

Molecular evolution in star-forming cores: From prestellar cores to protostellar cores

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Abstract. We investigate the molecular abundances in protostellar cores by solving the gas-grain chemical reaction network. As a physical model of the core, we adopt a result of one-dimensional radiation-hydrodynamics calculation, which follows the contraction of an initially hydrostatic prestellar core to form a protostellar core. Temporal variation of molecular abundances is solved in multiple infalling shells, which enable us to investigate the spatial distribution of molecules in the evolving core. The shells pass through the warm region of $T \sim 20\text{--}100$ K in several 10^4 yr and falls onto the central star in ~ 100 yr after they enter the region of $T > 100$ K. We found that the complex organic species such as HCOOCH_3 are formed mainly via grain-surface reactions at $T \sim 20\text{--}40$ K, and then sublimated to the gas phase when the shell temperature reaches their sublimation temperatures ($T \geq 100$ K). Carbon-chain species can be re-generated from sublimated CH_4 via gas-phase and grain-surface reactions. HCO_2^+ , which is recently detected towards L1527, are abundant at $r = 100\text{--}2,000$ AU, and its column density reaches $\sim 10^{11}$ cm^{-2} in our model. If a core is isolated and irradiated directly by interstellar UV radiation, photo-dissociation of water ice produces OH, which reacts with CO to form CO_2 efficiently. Complex species then become less abundant compared with the case of embedded core in ambient clouds. Although a circumstellar (protoplanetary) disk is not included in our core model, we can expect similar chemical reactions (i.e., production of large organic species, carbon-chains and HCO_2^+) to proceed in disk regions with $T \sim 20\text{--}100$ K.

Keywords. Stars: formation, ISM: molecules

1. Introduction

Chemical reactions in molecular clouds are classified to two categories: gas-phase reactions and grain-surface reactions. Figure 1 schematically summarizes the various chemical processes in molecular clouds. In the cold clouds of $T \sim 10$ K, chemical reactions are triggered by cosmic-ray ionization and proceed mainly via ion-molecule reactions in the gas-phase. While neutral-neutral reactions often have activation barrier, some neutral-neutral reactions do not, and play an important role in forming molecules such as N_2 . When atoms and molecules in the gas-phase collide with grains in the cold molecular clouds, they are efficiently adsorbed onto grain surfaces. The adsorbed species migrate on the grain surface, and react with each other when they meet. Grain-surface reactions are considered to contribute significantly to the formation of large organic species, because grain surfaces are relatively enriched with heavy-element species, and because the association reactions ($\text{AB} + \text{C} \rightarrow \text{ABC}$) are possible by depositing the excess energies on

the grain surfaces. In the gas-phase reactions, on the other hand, products are mostly broken apart ($AB + C \rightarrow A+BC$) to discard the excess energies as kinetic energies of the products.

Efficiency of each process in Figure 1 varies significantly with physical conditions. In prestellar cores with high density ($n_H \geq 10^5 \text{ cm}^{-3}$) and low temperature ($T \sim 10 \text{ K}$), the adsorption overwhelms the desorption, although the non-thermal desorption by impinging cosmic-rays, UV, and reaction heat of grain-surface reactions is not negligible. On the grain-surface, hydrogenation ($A + H \rightarrow AH$) should be the dominant reaction; since hydrogen atoms have relatively low mass and low binding energy onto the grain surface, they can migrate efficiently even under low temperatures.

In protostellar cores, on the other hand, hydrogen atoms are thermally sublimated to the gas phase. Heavy-element species, with higher binding energies (onto grains) than hydrogen atoms, are still on the grain surface, migrating to react with each other. As the core temperature rises further, the heavy element species are also sublimated according to their volatility. The sublimates undergo gas-phase reactions, which have been considered to be the dominant formation path of the large organic species observed in hot cores in high-mass star forming regions (e.g., Millar & Hatchell 1998). At the temperature of 200 K and density of $n_H \sim 10^6 \text{ cm}^{-3}$, for example, large organic species are formed from sublimated CH_3OH and H_2CO by the gas-phase reactions in $\sim 10^5$ yrs. But it is not clear if a similar process works in low-mass protostellar cores, in which large organic species are detected in recent years (Ceccarelli *et al.*, in this volume). In low-mass cores, the hot region ($T > 100 \text{ K}$) with sublimated CH_3OH must be much smaller than in high-mass cores, and thus the infalling material would pass through such hot regions in a shorter timescale than 10^5 yrs. Furthermore, the recent theoretical studies and laboratory experiments on chemical reactions suggest that the gas-phase formation of large organic species would not be as efficient as previously assumed (Horn *et al.* 2004, Geppert *et al.* 2006). Garrod & Herbst (2006) showed grain-surface reaction at warm temperature (\sim several 10 K) can be an alternative or supplemental formation path of large organic species.

Molecular evolution in star-forming cores can be predicted quantitatively by solving the chemical reaction network, which includes the chemical processes mentioned above and are described mathematically as rate equations,

$$\frac{dx(i)}{dt} = \sum_{j,k} \alpha_{j,k} x(j)x(k)n_H + \sum_j \beta_j x(j). \quad (1.1)$$

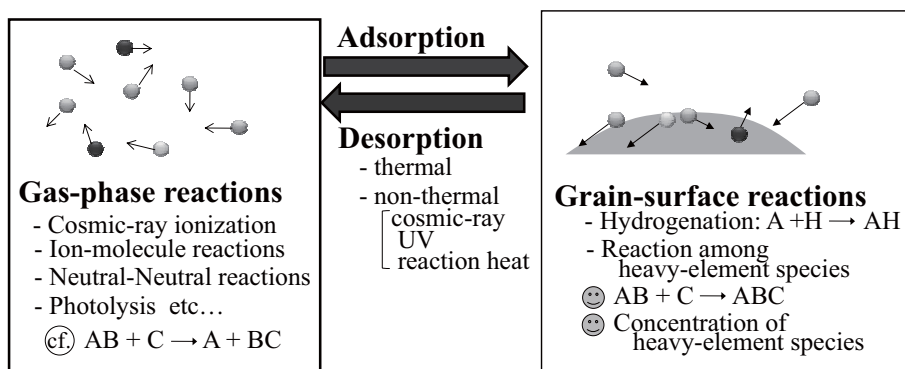


Figure 1. Chemical processes in molecular clouds.

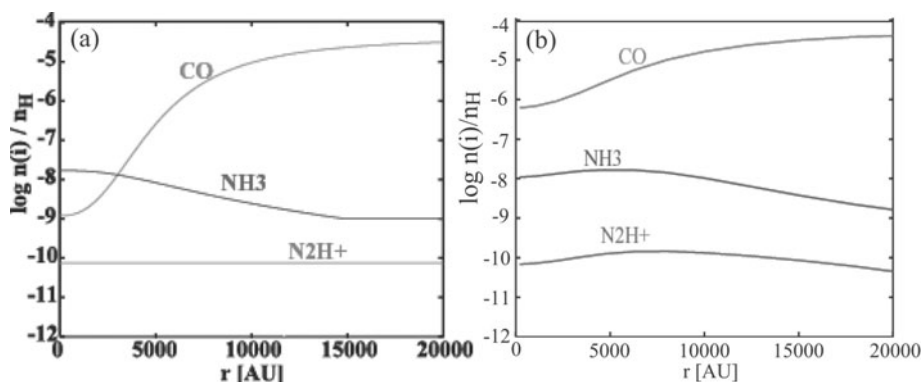


Figure 2. Radial distribution of molecular abundances in prestellar cores. Abundances derived from the observation of L1517B (Tafalla *et al.* 2002) (a) and our model (b).

The reaction rate coefficients, α and β , are in general a function of temperature, and $x(i)$ is the number density of species i relative to hydrogen nuclei (n_H). The solution of Eq. (1) then depends on density (n_H) and temperature, which are often assumed to be constant or a simple function of time. Garrod & Herbst (2006), for example, assumed that the density is constant and temperature rises as a square of the time ($T \propto t^2$). While such simple assumptions on the physical conditions are convenient in investigating the dependence of the chemical reaction network on physical parameters, it is also useful to adopt more realistic models of temperature and density evolution in order to evaluate the temporal and spatial variation of molecular abundances in star-forming cores. Masunaga & Inutsuka (2000), for example, constructed a one-dimensional (spherical) non-gray radiation-hydrodynamics code to calculate the physical evolution of a dense cloud core to form a protostar. Combining the chemical reaction network model with such radiation-hydrodynamics models, we can investigate temporal and spatial variation of molecular abundances in a star forming core (Aikawa *et al.* 2008).

In the following, we will present our model results on a collapsing prestellar core, which reproduces the observed chemical fractionation of C-bearing and N-bearing species (§2). In §3 we investigate the molecular evolution from a prestellar core to a protostellar core, and compare our model results with the observations of large organic species, carbon-chain species and HCO_2^+ in protostellar cores.

2. Prestellar Cores

Over the last several years, radio astronomers found chemical fractionation in prestellar cores. While the intensity map of dust continuum and nitrogen-bearing species (N_2H^+ and NH_3) are centrally peaked, the map of carbon-bearing species show a central hole. Comparison of these intensity maps revealed that C-bearing species are heavily depleted at the central regions, while N-bearing species are not. Figure 2a shows the radial distribution of molecular abundances in L1517B (Tafalla *et al.* 2002).

In order to investigate the cause of this fractionation, we calculated molecular evolution in prestellar cores by combining the chemical reaction network model with the gravitational contraction of a spherical core. Since the radiation cooling is more efficient than the contraction heating at this evolutionary stage ($n_H \leq 10^7 \text{ cm}^{-3}$), we adopt the isothermal collapse model to derive the temporal variation of density in infalling shells, in which the chemical reaction network is solved (Aikawa *et al.* 2005). Our present model includes two important updates from Aikawa *et al.* (2005). Firstly, we set the adsorption

energy of CO and N₂ to be $E_{\text{ads}}(\text{CO})=1180$ K and $E_{\text{ads}}(\text{N}_2)=1060$ K. While previous model works often assumed significantly lower adsorption energy of N₂ than that of CO, e.g. $E_{\text{ads}}(\text{CO})=1780$ K and $E_{\text{ads}}(\text{N}_2)=750$ K, Öberg *et al.* (2005) found via laboratory experiments that the adsorption energy E_{ads} of CO and N₂ are similar. Secondly, we assume N₂H⁺ recombination produces NH + H (10%) and N₂+H (90%), following the recent laboratory experiment (Geppert priv. com.). Previously, the branching ratio was claimed to be 65% for NH + H and 35% for N₂+H (Geppert *et al.* 2004).

Figure 2b shows radial distributions of molecular abundances in our model when the central density of the core is similar to that of L1517B ($\sim 10^5$ cm⁻³). The model shows a reasonable agreement with the observation; in the central region, NH₃ abundance is slightly enhanced, and N₂H⁺ abundance is almost constant, while CO is depleted. In our model, CO is adsorbed onto grains. Since CO is the main reactant of N₂H⁺, the rate of N₂H⁺ destruction decreases as CO depletes. Although N₂ can freeze-out onto grains with a similar adsorption energy of CO, slow (i.e. long-lasting) N₂ formation also helps to keep the abundance of N-bearing species undepleted (Aikawa *et al.* 2005, Maret *et al.* 2006).

3. Protostellar cores

3.1. Physical and chemical evolution

As the core contracts further, the contraction heating overwhelms the radiation cooling, and core temperature rises. Eventually, the protostar is formed, which further heats the envelope (i.e. protostellar core). In order to investigate the molecular evolution in protostellar cores, we adopt a core model of Masunaga & Inutsuka (2000). The model starts from a hydrostatic cloud core with the central density of $n_{\text{H}} \sim 6 \times 10^4$ cm⁻³, and follows the core evolution until 9.3×10^4 yrs after the protostar is born ($\equiv t_{\text{final}}$). Although the original paper by Masunaga & Inutsuka (2000) mainly discusses the formation of a protostar, the model gives the radial distribution of density, temperature and infall velocity in the entire core as a function of time. Figure 3 shows the structure of the core model at assorted evolutionary stages. It should be noted that CO and large organic species such as CH₃OH are sublimated to the gas phase at ~ 20 K and ~ 100 K, respectively. At the moment of protostar birth ($t = 0$), for example, CO is sublimated inside ~ 100 AU.

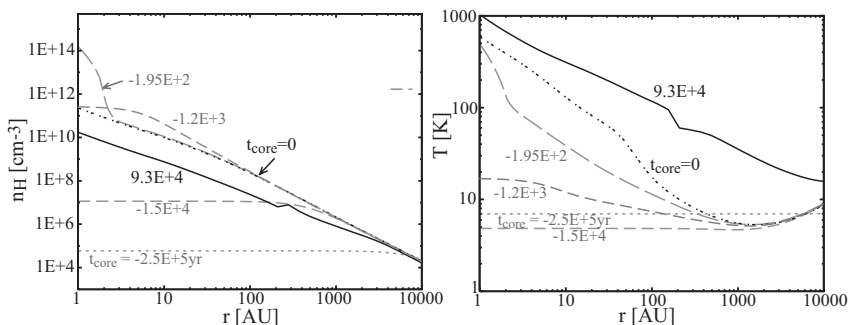


Figure 3. Radial distribution of density and temperature in a star-forming core at assorted evolutionary stages. The protostar is born at $t = 0$; $t < 0$ (gray lines) corresponds to the prestellar stage, while $t > 0$ (black lines) corresponds to the protostellar stage. The core is initially almost isothermal, but the central region starts to warm up as the contraction heating overwhelms the radiation cooling ($t \sim -1.2 \times 10^3$ yr). The first core, i. e. an AU-sized hydrostatic core, appears at $t \sim$ several 100 yr, which then collapses to form a protostar.

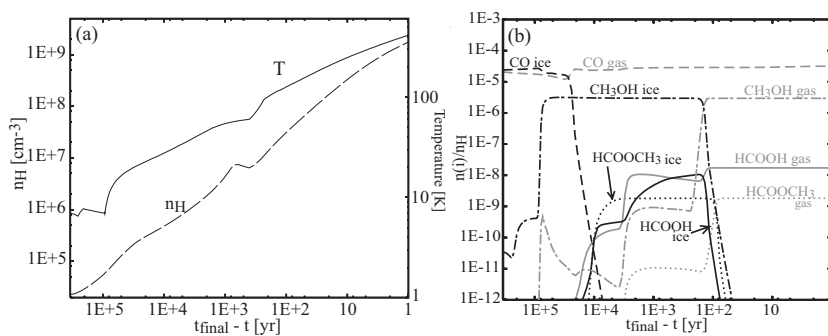


Figure 4. (a) Temporal variation of density and temperature in the infalling shell which reaches $r = 2.5$ AU at t_{final} . (b) Evolution of molecular abundances in the shell. Black lines represent the species adsorbed onto grain surfaces (i.e. ice), while gray lines represent gaseous species.

Based on the core model described above, we derive the temporal variation of density and temperature in infalling shells; Figure 4a shows such variation in the shell which reaches $r = 2.5$ AU at t_{final} . The temporal variation accelerates as the shell migrates towards the inner radius, where the infall velocity, temperature, and density steeply increases inwards. In order to highlight the rapid temporal variation near and at the final stage, the horizontal axis of Figure 4 is set to be the logarithm of $t_{\text{final}} - t$. We can see that once the shell enters the region of $T \geq 100$ K, it falls onto the central star within ~ 100 yrs, which is much shorter than the timescale ($\sim 10^5$ yr) needed to form the large organic species from sublimated CH $_3$ OH via gas-phase reactions (§1).

Figure 4b shows the molecular evolution in the infalling shell. As initial molecular abundances, we adopt the “low-metal” elemental abundances and assumed that the species are in the form of atoms or atomic ions except for hydrogen, which is in its molecular form. Assuming that the initial cloud core is stable for 1×10^6 yrs, we solved the chemical reaction network with the fixed density and temperature over this period, which sets the initial molecular abundances for the collapse stage shown in Figure 4. We can see that large organic species such as HCOOCH $_3$ are mainly formed on the grain surfaces at $T \sim 20$ – 40 K, and then sublimate to the gas-phase when the shell temperature reaches their sublimation temperatures. One exception is HCOOH, which is formed by the gas-phase reaction of sublimated H $_2$ CO with OH.

Following the molecular evolution in multiple infalling shells, we obtain spatial distribution of molecules in a protostellar core. Figure 5 shows the radial distribution of molecular abundances at t_{final} . Gaseous organic species are abundant inside about 100 AU, where the temperature exceeds their sublimation temperatures. Formic acid is again an exception; it extends beyond its sublimation radius because of the gas-phase formation. The organic species on grain surfaces are abundant at 100 AU–1,000 AU.

3.2. Effect of UV radiation: embedded core vs isolated core

Although we are interested in the central region of protostellar cores, the shells are initially (in the prestellar core stage) at outer radii $r \sim 10^4$ AU. In the model described in the previous subsection, we assumed that the model core is embedded in ambient clouds, and set the visual extinction of $A_{\text{v}} = 3$ mag at the outer boundary of the core. But this assumption is arbitrary; Bok Globules, for example, are not surrounded by ambient clouds and thus can be exposed directly to interstellar UV radiation. In order

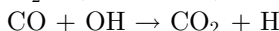
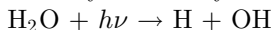
Table 1. Gas-phase molecular abundances in IRAS 16293-2422 and model results.

Species	IRAS 16293-2422	model	
		embedded	isolated ^a
H ₂ CO	1.0(-7) ^b , 1.1(-7) ^c	2.8(-6)	1.3(-11)
CH ₃ OH	1.0(-7) ^d , 9.4(-8) ^c	3.0(-6)	8.5(-8)
HCOOCH ₃	2.5-5.5(-7) ^e , 2.6-4.3(-9) ^f , > 1.2(-8) ^g	1.8(-9)	3.2(-11)
HCOOH	6.2(-8) ^e , 2.5(-9) ^g	1.7(-8)	2.1(-8)
CH ₃ OCH ₃	2.4(-7) ^e , 7.6(-8) ^c	3.5(-10)	5.8(-12)
CH ₃ CN	1.0(-8) ^e , 7.5(-9) ^h	3.0(-8)	1.3(-8)

Notes:^a Gas-phase abundances at $r = 30.6$ AU are listed, because the abundances are mostly constant at $r \lesssim 100$ AU.^b Maret *et al.* (2004), ^c Chandler *et al.* (2005), ^d Maret *et al.* (2005), ^e Cazaux *et al.* (2003), ^f Kuan *et al.* (2004),^g Remijian & Hollis (2006), ^h Schöier *et al.* (2002).

to investigate the effect of UV radiation on core chemistry, we calculated a model with $A_V = 0$ mag at the outer core edge.

We found that the isolated core model evolves to be a protostellar core with higher abundance of CO₂ and lower abundance of large organic molecules compared with the embedded core model (Aikawa *et al.* 2008). While the shells are at outer radii ($\sim 10^4$ AU) of the prestellar core, CO₂ are efficiently formed by the reactions of

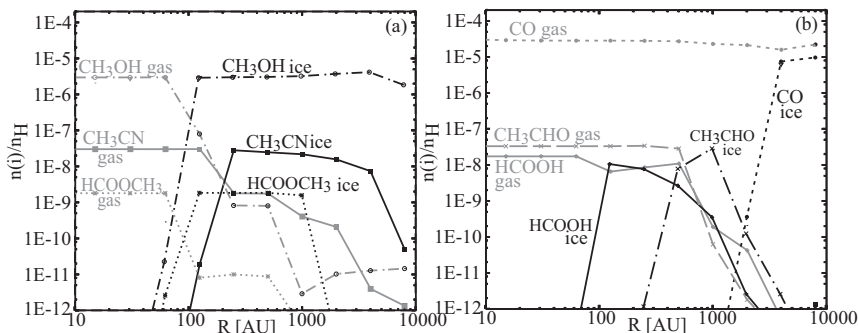


on the grain surfaces. In the observation of protostellar cores, molecular abundances are known to vary among objects. Our results indicate that such variation could, at least partially, result from the incident UV radiation at the prestellar stage.

3.3. Comparison with observation

Table 1 compares molecular abundances in our embedded and isolated core models and those estimated from the observations of a class 0 object IRAS 16293-2522. It should be noted that the estimated abundances vary significantly depending on the assumptions on temperature and density distribution in the protostellar core. Considering such uncertainties, our models are in reasonable agreement with the observations.

In order to understand the chemical conditions in protostellar cores, it is also important to observe species other than large organic species. Recently, interesting chemical species are detected towards L1527, which is considered to be in the transition phase from class 0 to class I. Firstly, Sakai *et al.* (2008a) detected carbon chains (e.g., C₄H, C₄H₂, and C₃H₂); C₄H emission extends over the 40'' scale, and the line width toward the core center is broader than those toward the surrounding positions, indicating that C₄H emission

**Figure 5.** Radial distribution of molecules in our protostellar core model at t_{final} .

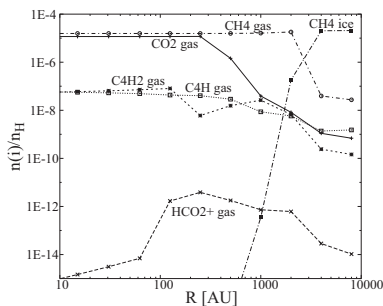


Figure 6. Radial distribution of carbon chains and HCO_2^+ in our protostellar core model at t_{final} .

is continuously distributed from the infalling envelope to the inner part of the core. Existence of carbon-chain species towards L1527 was a surprise, because carbon-chain species are usually associated with the early stages of cloud cores when the dominant form of carbon changes from atomic carbon to CO. Sakai *et al.* (2008a) proposed that carbon chains could regenerate from sublimated CH_4 in the warm central region. Indeed, carbon chains are abundant in our model (Figure 6). Methane is produced in the prestellar core phase and sublimate as the infalling shell temperature rises to 25 K. Then a fraction of sublimated CH_4 is transformed to carbon chains via both gas-phase and grain-surface reactions.

Secondly, Sakai *et al.* (2008b) detected HCO_2^+ towards L1517. While CO_2 ice is known to be abundant in protostellar cores (ex. 22% relative to H_2O ice towards Elias 29) (Ehrenfreund & Shutte 2000), the abundance of gaseous CO_2 has not been well-constrained. Since HCO_2^+ is formed by the reaction of $\text{H}_3^+ + \text{CO}_2 \rightarrow \text{HCO}_2^+ + \text{H}_2$, it can be a tracer of gaseous CO_2 . Although the spatial distribution of HCO_2^+ is not well-constrained from the observation, Sakai *et al.* (2008b) argue, with a help of simple chemical analysis, that the HCO_2^+ seems to extend upto ~ 2000 AU. The column density of HCO_2^+ is estimated to about $7.6 \times 10^{10} \text{ cm}^{-2}$.

In comparison, distribution of HCO_2^+ in our model is shown in Figure 6; it is abundant at radius of $100 \text{ AU} < r < 2000 \text{ AU}$. Column density of HCO_2^+ reaches 10^{10} – 10^{11} cm^{-2} towards the core center (impact parameter $\leq 2000 \text{ AU}$), which is consistent with the observation. In our model, gaseous CO_2 and HCO_2^+ increase at the region of $\sim 25 \text{ K}$, where some carbon-chains react with O_2 to form CO_2 in the gas-phase. CO_2 ice is abundantly formed by grain-surface reactions, but they do not desorb efficiently until the shell temperature reaches $\sim 70 \text{ K}$. In the central region of $r \leq 100 \text{ AU}$, on the other hand, CO_2 gas is abundant but HCO_2^+ is not, because NH_3 and H_2O , which have larger proton affinity than CO_2 , are sublimated to the gas phase.

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Discussion

BOTTINELLI: I was wondering what abundance of ammonia you had in your model and whether you can see difference in the chemistry between ammonia rich and ammonia poor ice.

AIKAWA: In our model the ammonia abundance is not assumed. It is determined by the chemical reaction network, and is about 10% of H₂O ice in outer radius.

CUPPEN: I have a question about the accretion model, your updated version. You show that ammonia is in agreement with the observations. Can you comment on the CO?

AIKAWA: CO in the gas phase is determined by the efficiency of cosmic ray desorption. This is one of the assumptions in the model. Observationally, we do not have a sensitivity to distinguish between 10⁻⁷ abundance to 10⁻⁶ abundance. Observers can only say that it is lower than 10⁻⁶.

CECCARELLI: I saw that your model prediction of the abundance of large organic molecules is around 100 times lower than we observed.

AIKAWA: In our model we need methanol to produce these molecules. I want to emphasize that there is no assumptions on initial conditions or ice abundance. I just started with a very simple protostellar core and solved detailed but still simple chemical reaction network, to see how far we can go and how much organic species we can make. If you need more, I have to work on the chemical network to make more organics from methanol. Another solution could be axisymmetric model, rather than spherical. Then the fluid parcels can stay longer in warm regions (of circumstellar disk) in which large organic molecules are formed.