

# A STUDY OF GAMMA SPECTRAL BREAK IN AGN

A. MARCOWITH, G. HENRI and G. PELLETIER

*Laboratoire d'Astrophysique  
Observatoire de Grenoble BP59X  
F38 041 GRENOBLE Cedex FRANCE*

Since its launch, CGRO has detected more than 20  $\gamma$ -ray emitting AGN, most of them associated with powerful, radio-loud, flat-spectrum objects, exhibiting VLBI superluminal motions. In the case of 3C279, the huge value of the apparent luminosity ( $\sim 10^{48} \text{erg.s}^{-1}$ ) and the variability time-scale of a few days (Hartmann *et al.*, 1992) gives a very large compactness  $\ell_{app} \simeq 200$ , that is, the medium should be completely thick to  $\gamma$ -rays. This contradiction can be explained if the  $\gamma$ -rays originate from a relativistic jet pointing at a small angle with respect to the line of sight (Maraschi *et al.*, 1992). However, the still large value of compactness suggests the existence of an inner, more compact region where pair production can take place efficiently (Henri *et al.*, 1993). This supports the so-called "two-flow" model, where the superluminal motion is attributed to the expansion of a relativistic pair plasma heated by a MHD jet from an accretion disk (Sol *et al.*, 1989). Hence we propose to interpret the spectral break observed in many objects around a few MeV (Lichti *et al.*, 1993) by an opacity effect due to *photon-photon absorption by pair production*.

Assuming that photons of energy  $\varepsilon_1 m_e c^2$  interact only with X photons of energy  $m_e c^2 / \varepsilon_1$ , and for a cylindrical jet of transverse radius  $r(z)$ , the opacity to pair production is  $\tau_{\gamma\gamma}(z) \simeq \sigma_T \dot{n}(1/\varepsilon_1) r^2(z) / c \varepsilon_1$  where  $\dot{n}(1/\varepsilon_1)$  is the density of photons produced per second by Inverse Compton process at energy  $1/\varepsilon_1$ . The total luminosity can then be evaluated as  $L_\gamma \simeq m_e c^2 \int_{z_\gamma}^\infty \varepsilon_1 \dot{n}(\varepsilon_1, z) S(z) dz$ , where  $z_\gamma$  is defined by  $\tau_{\gamma\gamma}(z_\gamma) = 1$ . The jet can become optically thin to  $\gamma$ -rays only by a decrease of soft photons and/or pair flux along  $z$ . Then one obtains  $L_\gamma = L_{thin} \varepsilon_1^{-\Delta\alpha}$ , where  $\Delta\alpha = \frac{s-1}{2}$  for a power-law distribution of electron with an index  $s$  and an exponential decrease of soft photons. The break energy is equal to 511 keV blueshifted by the Doppler factor. Pair annihilation photons can produce a bump at this break energy. A crucial prediction of the model is that variability at high energy should lag the low energy one, in contrast with optically thin models.

## References

- Hartmann, R. C., *et al.*: 1992, *ApJ Lett.*, **385**, L1  
Lichti, G. G., *et al.*: 1993, *these proceedings*  
Henri, G., Pelletier, G., and Roland, J.: 1993, *ApJ Lett.*, **404**, L41  
Maraschi, L., Ghisellini, G., and Celotti, A.: 1992, *ApJ*, **397**, L5  
Sol, H., Pelletier, G., and Asseo, E.: 1989, *MNRAS*, **237**, 411