# A CHARACTERIZATION AND A CLASS OF DISTRIBUTION FUNCTIONS FOR THE STIELTJES-WIGERT POLYNOMIALS 

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1. In his classic memoir on the moment problem that bears his name, Stieltjes [2] exhibited

$$
\begin{gather*}
q^{-(n+1)^{2 / 2}}=\pi^{-1 / 2} \int_{0}^{\infty} x^{n}[1+A \sin (2 \pi \log x)] \exp \left(-\log ^{2} x\right) d x \\
\left(q=\exp \left(-\frac{1}{2}\right), \quad|A| \leq 1, \quad n=0,1,2, \ldots\right) \tag{1}
\end{gather*}
$$

as an example of an indeterminate (Stieltjes) moment sequence.
Stieltjes also obtained the corresponding $S$-fraction and thus implicitly obtained the three-term recurrence formula satisfied by the corresponding orthogonal polynomials.

Later, Wigert [4] found an explicit formula for the orthonormal polynomials corresponding to the more general weight function

$$
\begin{gather*}
w(x)=\gamma \pi^{-1 / 2} \exp \left(-\gamma^{2} \log ^{2} x\right), \quad x>0 \\
\left(\gamma^{2}=-(2 \log q)^{-1}, \quad 0<q<1\right) . \tag{2}
\end{gather*}
$$

Following Szegö [3], we call these polynomials the Stieltjes-Wigert polynomials and denote the orthonormal polynomial of degree $n$ by $S_{n}(x)$.

Now the weight function (2) satisfies the identity

$$
\begin{equation*}
x w(x)=q^{-1 / 2} w(q x), \quad x>0 . \tag{3}
\end{equation*}
$$

This shows that $\left\{S_{n}(q x)\right\}$ is an orthogonal polynomial sequence (OPS) with respect to the weight function $x w(x)$ on $(0, \infty)$.

This raises the question of whether there are other OPS with this property but it is quickly seen that the answer is negative. More precisely, suppose $\left\{P_{n}(x)\right\}$ is an OPS with respect to a distribution, $d \psi(x)$, on $[0, \infty)$. Suppose also that $\left\{P_{n}(a x)\right\}$ is an OPS with respect to $x d \psi(x)$. Writing

$$
\begin{equation*}
\mu_{n}=\int_{0}^{\infty} x^{n} d \psi(x) \quad(n=0,1,2, \ldots) \tag{4}
\end{equation*}
$$

it follows that $\left\{a^{-n} \mu_{n}\right\}_{0}^{\infty}$ is a moment sequence corresponding to $\left\{P_{n}(a x)\right\}$. But since $\left\{P_{n}(a x)\right\}$ is also an OPS with respect to $x d \psi(x)$, then $\left\{P_{n}(a x)\right\}$ is orthogonal with respect to $\left\{\mu_{n+1}\right\}_{0}^{\infty}$. Thus there exists $c>0$ such that

$$
\mu_{n+1}=c a^{-n} \mu_{n} \quad(n \geq 0)
$$

hence

$$
\mu_{n}=c^{n} a^{-n(n-1) / 2} \mu_{0}=\left(c a^{3 / 2}\right)^{n} a^{-(n+1)^{2 / 2}}\left(a^{1 / 2} \mu_{0}\right) .
$$

Now the condition

$$
\Delta_{2} \equiv \mu_{2} \mu_{0}-\mu_{1}^{2}>0
$$

is necessary for (4) to be a Stieltjes moment sequence and this requires $0<a<1$. Hence imposing the latter condition and writing $a=q$, we see upon referring to (2), which has the moments, $q^{-(n+1)^{2} / 2}(n \geq 0)$, that apart from a factor independent of $x$,

$$
P_{n}\left(c q^{3 / 2} x\right)=S_{n}(x) .
$$

2. Although the generalization first suggested by (3) does not lead to any new OPS, it does lead to a rather interesting class of new distribution functions having the same moments as (2).

Consider the generalization of (3) to general distribution functions:

$$
\begin{align*}
\int_{0}^{x} t d \psi(t) & =c \int_{0}^{x} d \psi(q t), \quad x>0 \\
\psi(t) & =0 \quad \text { for } x \leq 0 \tag{5}
\end{align*}
$$

Now if (5) has finite moments $\mu_{n}$, then the preceding discussion still applies to yield

$$
\begin{equation*}
\mu_{n}=q^{-(n+1)^{2 / 2}} \tag{6}
\end{equation*}
$$

provided

$$
\psi(\infty) \equiv \mu_{0}=q^{-1 / 2}, \quad c=q^{-3 / 2} .
$$

Moreover, if $\psi$ satisfies (5), then $\psi$ is continuous at the origin and has jump $J$ at $\tau$ if and only if it also has a jump $c^{-1} \tau J$ at $q \tau$. That is, if $\mathscr{D}$ denotes the set of discontinuities of a nondecreasing solution of (5), then

$$
\tau \in \mathscr{D} \Leftrightarrow \tau>0 \quad \text { and } \quad q^{k} \tau \in \mathscr{D} \quad(k=0, \pm 1, \pm 2, \ldots) .
$$

Since $\mathscr{D}$ is at most denumerable, it follows that there is an at most denumerable set $T$ such that

$$
\begin{equation*}
\mathscr{D}=\left\{q^{k} \tau \mid \tau \in T, \quad k=0, \pm 1, \pm 2, \ldots\right\} . \tag{7}
\end{equation*}
$$

Assuming $\mathscr{D} \neq \phi$, write

$$
T=\left\{\tau_{1}, \tau_{2}, \ldots\right\}=\left\{\tau_{v}\right\}_{v=1}^{N}, \quad 1 \leq N \leq \infty .
$$

If $j_{v}$ denotes the jump of $\psi$ at $\tau_{v}$, then the jump at $q^{k} \tau_{v}$ is given by

$$
\begin{equation*}
J_{v k}=c^{-k} \tau_{v}^{k} q^{k(k-1) / 2} j_{v}, \quad(k=0, \pm 1, \pm 2, \ldots ; v=1,2, \ldots, N) . \tag{8}
\end{equation*}
$$

Conversely, consider now an arbitrary (finite or infinite) sequence $\left\{\tau_{v}\right\}_{v=1}^{N}$ of distinct positive numbers. Let

$$
f(x) \equiv f(q ; x)=\sum_{k=-\infty}^{\infty} q^{k^{2} / 2} x^{k}, \quad 0<q<1, \quad x \neq 0
$$

and let $\left\{j_{v}\right\}_{v=1}^{N}$ be any sequence of positive numbers such that

$$
\sum_{v=1}^{N} f\left(\tau_{v}\right) j_{v}<\infty
$$

It is easily verified that

$$
f(q x)=q^{-1 / 2} x^{-1} f(x)
$$

hence

$$
\begin{equation*}
f\left(q^{n+1} \tau_{v}\right)=\left[q^{(n+1)^{2} / 2} \tau_{v}^{n}\right]^{-1} f\left(\tau_{v}\right) . \tag{9}
\end{equation*}
$$

Now define

$$
\begin{equation*}
\phi(x)=\sum_{q^{k} \tau_{v} \leq x} J_{v k} \quad(\nu=1,2,3, \ldots, N ; \quad k=0, \pm 1, \pm 2, \ldots) \tag{10}
\end{equation*}
$$

where $\phi(x)=0$ for $x \leq 0$ and $J_{v k}$ is given by (8) with $c=q^{-3 / 2}$. Then

$$
\begin{aligned}
\int_{0}^{\infty} x^{n} d \phi(x) & =\sum_{v=1}^{\infty} \sum_{k=-\infty}^{\infty}\left(q^{k} \tau_{v}\right)^{n} c^{-n} \tau_{v}^{k} q^{k(k-1) / 2} j_{v} \\
& =\sum_{\nu=1}^{N} f\left(\tau_{v} q^{n+1}\right) \tau_{v}^{n} j_{v}
\end{aligned}
$$

According to (9)

$$
\int_{0}^{\infty} x^{n} d \phi(x)=q^{-(n+1)^{2} / 2} \sum_{v=1}^{N} f\left(\tau_{v}\right) j_{v}
$$

Thus the jump function $\psi$ defined by

$$
\begin{equation*}
\psi(x)=K^{-1} \phi(x), \quad K=\sum_{v=1}^{N} f\left(\tau_{v}\right) j_{v}, \tag{11}
\end{equation*}
$$

has the moments (6).
We thus obtain a continuum of solutions to this Stieltjes moment problem, each of which is a jump function whose points of discontinuity form a set of the type (7) (and whose spectrum is thus the closure of $\mathscr{D}$ ). In particular, we can construct explicit solutions of the moment problem whose points of discontinuity form an everywhere dense subset of $[0, \infty)$. We remark that none of the above solutions are extremal solutions so by a theorem of M. Riesz [1], $\left\{S_{n}(x)\right\}$ is not complete in the corresponding $L^{2}$ spaces.
3. As a consequence of the above, we note a minor fact concerning the zeros of the Stieltjes-Wigert polynomials (about which very little seems to be known).

It is well known that on any open interval in which a distribution function is constant, the corresponding orthogonal polynomials each have at most one zero (see Szegö [3, Theorem 3.41.2]). If we take $T=\{\tau\}(\tau>0)$ and construct the corresponding distribution function, it follows that $S_{n}(x)$ has at most one zero in every open interval of the form

$$
\left(q^{k+1} \tau, q^{k} \tau\right) \quad(k=0, \pm 1, \pm 2, \ldots)
$$

for every $\tau>0$. (By a slightly more refined analysis, it is possible to extend this conclusion to the corresponding half-open intervals.)

## References

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2. T. J. Stieltjes, Recherches sur les fractions continues, Oeuvres, Tome II, Noordhoff, Groningen, 1918, 402-566.
3. G. Szegö: Orthogonal Polynomials, Colloq. Publ., Vol. 23, Amer. Math. Soc., New York, 1939.
4. S. Wigert, Sur les polynomes orthogonaux et l'approximation des fonctions continues. Ark. Mat. Astr. Fys. (18) 17 (1923), 15 pp.

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