# UNIQUE CONTINUATION AT INFINITY OF SOLUTIONS TO SCHRÖDINGER EQUATIONS WITH COMPLEX-VALUED POTENTIALS 

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(Received 13th March 1997)

Dedicated to the memory of Professor Olgierd A. Biberstein (1921-1997)


#### Abstract

We obtain optimal $L^{2}$-lower bounds for nonzero solutions to $-\Delta \Psi+V \Psi=E \Psi$ in $\mathbf{R}^{n}, n \geq 2, E \in \mathbf{R}$, where $V$ is a measurable complex-valued potential with $V(x)=O\left(|x|^{-c}\right)$ as $|x| \rightarrow \infty$, for some $\epsilon \in \mathbf{R}$. We show that if $3 \delta=\max \{0,1-2 \epsilon\}$ and $\exp \left(\tau|x|^{1+\delta}\right) \Psi \in L^{2}\left(R^{n}\right)$ for all $\tau>0$, then $\Psi$ has compact support. This result is new for $0<\epsilon<1 / 2$ and generalizes similar results obtained by Meshkov for $\epsilon=0$, and by Froese, Herbst, M. Hoffmann-Ostenhof, and T. Hoffmann-Ostenhof for both $\epsilon \leq 0$ and $\epsilon \geq 1 / 2$. These $L^{2}$-lower bounds are well known to be optimal for $\epsilon \geq 1 / 2$ while for $\epsilon<1 / 2$ this last is only known for $\epsilon=0$ in view of an example of Meshkov. We generalize Meshkov's example for $\epsilon<1 / 2$ and thus show that for complex-valued potentials our result is optimal for all $\epsilon \in \mathbf{R}$.


1991 Mathematics subject classification: 35J10, 35B40, 35B60, 81C05.

## 1. Introduction

Let $\epsilon \in \mathbf{R}$ be given and suppose $V$ is a measurable complex-valued function on $\mathbf{R}^{n}$ that satisfies

$$
\begin{equation*}
V(x)=O\left(|x|^{-t}\right), \quad \text { as } \quad|x| \rightarrow \infty \tag{1.1}
\end{equation*}
$$

In this paper we investigate the fastest possible rate of decay of the solutions to

$$
\begin{equation*}
-\Delta \Psi+V \Psi=E \Psi \tag{1.2}
\end{equation*}
$$

on $\mathbf{R}^{n}$, where $\Delta$ is the Laplacian on $\mathbf{R}^{n}, n \geq 2$, and $E \in \mathbf{R}$. Without further assumptions on $V$ we prove:

Theorem 1. Let $V, E$, and $\epsilon$ be as in (1.1) and (1.2) and let $3 \delta=\max \{0,1-2 \epsilon\}$. Let $\Psi \in H_{\mathrm{loc}}^{2}\left(\mathbf{R}^{n}\right)$ be a nonzero solution of (1.2) that satisfies

[^0]\[

$$
\begin{equation*}
\exp \left(\tau|x|^{1+\delta}\right) \Psi \in L^{2}\left(\mathbf{R}^{n}\right) \tag{1.3}
\end{equation*}
$$

\]

for all $\tau>0$. Then $\Psi$ has compact support.

We view this theorem as a unique continuation at infinity result and prove it through a Carleman-like estimate that generalizes work of Meshkov [10] for $\epsilon=0$. A similar result for both $\epsilon \leq 0$ and $\epsilon \geq 1 / 2$ has been obtained using different methods by Froese, Herbst, T. Hoffmann-Ostenhof, and M. Hoffmann-Ostenhof [8]. It is well known that Theorem 1 is optimal for $\epsilon \geq 1 / 2$ while for $\epsilon<1 / 2$ this last is only known for $\epsilon=0$ in view of an example due to Meshkov [10]. We generalize Meshkov's example for $\epsilon<1 / 2$ and prove:

Theorem 2. Let $\epsilon<1 / 2$ and $\delta>0$ satisfy $2 \epsilon+3 \delta=1$. Then there exist a continuous complex-valued function $V$ on $\mathbf{R}^{2}$ satisfying (1.1), and a $C^{2}$-function $\Psi$ which does not have compact support and satisfies $\Delta \Psi=V \Psi$ on $\mathbf{R}^{2}$ and

$$
\begin{equation*}
\Psi(x)=O\left(\exp \left(-\beta|x|^{1+\delta}\right)\right), \quad \text { as } \quad|x| \rightarrow \infty, \tag{1.4}
\end{equation*}
$$

for some $\beta>0$.

Thus, for complex-valued potentials Theorem 1 is optimal for all $\epsilon \in \mathbf{R}$.
The above results are closely related to the following question posed by B. Simon [12]. Let $V$ be a real-valued potential and suppose that $\Psi \not \equiv 0$ satisfies $(-\Delta+V-E) \Psi=0$ in $\mathbf{R}^{n}$. Is it true that if $-\Delta+V$ does not have compact resolvent, then $\exp (\tau|x|) \Psi \notin L^{2}\left(\mathbf{R}^{n}\right)$ for $\tau>0$ sufficiently large? The answer to this question is not known but a positive response is suggested by the sharp exponential upper and lower bounds of different kinds already established for several classes of potentials [1, 2, 3, 4, $5,6,7,8$ ]. The results for $0<\epsilon<1 / 2$ presented in this paper show that a proof of an affirmative answer to Simon's question has to use in an essential way the fact that $V$ is real-valued, even if $V(x)$ goes to zero as $|x|$ goes to infinity.

The general strategy to prove Theorems 1 and 2 is that of Meshkov [10]; however, the author also benefited from [11]. In Section 2 we obtain Carleman-like estimates near infinity and use these to prove Theorem 1. In Section 3 we generalize Meshkov's example.

## 2. Carleman-like estimates in exterior domains

Theorem 1 will be derived from the following Carleman-like estimate at infinity.

Lemma 2.1. Let $\rho>0, E \in \mathbf{R}$, and $\delta \geq 0$ be given. Set $\Omega_{\rho}=\left\{x \in \mathbf{R}^{n}:|x|>\rho\right\}$ and fix $\kappa$ such that $-3 \delta<\kappa \leq 2-\delta$. Then there exist a constant $K$ independent of $\tau$ and $\tau_{0}>0$ such that for any $v \in C_{0}^{\infty}\left(\Omega_{\rho}\right)$ and $\tau>\tau_{0}$ we have

$$
\begin{equation*}
\int_{\Omega_{0}}|x|^{3 \delta-n+\kappa}|v(x)|^{2} \exp \left(2 \tau|x|^{1+\delta}\right) d x \leq \frac{K}{\tau^{3}} \int_{\Omega_{\rho}}|x|^{x+1-n}|(\Delta+E) v(x)|^{2} \exp \left(2 \tau|x|^{1+\delta}\right) d x \tag{2.1}
\end{equation*}
$$

Remark. The constant $E$ in (2.1) is important only for $0 \leq \delta<1 / 3$.
Proof. Since $E \in \mathbf{R}$ we may assume without loss of generality that $v$ is real-valued. Set $\quad r=|x|, \omega=x /|x|$, and $\alpha=1+\delta$. For $\tau>0 \quad$ set $\quad u=\exp \left(\tau r^{\alpha}\right) v \quad$ and $\Delta_{\mathrm{t}}=\exp \left(\tau r^{\alpha}\right) \Delta \exp \left(-\tau r^{\alpha}\right)$, where $\exp \left(\tau r^{2}\right)$ and $\exp \left(-\tau r^{\alpha}\right)$ are multiplication operators. Using this notation we find that (2.1) is equivalent to

$$
\begin{equation*}
\int r^{3 \delta+\kappa-1}|u|^{2} d r d \omega \leq \frac{K}{\tau^{3}} \int r^{\kappa}\left|\left(\Delta_{\tau}+E\right) u\right|^{2} d r d \omega \tag{2.2}
\end{equation*}
$$

where $d \omega$ denotes the Lebesgue measure on $S^{n-1}$. Defining $L=2 \tau \alpha r^{\alpha-1}+\partial_{r}$ we have

$$
L+\Delta_{\tau}=\partial_{r r}+\tau^{2} \alpha^{2} r^{2 x-2}+\frac{n-1}{r} \partial_{r}-\tau \alpha(n+\alpha-2) r^{\alpha-2}+\frac{1}{r^{2}} \Lambda,
$$

where $\Lambda$ is the Laplace-Beltrami operator on the unit sphere $S^{n-1}$. Setting $u_{r}=\partial_{r} u$ and integrating by parts with respect to $r$ we obtain

$$
\begin{aligned}
\int r^{\kappa}\left|\left(\Delta_{\tau}+E\right) u\right|^{2} d r d \omega \geq & \int r^{\kappa}\left(|L u|^{2}-2 L u\left(L+\Delta_{\tau}+E\right) u\right) d r d \omega \\
= & \int\left(4 \tau^{2} \alpha^{2} r^{2 \alpha+\kappa-2}+2 \tau \alpha(\alpha+\kappa-2 n+1) r^{\alpha+\kappa-2}\right) u_{r}^{2} d r d \omega \\
& +2 \tau \alpha \int\left[\left(\tau^{2} \alpha^{2}(3 \alpha+\kappa-3) r^{3 \alpha+\kappa-4}\right.\right. \\
& \left.-\tau \alpha(2 \alpha-3+\kappa)(n+\alpha-2) r^{2 \alpha+\kappa-4}\right) u^{2} \\
& \left.+(\alpha+\kappa-3) r^{\alpha+\kappa-4} u \Lambda u-\alpha E(\alpha+\kappa-1) r^{\alpha+\kappa-2} u^{2}\right] d r d \omega \\
= & \int 2 \tau^{3} \alpha^{3}(3 \delta+\kappa) r^{3 \delta+\kappa-1} \\
& {\left[1-\frac{1}{\tau \alpha(3 \delta+\kappa)}\left(\frac{(2 \alpha+\kappa-3)(n+\alpha-2)}{r^{2}}+\frac{E(\delta+\kappa)}{\tau \alpha r^{2 \delta}}\right)\right] u^{2} d r d \omega } \\
& +\int 2 \tau \alpha\left[2 \tau \alpha r^{+2 \delta}\left(1+\frac{\alpha+\kappa-2 n+1}{2 \tau \alpha r^{2}}\right) u_{r}^{2}+(\delta+\kappa-2) r^{\alpha+\kappa-4} u \Lambda u\right] d r d \omega .
\end{aligned}
$$

Since $-3 \delta<\kappa \leq 2-\delta$ and $\Lambda$ is a negative operator on $L^{2}\left(S^{n-1}, d \omega\right)$ we have

$$
\int r^{\kappa}\left|\left(\Delta_{\tau}+E\right) u\right|^{2} d r d \omega \geq \int 2 \tau^{3} \alpha^{3}(3 \delta+\kappa) r^{3 \delta+\kappa-1} u^{2}\left(1+O\left(\frac{1}{\tau}\right)\right) d r d \omega
$$

from which (2.2) follows.

Proof of Theorem 1. Theorem 1 follows from (2.1) using standard arguments that we sketch here for the sake of completeness. Let $V, E, \epsilon, \delta$ and $\Psi$ be as in Theorem 1 . Assuming that

$$
\begin{equation*}
\int_{\mathbf{R}^{n}}|\Psi(x)|^{2} \exp \left(2 \tau|x|^{1+\delta}\right) d x<\infty \tag{2.3}
\end{equation*}
$$

for all $\tau>0$, we will prove that $\Psi$ has compact support. Using $L^{2}$-interior estimates [9], it follows from (2.3) that

$$
\begin{equation*}
\int_{\mathbf{R}^{n}}\left|D^{\beta} \Psi(x)\right|^{2} \exp \left(2 \tau|x|^{1+\delta}\right) d x<\infty \tag{2.4}
\end{equation*}
$$

for all $\tau>0$ and all multi-indices $\beta$ such that $|\beta| \leq 2$. Let $\kappa$ be as in Lemma 2.1 and fix $\rho \geq 0$ so that $\left.|V(x \mid) \leq C| x\right|^{-\epsilon}$ for $x \in \Omega_{\rho}$. A simple estimate, using (2.1) and $3 \delta=\max \{0,1-2 \epsilon\}$, shows that there exist a constant $K$ independent of $\tau$ and $\tau_{0}>0$ such that for any $v \in C_{0}^{\infty}\left(\Omega_{\rho}\right)$ and $\tau>\tau_{0}$ we have

$$
\begin{equation*}
\int_{\Omega_{\rho}}|x|^{3 \delta-n+\kappa}|v(x)|^{2} \exp \left(2 \tau|x|^{1+\delta}\right) d x \leq \frac{K}{\tau^{3}} \int_{\Omega_{\rho}}|x|^{\kappa+1-n}|(\Delta-V+E) v(x)|^{2} \exp \left(2 \tau|x|^{1+\delta}\right) d x . \tag{2.5}
\end{equation*}
$$

Let $h$ be a $C^{\infty}$-function on $\mathbf{R}^{n}$ which takes values between 0 and 1 , vanishes on $|x| \leq \rho+1 / 2$, and equals 1 on $|x| \geq \rho+1$. Let $\phi \in C_{0}^{\infty}\left(\mathbf{R}^{n}\right)$ be a function which equals 1 on $|x| \leq 1$ and set $\phi_{\lambda}(x)=\phi(x / \lambda)$ for $\lambda>0$. An approximation argument shows that (2.5) holds for every $v \in H^{2}\left(\mathbf{R}^{n}\right)$ with compact support contained in $\Omega_{\rho}$. Hence (2.5) holds on $\Omega_{\rho}$ for $v_{\lambda}=\phi_{\lambda} \Psi h$ and therefore, using (2.4), for $v=\Psi h$. Since $\Psi$ satisfies (1.2) and $h(x)=1$ on $\Omega_{\rho+1}$ we obtain

$$
\int_{\Omega_{\rho+1}}|x|^{3 \delta-n+\kappa}|\Psi(x)|^{2} d x \leq \frac{K}{\tau^{3}} \int_{\rho \leq|x| \leq \rho+1}|x|^{\kappa+1-n}|(\Delta-V+E) v(x)|^{2} d x
$$

Letting $\tau$ go to infinity in this last estimate we find that $\Psi \equiv 0$ on $\Omega_{\rho+1}$.

## 3. Examples

Although the essential idea of the construction given below is that of Meshkov [10], we present the details for the reader's convenience. For $\rho>0$ we will denote by $A(\alpha, \beta)$ the annulus in $\mathbf{R}^{2}$ defined by $\rho+\alpha \rho^{(1-\delta) / 2} \leq r \leq \rho+\beta \rho^{(1-\delta) / 2}$.

Lemma 3.1. Let $\delta>0$ and $\epsilon \in \mathbf{R}$ satisfy $2 \epsilon+3 \delta=1$. For a fixed and large $\rho>0$ let $n$ and $k$ be positive integers such that $\left|n-\rho^{1+\delta}\right| \leq 1$ and $\left|k-6(1+\delta) \rho^{(1+\delta) / 2}\right| \leq$ $1+20 \delta(1+\delta)$. Then there exist complex-valued functions $u$ and $V$ on $A(0,6)$ possessing the following properties:

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(a) The function $u$ is of class $C^{2}$ and satisfies

$$
\begin{equation*}
\Delta u+V u=0 \quad \text { on } A(0,6) \tag{3.1}
\end{equation*}
$$

(b) There exists a constant $C$ independent of $\rho, n$, and $k$ such that

$$
\begin{equation*}
|V(r, \theta)| \leq C r^{-c} \quad \text { on } A(0,6) \tag{3.2}
\end{equation*}
$$

(c) For a constant $a>0$ we have

$$
u(r, \theta)= \begin{cases}r^{-n} \exp (-i n \theta) & \text { on } A(0,0.1) \\ a r^{-n-k} \exp i(-n-k) \theta & \text { on } A(5.9,6)\end{cases}
$$

Therefore $V \equiv 0$ on the annuli $A(0,0.1)$ and $A(5.9,6)$.
(d) Let $m(r)=\max \{|u(r, \theta)|, 0 \leq \theta \leq 2 \pi\}$. Then

$$
\log m(r)-\log m(\rho) \leq \log 2-\frac{1}{6} \int_{\rho}^{r} t^{\delta} d t
$$

for $\rho \leq r \leq \rho+6 \rho^{(1-\delta) / 2}$.

Proof of Theorem 2. First we fix a large $\rho_{1}>0$ and set $\rho_{j+1}=\rho_{j}+6 \rho_{j}^{(1-\delta) / 2}$ for $j=1,2, \ldots$. Then we set $n_{j}=\left[\rho_{j}^{1+\delta}\right]$, where $[x]=\max \{n \in \mathbf{Z}: n \leq x\}$, and $k_{j}=n_{j+1}-n_{j}$. For $j=1,2, \ldots$ we have $n_{j}=\rho_{j}^{1+\delta}-\gamma_{j}$, with $0 \leq \gamma_{j}<1$, and

$$
\begin{aligned}
k_{j} & =\rho_{j+1}^{1+\delta}-\rho_{j}^{1+\delta}+\gamma_{j}-\gamma_{j+1} \\
& =\rho_{j}^{1+\delta}\left(1+6 \rho_{j}^{-(1+\delta) / 2}\right)^{1+\delta}-\rho_{j}^{1+\delta}+\gamma_{j}-\gamma_{j+1} \\
& =6(1+\delta) \rho_{j}^{(1+\delta) / 2}+18 \delta(1+\delta)+O\left(\rho_{j}^{-(1+\delta) / 2}\right)+\gamma_{j}-\gamma_{j+1} .
\end{aligned}
$$

Therefore if $\rho_{1}$ is large we may assume that $\left|k_{j}-6(1+\delta) \rho_{j}^{(1+\delta) / 2}\right| \leq 1+20 \delta(1+\delta)$. For $j=1,2, \ldots$, let $a_{j}$ be constants and $u_{j}$ and $V_{j}$ be functions constructed on $\rho_{j} \leq r \leq \rho_{j+1}$ as in Lemma 3.1. Then $u_{j}\left(\rho_{j}, \theta\right)=\rho_{j}^{-n_{j}} \exp \left(-i n_{j} \theta\right)$ and $u_{j}\left(\rho_{j+1}, \theta\right)=a_{j} \rho_{j+1}^{-n_{j+1}} \exp \left(-i n_{j+1} \theta\right)$. Since $\rho_{j} \rightarrow \infty$ as $j \rightarrow \infty$, then for $r>\rho_{1}$ we set $V(r, \theta)=V_{j}(r, \theta)$ and $\Psi(r, \theta)=A_{j} u_{j}(r, \theta)$ for $\rho_{j} \leq r \leq \rho_{j+1}$, where we define $A_{j}=a_{0} a_{1} \ldots a_{j-1}$, for $j=1,2, \ldots$, and $a_{0}=1$. Clearly $V$ satisfies (1.2) and $\Psi$ is of class $C^{2}$ and satisfies $-\Delta \Psi+V \Psi=0$ on $\Omega_{\rho_{1}}$. To prove that $\Psi$ satisfies (1.4) we set $\mu(r)=\max \{|\Psi(r, \theta)|: 0 \leq \theta \leq 2 \pi\}$ for $r>\rho_{1}$ and pick $l$ such that $\rho_{l} \leq r \leq \rho_{l+1}$. Then

$$
\log \mu(r)=\left(\log m_{l}(r)-\log m_{l}\left(\rho_{l}\right)+\ldots+\left(\log m_{1}\left(\rho_{2}\right)-\log m_{1}\left(\rho_{1}\right)\right)+\log m_{1}\left(\rho_{1}\right)\right.
$$

where $m_{j}(r)$ is as in (d) of Lemma 3.1. Using (d) of this lemma we find that

$$
\log \mu(r) \leq l \log 2-\frac{1}{6} \int_{\rho_{1}}^{r} t^{\delta} d t+\log m\left(\rho_{1}\right)
$$

Thus if $\delta>1$ then $\log \mu(r) \leq C r^{(1+\delta / 2}-c r^{1+\delta}$, and if $0<\delta \leq 1$ then $\log \mu(r) \leq C r-c r^{1+\delta}$, where $C$ and $c$ are positive constants. Therefore, since $\delta>0$, for $r$ sufficiently large we have

$$
0<\mu(r) \leq C \exp \left(-\beta r^{1+\delta}\right)
$$

for some $\beta>0$. The functions $V$ and $\Psi$ defined above can easily be extended to $\mathbf{R}^{2}$ in a way that Theorem 2 is satisfied.

Proof of Lemma 3.1. We will smoothly modify in four steps the function $u_{1}=r^{-n} \exp (-i n \theta)$ into a function $u$ on $A(0,6)$ that satisfies (a), (b), (c), and (d). In this proof $C$ denotes a positive constant independent of $\rho, k$, and $n$.
I. The annulus $A(0,2)$. For $m=0,1, \ldots, 2 n+2 k-1$ we set $\theta_{m}=m T$, where $T \equiv \pi /(n+k)$. Let $f$ be a smooth $T$-periodic function on $\mathbf{R}$ such that $\int_{0}^{T} f(\theta) d \theta=0$, $f(\theta)=-4 k$ on $[0, T / 5] \cup[4 T / 5, T]$, and $-4 k \leq f(\theta) \leq 5 k$ and $\left|f^{\prime}(\theta)\right| \leq C k / T$, for $0 \leq \theta \leq T$. Set

$$
\Phi(\theta)=\int_{0}^{0} f(t) d t
$$

Clearly $\Phi$ is $T$ - and $2 \pi$-periodic, and $\Phi\left(\theta_{m}\right)=0$. In addition, for $\theta \in \mathbf{R}$ we have

$$
\begin{equation*}
|\Phi(\theta)| \leq 5 k /(n+k), \quad\left|\Phi^{\prime}(\theta)\right| \leq 5 k, \quad \text { and } \quad\left|\Phi^{\prime \prime}(\theta)\right| \leq C k n, \tag{3.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\Phi(\theta)=-4 k\left(\theta-\theta_{m}\right) \equiv-4 k \theta+b_{m}, \quad \text { for }\left|\theta-\theta_{m}\right| \leq T / 5 . \tag{3.4}
\end{equation*}
$$

Set $F(\theta)=(n+2 k) \theta+\Phi(\theta), b=\left(\rho+\rho^{(1-\delta) / 2}\right)^{-2 k}$, and $u_{2}=-b r^{-n+2 k} \exp (i F(\theta))$. Note that $\left|u_{1}(r, \theta)\right|=\left|u_{2}(r, \theta)\right|$ for $r=\rho+\rho^{(1-\theta) / 2}$; in addition, it follows from (3.4) that $u_{2}=-b r^{-(n-2 k)} \exp \left(i(n-2 k) \theta+i b_{m}\right)$ on the sectors

$$
S_{m} \equiv\left\{(r, \theta):\left|\theta-\theta_{m}\right| \leq T / 5\right\}, \quad m=0,1, \ldots, 2 n+2 k-1
$$

On $A(0,1 / 3)$ we have

$$
\begin{aligned}
\frac{\left|u_{2}(r, \theta)\right|}{\left|u_{1}(r, \theta)\right|} & =\frac{r^{2 k}}{\left(\rho+\rho^{(1-\delta) / 2}\right)^{2 k}} \\
& \leq\left(1+\frac{2}{3} \frac{1}{\rho^{(1+\delta) / 2}+\frac{1}{3}}\right)^{-2 k}
\end{aligned}
$$

Hence using the assumptions on $k$ and $\rho$ we obtain

$$
\begin{equation*}
\left|u_{2}(r, \theta)\right| \leq \exp (-8)\left|u_{1}(r, \theta)\right| \quad \text { on } A(0,1 / 3) \tag{3.5}
\end{equation*}
$$

Similarly

$$
\begin{equation*}
\left|u_{2}(r, \theta)\right| \geq \exp (8)\left|u_{1}(r, \theta)\right| \quad \text { on } A(5 / 3,2) . \tag{3.6}
\end{equation*}
$$

Let $\psi_{i}(r), i=1,2$, be $C^{\infty}$-functions taking values between 0 and 1 such that $\psi_{1}$ vanishes for $r \geq \rho+1.9 \rho^{(1-\delta) / 2}$ and equals 1 for $r \leq \rho+(5 / 3) \rho^{(1-\delta) / 2}, \psi_{2}$ vanishes for $r \leq \rho+0.1 \rho^{(1-\delta) / 2}$ and equals 1 for $r \geq \rho+(1 / 3) \rho^{(1-\delta) / 2}$, and

$$
\begin{equation*}
\left|\psi_{i}^{(p)}(r)\right| \leq C \rho^{-p(1-\delta) / 2}, \quad r \geq 0 ; \quad i=1,2 ; p=1,2 . \tag{3.7}
\end{equation*}
$$

Define $u=\psi_{1} u_{1}+\psi_{2} u_{2}$. Clearly $u$ is harmonic in $S \equiv A(1 / 3,5 / 3) \cap\left(U S_{m}\right)$. Now set

$$
V(r, \theta)= \begin{cases}0 & (r, \theta) \in S \\ \Delta u / u & \text { otherwise }\end{cases}
$$

Clearly (3.1) holds on $A(0,2)$. Next we show that $|u|>0$ on $A(0,2) \backslash S$ and that (3.2) holds on $A(0,2)$.

On $A(0,1 / 3)$ we have

$$
\begin{equation*}
\Delta u=\psi_{2} \Delta u_{2}+2 \psi_{2}^{\prime} \partial_{r} u_{2}+\left(\psi_{2}^{\prime} / r+\psi_{2}^{\prime \prime}\right) u_{2} \tag{3.8}
\end{equation*}
$$

and, using (3.5),

$$
\begin{equation*}
|u| \geq\left|u_{1}\right|-\left|u_{2}\right| \geq \frac{1}{2}\left|u_{1}\right| \geq \exp (7)\left|u_{2}\right|>0 . \tag{3.9}
\end{equation*}
$$

Similarly on $A(5 / 3,2)$ we have

$$
\begin{equation*}
\Delta u=\Delta u_{2}+2 \psi_{1}^{\prime} \partial_{r} u_{1}+\left(\psi_{1}^{\prime} / r+\psi_{1}^{\prime \prime}\right) u_{1} \tag{3.10}
\end{equation*}
$$

and, using (3.6),

$$
\begin{equation*}
|u| \geq\left|u_{2}\right|-\left|u_{1}\right| \geq \exp (7)\left|u_{1}\right|>0 . \tag{3.11}
\end{equation*}
$$

A short calculation shows that on $A(0,2)$ we have

$$
\Delta u_{2}=\left[\frac{\left(4 k+2 n+\Phi^{\prime}\right) \Phi^{\prime}-8 k n}{r^{2}}+\frac{i \Phi^{\prime \prime}}{r^{2}}\right] u_{2}
$$

Using (3.3) we obtain

$$
\left|\Delta u_{2}\right| \leq \frac{C k n}{r^{2}}\left|u_{2}\right| .
$$

Thus, by the assumptions on $k, n, \epsilon$, and $\delta$, we have

$$
\begin{equation*}
\left|\Delta u_{2}\right| \leq C r^{-c}\left|u_{2}\right| \quad \text { on } A(0,2) . \tag{3.12}
\end{equation*}
$$

We also have

$$
\left|\psi_{2}^{\prime} \partial_{r} u_{2}\right|=\left|\psi_{2}^{\prime} \frac{n-2 k}{r} u_{2}\right| \leq C \frac{n}{r} \rho^{-(1-\delta) / 2}\left|u_{2}\right| \leq C r^{-c}\left|u_{2}\right|
$$

and

$$
\left|\left(\psi_{2}^{\prime} / r+\psi_{2}^{\prime \prime}\right) u_{2}\right| \leq C\left(\rho^{-(1-\delta) / 2} / r+\rho^{-(1-\delta)}\right)\left|u_{2}\right| \leq C r^{-\epsilon}\left|u_{2}\right| .
$$

Combining these three last estimates, (3.8), and (3.9) we find that (3.2) holds on $A(0,1 / 3)$. Analogously, using (3.10), (3.11), and (3.12), we have that (3.2) holds on $A(5 / 3,2)$. It remains to show that $|u|>0$ and that (3.2) holds on

$$
P_{m}=\left\{(r, \theta): \theta_{m}+\frac{T}{5} \leq \theta \leq \theta_{m}+\frac{4 T}{5}\right\} \cap A(1 / 3,5 / 3), \quad m=0, \ldots, 2 n+2 k-1 .
$$

For this purpose we set $G(\theta)=F(\theta)+n \theta$. On the annular sectors $P_{m}$ we have

$$
\begin{equation*}
|u|=\left|u_{1}+u_{2}\right|=\left|u_{2}\right|\left|\exp (i G(\theta))-\frac{1}{b r^{2 k}}\right| \tag{3.13}
\end{equation*}
$$

We will show now that for some $\eta>0$ we have

$$
\begin{equation*}
\left|\exp (i G(\theta))-\frac{1}{b r^{2 k}}\right| \geq \eta, \quad(r, \theta) \in P_{m}, \quad m=0, \ldots, 2 n+2 k-1 \tag{3.14}
\end{equation*}
$$

Using this last, (3.12), and the fact that $\Delta u=\Delta u_{2}$ on $P_{m}$ we obtain $|u|>\eta\left|u_{2}\right|$ and therefore (3.2) holds on $P_{m}$. To prove (3.14) note that $G(\theta)=2(n+k) \theta+\Phi(\theta)$ and $G^{\prime}(\theta)=2(n+k)+f(\theta)$. Hence by the assumptions of $f, k$, and $n$ we may assume that $G^{\prime}(\theta)>n>0$. Since $G\left(\theta_{m}\right)=2 \pi m$ and $G\left(\theta_{m+1}\right)=2 \pi(m+1)$ we conclude that

$$
2 \pi m+\frac{n T}{5} \leq G(\theta) \leq 2 \pi(m+1)-\frac{n T}{5} \quad \text { for } \quad \theta_{m}+\frac{T}{5} \leq \theta \leq \theta_{m}+\frac{4 T}{5}
$$

Using the definition of $T$ and the assumptions on $k$ and $n$ we find that

$$
2 \pi m+\frac{\pi}{7} \leq G(\theta) \leq 2 \pi(m+1)-\frac{\pi}{7} \quad \text { for } \quad \theta_{m}+\frac{T}{5} \leq \theta \leq \theta_{m}+\frac{4 T}{5}
$$

It follows from this last estimate and (3.13) that (3.14) holds with $\eta=\sin (\pi / 7)$.
II. On $A(2,3)$ we deform $u_{2}$ into $u_{3}=-b r^{-n+2 k} \exp i(-n+2 k) \theta$. Let $\psi(r)$ be a $C^{\infty}$ function which takes values between 0 and 1 , equals 1 for $r \leq \rho+(7 / 3) \rho^{(1-\delta) / 2}$, vanishes for $r \geq \rho+(8 / 3) \rho^{(1-\delta) / 2}$, and satisfies (3.7). On $A(2,3)$ we set $u=-b r^{-n+2 k} \exp i(\psi(r) \Phi(\theta)$ $+(n+2 k) \theta)$ and $V=\Delta u / u$. A short calculation shows that

$$
\Delta u=\left[\left(\frac{-n+2 k}{r}+i \psi^{\prime} \Phi\right)^{2}+i \Phi\left(\frac{\psi^{\prime}}{r}+\psi^{\prime \prime}\right)-\frac{\left(n+2 k+\psi \Phi^{\prime}\right)^{2}}{r^{2}}+\frac{i \psi \Phi^{\prime \prime}}{r^{2}}\right] u .
$$

Using (3.3), the assumptions on $\psi^{(p)}$, and the assumptions on $k, n, \epsilon$, and $\delta$ we find that $\psi^{\prime} \Phi=O(1 / r)$, that $\psi \Phi^{\prime}=O(k)$, that $\psi \Phi^{\prime \prime}=O(k n)$, and that $\Phi\left(\psi^{\prime} / r+\psi^{\prime \prime}\right)=O\left(r^{-c}\right)$. Hence

$$
\Delta u=\left[\frac{O(n k)}{r^{2}}+O\left(r^{-\epsilon}\right)\right] u .
$$

Using again the assumptions on $k, n, \epsilon$, and $\delta$ we find that (3.2) holds on $A(2,3)$.
III. On $A(3,4)$ we deform $u_{3}$ into $u_{4}=-b r^{-(n+2 k)} \exp i(n+2 k) \theta$, where $b$ is as in I and $d \equiv\left(\rho+3 \rho^{(1-\delta) / 2}\right)^{4 k}$. Let $\psi$ be a $C^{\infty}$-function which takes values between 0 and 1 , equals 1 for $r \leq \rho+(10 / 3) \rho^{(1-\delta) / 2}$, vanishes for $r \geq \rho+(11 / 3) \rho^{(1-\delta) / 2}$, and satisfies (3.7). Next we define $h(r)=\psi(r)+(1-\psi(r)) d r^{-4 k}$. It is easily verified using the assumptions on $\psi, k$, and $\delta$ that $h$ satisfies (3.7) and that

$$
h(r) \geq d r^{-4 k} \geq\left(1+\frac{1}{\rho^{(1+\delta) / 2}+1}\right)^{-4 k} \geq \exp (-25(1+\delta))
$$

Now we set $u=u_{3} h$ and $V=\Delta u / u$, and verify as above that (3.2) holds on $A(3,4)$. In addition, on $A(11 / 2,4)$ we have $u=-b d r^{-(n+2 k)} \exp i(n+2 k) \theta$.
IV. Finally on $A(4,6)$ we deform $u_{4}$ into $u_{5}=a r^{-n-k} \exp i(-n-k) \theta$, where $a \equiv b d\left(\rho+5 \rho^{(1-\delta) / 2}\right)^{-k}$ and $b$ and $d$ are as in III. Note that $a$ has been chosen so that $\left|u_{4}(r, \theta)\right|=\left|u_{5}(r, \theta)\right|$ for $r=\rho+5 \rho^{(1-\delta) / 2}$. Let $\psi_{i}(r), i=1,2$, be $C^{\infty}$-functions taking values between 0 and 1 and satisfying (3.7), such that $\psi_{1}$ vanishes for $r \geq \rho+5.9 \rho^{(1-\delta / 2 / 2}$ and equals 1 for $r \leq \rho+(17 / 3) \rho^{(1-\delta / 2}$, and $\psi_{2}$ vanishes for $r \leq \rho+4.1 \rho^{(1-\delta) / 2}$ and equals 1 for $r \geq \rho+(13 / 3) \rho^{(1-\delta) / 2}$. Now on $A(4,6)$ we set $u=\psi_{1} u_{4}+\psi_{2} u_{5}$. It is clear that $u$ is harmonic on $A(13 / 3,17 / 3)$. Therefore we set $V=0$ on this annulus. We verify as in I that $V=\Delta u / u$ satisfies (3.2) on the remaining points of $A(4,6)$.

To finish this proof we set $m(r)=\max \{|u(r, \theta)|, 0 \leq \theta \leq 2 p\}$ and

$$
M(r)= \begin{cases}r^{-n} & \rho \leq r \leq \rho+\rho^{(1-\delta) / 2} \\ b r^{-n+2 k} & \rho+\rho^{(1-\delta) / 2} \leq r \leq \rho+3 \rho^{(1-\delta) / 2} \\ b r^{-n+2 k} h(r) & \rho+3 \rho^{(1-\delta) / 2} \leq r \leq \rho+4 \rho^{(1-\delta) / 2} \\ b d r^{-n-2 k} & \rho+4 \rho^{(1-\delta) / 2} \leq r \leq \rho+5 \rho^{(1-\delta) / 2} \\ a r^{-n-k} & \rho+5 \rho^{(1-\delta) / 2} \leq r \leq \rho+6 \rho^{(1-\delta) / 2}\end{cases}
$$

where $a, b$, and $d$ are as in IV. It is clear that $M(r)$ is a continuous piecewise smooth function on $\left[\rho, \rho+6 \rho^{(1-\delta) / 2}\right]$ that satisfies $m(r) \leq 2 M(r), m(\rho)=M(\rho)$, and

$$
\frac{d}{d r} \log M(r)=\frac{-n+O(k)}{r} \leq-\frac{\rho^{1+\delta}}{2 r} \leq-\frac{1}{6}\left(\rho+6 \rho^{(1-\delta) / 2}\right)^{\delta} \leq-\frac{1}{6} r^{\delta} .
$$

## Therefore

$$
\log m(r)-\log m(\rho) \leq \log 2+\int_{\rho} \frac{d}{d r} \log M(t) d t \leq \log 2-\frac{1}{6} \int_{\delta}^{r} d t
$$

which proves Lemma 3.1.
Acknowledgements. I would like to thank Professor I. Herbst for suggesting this problem as well as Meshkov's paper to me, and for very helpful conversations. I would also like to thank Professor L. Hörmander for useful comments.

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[^0]:    * Supported in part by CONACyT, Mexico.

