nu)=\frac{1}{2}\|\nabla \Psi(\nu)\|=\frac{1}{2} \sqrt{\sum_{i=1}^{n}\left(\psi^{\prime}\left(v_{i}\right)\right)} .
$$

Since $\Psi(\nu)$ is strictly convex and minimal at $\nu=e$ we have

$$
\Psi(\nu)=0 \quad \Longleftrightarrow \quad \delta(\nu)=0 \quad \Longleftrightarrow \quad \nu=e .
$$

In other words, the proximity measure $\delta(\nu)$ is zero if and only if the current iterates are the $\mu$-centres. For any kernel function in the mentioned class, iteration bounds for both small-update and large-update methods can be obtained by using the following scheme:

Step 1. Solve the equation $-\psi^{\prime}(t) / 2=s$ to get $\rho(s)$, the inverse function of $-\psi^{\prime}(t) / 2, t \in(0,1]$. If the equation is hard to solve, derive a lower bound for $\rho(s)$.
Step 2. Calculate the decrease of $\Psi(\nu)$ in terms of $\delta$ for the default step size $\tilde{\alpha}=1 / \psi^{\prime \prime}(\rho(2 \delta))$ from $f(\tilde{\alpha})=\delta^{2} / \psi^{\prime \prime}(\rho(2 \delta))$.
Step 3. Solve the equation $\psi(t)=s$ to get $Q(s)$, the inverse function of $\psi(t), t \geq 1$. If the equation is hard to solve, derive lower and upper bound for $Q(s)$.
Step 4. Derive a lower bound for $\delta$ in term of $\Psi(\nu)$ by using $\delta(\nu) \geq \psi^{\prime}(Q(\Psi(\nu))) / 2$.
Step 5. Using the result of Steps 3 and 4 find positive constants $\kappa$ and $\gamma$, with $\gamma \in(0,1]$, such that $f(\tilde{\alpha}) \leq-\kappa \Psi(\nu)^{1-\gamma}$.
Step 6. Calculate the upper bound for $\Psi_{0}$ from

$$
\Psi_{0} \leq L_{\psi}(n, \theta, \tau)=n \psi\left(\frac{Q(\tau / n)}{\sqrt{1-\theta}}\right)
$$

Step 7. Derive an upper bound for the total number of iterations from

$$
\frac{\Psi_{0}^{\gamma}}{\kappa \gamma} \frac{1}{\theta} \log \frac{n}{\epsilon}
$$

Step 8. To calculate a complexity bound for large-update algorithms set $\tau=O(n)$ and $\theta=\Theta(1)$ and for small-update method algorithms set $\tau=O(1)$ and $\theta=\Theta(1 / \sqrt{n})$.

In the next section, we apply this scheme to obtain an iteration bound for the algorithm generated by the kernel function introduced in (1.4).

## 4. Iteration bounds of the algorithm

Since we are unable to get explicit expressions for the inverse functions $\rho$ and $Q$ in the above steps, we recall two lemmas from [2].

Lemma 4.1. Let $Q:[0, \infty] \rightarrow[1, \infty]$ be the inverse function of $\psi(t)$ for $t \geq 1$. Then we have

$$
\sqrt{1+2 s} \leq Q(s) \leq 1+\sqrt{2 s}, \quad s \geq 0
$$

LEMMA 4.2. Let $\psi_{b}(t)$ be the barrier term of $\psi(t)\left(\psi(t)=\left(t^{2}-1\right) / 2+\psi_{b}(t)\right)$ and let $\underline{\rho}:[0, \infty) \rightarrow(0,1]$ be the inverse function of the restriction of $-\psi_{b}^{\prime}(t)$ to the interval $(0,1]$. Then one has

$$
\rho(s) \geq \underline{\rho}(1+2 s) .
$$

Now by using these two lemmas, we derive some bounds for $\rho$ and $Q$.
Step 1. From the equation $-\psi_{b}^{\prime}(t)=s$ we have

$$
\begin{gathered}
-\psi_{b}^{\prime}(t)=\frac{e^{q(1 / t-1)}}{t^{2}}=s, \\
e^{q(1 / t-1)}=s t^{2} \Longleftrightarrow q\left(\frac{1}{t}-1\right)=\log s+2 \log t .
\end{gathered}
$$

From $0<t \leq 1$ we find that $\log t$ and as a result

$$
q\left(\frac{1}{t}-1\right) \leq \log s \quad \Longleftrightarrow \quad t=\underline{\rho}(s) \geq \frac{1}{1+q^{-1} \log s} .
$$

Now by using Lemma 4.2 , since $\rho(s) \geq \underline{\rho}(1+2 s)$, we derive a lower bound for $\rho(s)$ as follows:

$$
\begin{equation*}
\rho(s) \geq \frac{1}{1+q^{-1} \log (1+2 s)}, \quad s \geq 0 \tag{4.1}
\end{equation*}
$$

Step 2. The function $\psi^{\prime \prime}(t)$ is monotonically decreasing, hence

$$
f(\tilde{\alpha}) \leq-\frac{\delta^{2}}{\psi^{\prime \prime}(\rho(2 \delta))} \leq-\frac{\delta^{2}}{\psi^{\prime \prime}(\underline{\rho}(1+4 \delta))}
$$

Putting $t=\underline{\rho}(1+4 \delta)$, we have $t \leq 1$ and can write

$$
\begin{equation*}
f(\tilde{\alpha}) \leq-\frac{\delta^{2}}{\psi^{\prime \prime}(t)} \leq-\frac{\delta^{2}}{1+(q+2 t) t^{-4} e^{q(1 / t-1)}} \tag{4.2}
\end{equation*}
$$

Note that

$$
t=\underline{\rho}(1+4 \delta) \quad \Longleftrightarrow \quad(1+4 \delta)=-\psi_{b}^{\prime}(t)=\frac{e^{q(1 / t-1)}}{t^{2}}
$$

Substituting this equation in (4.2), we have

$$
f(\tilde{\alpha}) \leq-\frac{\delta^{2}}{1+(q+2 t) t^{-2}(1+4 \delta)} \leq-\frac{\delta^{2}}{1+3 q t^{-2}(1+4 \delta)}
$$

On the other hand,

$$
\frac{1}{t^{2}}=\frac{1}{\underline{\rho}(1+4 \delta)^{2}} \leq \frac{1}{\rho(2 \delta)^{2}} \leq\left(1+q^{-1} \log (1+4 \delta)\right)^{2}
$$

So

$$
\begin{equation*}
f(\tilde{\alpha}) \leq-\frac{\delta^{2}}{1+3 q(1+4 \delta)\left(1+q^{-1} \log (1+4 \delta)\right)^{2}} \tag{4.3}
\end{equation*}
$$

Step 3. By Lemma 4.1 the inverse function of $\psi(t)$ for $t \in[1, \infty)$ satisfies

$$
\begin{equation*}
\sqrt{1+2 \psi(t)} \leq Q(\psi(t)) \leq 1+\sqrt{2 \psi(t)} \tag{4.4}
\end{equation*}
$$

Step 4. Using $\delta(\nu) \geq \psi^{\prime}(Q(\Psi(\nu))) / 2$, we have

$$
\begin{equation*}
\delta \geq \frac{\psi^{\prime}(Q(\Psi(\nu)))}{2} \geq \frac{1}{2}\left(\sqrt{1+2 \Psi}-\frac{e^{q(1 / \sqrt{1+2 \Psi}-1)}}{(1+2 \Psi)^{2}}\right) . \tag{4.5}
\end{equation*}
$$

Note that $\Psi \geq 1$, so $(1+2 \Psi)^{3 / 2} \geq e^{q(1 / \sqrt{1+2 \Psi}-1)}$ and $\sqrt{1+2 \Psi} \leq \sqrt{3 \Psi}$. So we can write

$$
\begin{equation*}
\delta \geq \frac{1}{2}\left(\sqrt{1+2 \Psi}-\frac{1}{\sqrt{1+2 \Psi}}\right)=\frac{\Psi}{\sqrt{1+2 \Psi}} \geq \frac{\Psi}{\sqrt{3 \Psi}} \geq \sqrt{\frac{\Psi}{3}} \tag{4.6}
\end{equation*}
$$

Step 5. Let $\Psi_{0} \geq \Psi \geq \tau \geq 3$. We deduced that $\delta \geq 1$ and $\sqrt{\Psi} \leq \sqrt{3} \delta \leq 2 \delta$. Now by using (4.3) we have

$$
f(\tilde{\alpha}) \leq-\frac{\delta^{2}}{16 q \delta\left(1+q^{-1} \log (1+4 \delta)\right)^{2}} \leq-\frac{\delta}{16 q\left(1+q^{-1} \log (1+4 \delta)\right)^{2}}
$$

that is,

$$
\begin{equation*}
f(\tilde{\alpha}) \leq-\frac{\sqrt{\Psi}}{48 q\left(1+q^{-1} \log \left(1+\sqrt{\Psi_{0}}\right)\right)^{2}} \tag{4.7}
\end{equation*}
$$

Thus it follows that

$$
\begin{align*}
\Psi_{k+1} & \leq \Psi_{k}-\kappa\left(\Psi_{k}\right)^{1-\gamma}, \quad k=0,1, \ldots, K-1,  \tag{4.8}\\
\text { where } \quad \kappa & =\frac{1}{48 q\left(1+q^{-1} \log \left(1+\sqrt{\Psi_{0}}\right)\right)^{2}}, \quad \gamma=\frac{1}{2}
\end{align*}
$$

and $K$ denotes the number of inner iterations.

Step 6. From Lemma 4.1 we have $Q(\tau / n) \leq 1+\sqrt{2 \tau / n}$. As a consequence

$$
\Psi_{0} \leq L_{\psi}(n, \theta, \tau)=n \psi\left(\frac{Q(\tau / n)}{\sqrt{1-\theta}}\right) \leq n \psi\left(\frac{1+\sqrt{2 \tau / n}}{\sqrt{1-\theta}}\right)
$$

Since $\psi(t) \leq\left(t^{2}-1\right) / 2$ for $t \geq 1$,

$$
\begin{equation*}
\Psi_{0} \leq \frac{n^{2}}{2} \frac{\sqrt{2 \tau / n}+2 \tau / n}{1-\theta}=\frac{\tau+\sqrt{2 n \tau}}{1-\theta} \tag{4.9}
\end{equation*}
$$

Step 7. By inequality (4.8), the number of inner iterations is bounded by

$$
K \leq \frac{\Psi_{0}^{\gamma}}{\kappa \gamma}=96 q\left(1+q^{-1} \log \left(1+\sqrt{\Psi_{0}}\right)\right)^{2} \Psi_{0}^{1 / 2}
$$

Substituting (4.9) in this inequality gives

$$
K \leq 96 q\left(1+q^{-1} \log \left(1+\sqrt{\frac{\tau+\sqrt{2 n \tau}}{1-\theta}}\right)\right)^{2}\left(\frac{\tau+\sqrt{2 n \tau}}{1-\theta}\right)^{1 / 2} .
$$

Thus the total number of iterations is bounded above by

$$
\begin{equation*}
\frac{K}{\theta} \log \frac{n}{\epsilon} \leq 96 q\left(1+q^{-1} \log \left(1+\sqrt{\frac{\tau+\sqrt{2 n \tau}}{1-\theta}}\right)\right)^{2}\left(\frac{\tau+\sqrt{2 n \tau}}{1-\theta}\right)^{1 / 2} \frac{1}{\theta} \log \frac{n}{\epsilon} \tag{4.10}
\end{equation*}
$$

Step 8. For large-update methods set $\tau=O(n)$ and $\theta=\Theta(1)$. As a consequence,

$$
\Psi_{0} \leq \frac{\tau+\sqrt{2 n \tau}}{1-\theta}=O(n)
$$

By choosing

$$
q=\log \left(1+\sqrt{\frac{\tau+\sqrt{2 n \tau}}{1-\theta}}\right)=O(\log n)
$$

the total iteration bound in (4.10) becomes

$$
O(\sqrt{n}(\log n)) \log \frac{n}{\epsilon}
$$

Setting $\tau=O(1)$ and $\theta=\Theta(1 / \sqrt{n})$ for small-update methods, we can obtain the best known bound as follows.
By Lemma 2.4 in [2], we have $\psi(t)<\psi^{\prime \prime}(1)(t-1)^{2} / 2, t>1$. Since $\psi^{\prime \prime}(1)=q+3$, we achieve an upper bound for $\Psi_{0}$ of

$$
\begin{aligned}
\Psi_{0} & =\frac{n \psi^{\prime \prime}(1)}{2}\left(\frac{1+\sqrt{2 \tau / n}}{\sqrt{1-\theta}}-1\right)^{2} \leq \frac{n \psi^{\prime \prime}(1)}{2}\left(\frac{\theta+\sqrt{2 \tau / n}}{\sqrt{1-\theta}}\right)^{2} \\
& =\frac{q+3}{2} \frac{(\sqrt{2 \tau}+\theta \sqrt{n})^{2}}{1-\theta}
\end{aligned}
$$

where we also use that $1-\sqrt{1-\theta} \leq \theta$. Now $\Psi_{0}=O(q)$ and the iteration bound becomes

$$
O(q \sqrt{q n}) \log \frac{n}{\epsilon}
$$

By choosing the parameter $q$ as a constant which is independent of $n$, this is the best bound for small-update methods.

## 5. Concluding remarks

In this paper, we introduced a new kernel function and showed that it generates the best possible iteration bounds, both for small-update and for large-update methods. This paper was inspired as a result of the work on primal-dual interior-point methods (IPMs) for linear optimization based on kernel functions and the scheme for analyzing such methods. The new kernel function is a parameterized version of a kernel function introduced in [2].

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