

# **Quantum dark energy in a seven-dimensional universe**

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## **Abstract:**

This paper explains the dark energy and acceleration of the universe by quantizing the space in hidden dimensions, which provides the basis and background for the gravitational force through the curvature of space-time. Space-time is considered to be made of a four-dimensional elastic grid in a seven-dimensional universe in which matter also expands along with the universe. Each cube of the grid is considered a quantum of hidden three-dimensional space of Planck volume containing Planck charge, which makes the universe seven-dimensional. The dark energy is explained by the electrostatic repulsion between the Planck charges in each quantum of the hidden space. Mathematically, this electrostatic repulsion is related to the Hubble constant to explain the accelerated expansion, dark energy, and increase in the cosmological potential energy of matter. Expansion of space-time is considered not due to the creation of the new space but due to the stretching of the existing space-time itself like an elastic ruler where the proper length and the volume remains constant. The values of the Planck constant, gravitational constant, permittivity of free space, and Boltzmann constant are shown to vary owing to the expansion of space-time and hence provide falsifiable predictions for this theory. This theory builds a framework for the relativistic Newtonian theory of gravity, the relativistic MONDian (Modified Newtonian dynamics) gravity, identifies a valid reason for the transition of Newtonian gravity to MOND at  $a_0$  and eliminates the need for dark matter to explain the dynamics of galaxy clusters.

## 1. Introduction

In this theory of quantum dark energy in a seven-dimensional universe, Hubble expansion is considered to be due to the stretching of the existing space-time rather than the creation of a new space that is analogous to the markers on an elastic ruler stretching along with the expansion of the ruler, as opposed to the raisin bread model, where matter does not expand along with space. Therefore, the matter is considered to uniformly stretch along with the space-time and not become diluted with the expansion of the universe. As the matter also stretches with space, the expansion is nonobservable locally but can be observed through the redshift of the light coming from a far-off space of a lower stretch. At the outset, we can see that matter exists in different energy states based on the magnitude of the space-time stretch. The matter continues to move to higher potential energies as space-time expands, but the energy of the photon remains the same. Therefore, in a lower stretch space-time, a blue photon has less energy compared to the same blue photon emitted in a higher stretch space-time. Therefore, when light travels from a lower stretch of space-time to a higher stretch of space-time, it becomes redshifted without violating the law of energy conservation. The cosmological potential energy of the mass increases with the space-time expansion, whereas the energy of the photon does not change with the space-time expansion, and it does not gravitate, as it does not resist the space-time expansion because the mass of the photon is zero but is affected by the gravity of the other objects, which means that the energy of a photon does not curve the space-time around it but takes the path of the curved space around any matter. Thus, electromagnetic fields and electromagnetic-related potential energies whose forces are mediated by virtual photons do not curve the space-time in this model of the universe. As photons do not curve the space-time in this theory, the annihilation of matter would convert the cosmological potential energy, gravitational potential energy, and the rest mass to photons.

This theory proposes an absolute inertial frame, a way to identify it, and hence revives the concept of relativistic mass. In this theory, only the rest/relativistic mass or the mass density is considered to be responsible for gravitational force (gravitoelectric) and the relativistic mass current density (mass flux) for the gravitomagnetic force. Any other form of energy increases neither the inertial mass nor the gravitational mass in this model of the universe but only the rest/relativistic mass and the relativistic mass flux. Therefore, this theory only partially satisfies the mass-energy equivalence principle but upholds the weak equivalence principle.

The proper distance between any two points in space is considered to remain constant despite the stretching of space-time, similar to the markers on an elastic ruler. Therefore, the volume of the universe does not change with the expansion of space. The expansion of space-time and matter is considered to be due to electrostatic repulsion between the Planck charges present in the Planck volumes and hence discards dark energy. Space-time is considered to be made of a four-dimensional elastic grid in a seven-dimensional universe. Each cube of the grid is considered to be a quantum of three-dimensional space, equivalent to the Planck volume and containing Planck charge. These cubes can become loosely connected to the surrounding cubes for the formation of black holes. As the Hubble constant  $H_0$  decreases with time, the rate of acceleration of the universe is considered to be decreasing because of the contraction force of the space-time elastic grid opposing the electrostatic repulsion between Planck charges. The space-time membrane acts as a dielectric material between the Planck charges. Matter only exists as a probability wave function ( $\Psi$ ) in the four-dimensional space-time grid but not in the hidden three spatial dimensions. The presence of matter in space-time would increase the force required to expand the space-time grid as the matter also expands along with space.

The presence of matter increases the permittivity of space-time and hence reduces the electrostatic repulsion within the grid enveloped by the matter compared to the electrostatic repulsion outside the matter. As the expansion of the space-time surrounding the matter will be greater than the expansion of the matter due to the permittivity difference and the net compressing force on the matter from the surrounding Planck charges, space-time becomes naturally curved around any mass, and hence the gravitational force. Therefore, in this theory, except for the photons/electromagnetic fields, the gravitational force is considered a real force rather than a fictitious force, as in the general theory of relativity, due to the curvature of space-time. In this theory, the electrostatic potential energy between the Planck charges in the three-dimensional space is considered to be the same as the dark energy, which causes accelerated expansion of the universe. The electrostatic dark energy in three-dimensional space does not gravitate because this energy itself is the cause of the gravitational force in the four-dimensional space-time enveloping the three-dimensional Planck charges. The constancy of the speed of light should not limit the apparent velocity of the universe in this model, as it considers it to be applicable only for objects moving through space but not for the expansion of space-time. Beginning from the birth of the universe, the first law of thermodynamics is strictly followed in this model of the universe to uphold the law of conservation of energy, which includes energy conservation in dark energy and cosmological redshift. Linear momentum and angular momentum are also conserved in this theory.

This theory proposes a flat or zero-curvature, isotropic, and homogenous universe, where the rate of acceleration will continue to decrease proportionally to the age of the universe, and its velocity only becomes zero after an infinite amount of time. However, the universe will continue to accelerate, but only at a continuously decreasing rate that asymptotes to zero. Invoking the critical density ( $\Omega_0 = 1$ ) or cosmic inflation is not required to explain the flatness in this model of the universe, as the uniform expansion of the whole universe due to electrostatic repulsion explains why the universe is flat rather than closed or open. As there is no increase in the volume of the universe with time, the Big Bang should be replaced with the Big Repulsion. This theory makes the gravity electrostatic background dependent to explain the curvature of space-time, dark energy, and cosmological redshift. The electrostatic background provides an additional background to the quantum fields on top of the gravitational background, which mandates an absolute inertial frame and universal time, which is Hubble time in this model of the universe. Therefore, this theory makes the universe three-layered.

- 1) Electrostatic or the power layer
- 2) Space-time or the gravitational layer
- 3) Quantum fields layer

The first layer acts as a power or energy source of the universe. The gravitational layer creates gravity, and the quantum field layer enables matter, energy, and the rest of the fundamental forces of nature to work on top of the space-time or the gravitational layer.

This theory also eliminates the cosmic event horizon and the cosmic scale factor ( $a$ ) because the proper length and volume remain constant despite the expansion of space-time. Therefore, we should be able to see light from any part of the universe without any distance limit, such as a cosmic event horizon. The cosmological constant  $\Lambda$  in the gravitational field equations to factor in the dark energy is not required in this model of the universe, as dark energy does not gravitate and is isolated to the electrostatic background, which is handled independently from gravity.

The values of the Planck constant, gravitational constant, Boltzmann constant, and permittivity of free space are shown to vary owing to the expansion of the space-time grid, proving the existence of an absolute frame of reference, and hence providing falsifiable predictions for this theory. The aforementioned constants are also shown to vary with the gravitational potential and in all the moving inertial frames and hence enable us to identify the frames through the change in values of the constants compared to the absolute inertial frame or an inertial frame of a different velocity. An absolute inertial frame is one where the values of the constants are minimum or maximum depending on how they vary with the gravitational potential or the velocity of the inertial frame. Therefore, this theory completely resolves the twin paradox: the person in the frame with the change in physical constants due to velocity will age less than the one in the frame with no change in the physical constants after they meet. However, this theory upholds the invariance of the speed of light in all internal and noninertial frames and hence upholds the special theory of relativity except for the absolute reference frame.

In this theory, the constants that do not change with the expansion of the universe, the change in gravitational potential, or the velocity of the frame are listed below but are not limited to.

- 1) Speed of light
- 2) Planck length
- 3) Planck time
- 4) Planck temperature
- 5) Electric charge
- 6) Fine-structure constant ( $\alpha$ )
- 7) Rest mass

Black hole singularities do not exist in this model of the universe as the mass is converted to pure informational entropy at the event horizon and the space-time terminates at the event horizon and cannot be extended beyond the event horizon as in the general theory of relativity as the effective radial length, time and the tangible mass becomes zero at the event horizon and hence discards the Riemannian geometry of space-time. Additionally, charged black holes do not exist in this model of the universe, as the permittivity of free space becomes infinity at the event horizon. Therefore, stationary black holes can only have two properties, namely, informational (entropic) mass and angular momentum, and hence only partially satisfy the no-hair theorem.

This theory also builds a framework for the covariant/relativistic Newtonian theory of gravity. Relativistic Newtonian gravity proposed in this theory produces similar or even better results in closed form than the general theory of relativity without using weak field approximations for the below listed.

- 1) Black hole radius (Schwarzschild radius)
- 2) Photon sphere
- 3) Gravitational lensing
- 4) Perihelion precession of Mercury
- 5) Shapiro time delay

This theory generates the relativistic Poisson equations and Maxwell-like equations for gravity or GEM (Gravitoelectromagnetism) equations to make them work in strong gravitational fields and relativistic speeds as well, which can explain the frame-dragging effect, orbital precession, geodetic effect, and gravitational waves.

This theory also generates the equivalent of the Schwarzschild and FLRW (Friedmann–Lemaître–Robertson–Walker) metrics. The expansion of the universe proposed in this theory identifies a valid reason for the transition of Newtonian gravity to MOND at  $a_0$  ( $\sim 1.2 \times 10^{-10} \text{ m/s}^2$ ) and hence makes a case for the absolute reference frame. However, this theory establishes that transitioning to a deep-MOND regime is only possible in the radial space around the black holes but not around the ordinary matter and hence explains the missing gravitational lensing around the gaseous part of the Bullet cluster (1E 0657-56), and eliminates the need for dark matter to explain the dynamics of galaxy clusters using an updated virial theorem. This theory also builds a framework for relativistic MONDian gravity.

In addition, this new model could act as a precursor to theories explaining baryogenesis and primordial nucleosynthesis based on how the initial extremely high expansion energy of space-time interacted with the quantum fields to create matter. However, it still needs to be seen if the high value of the gravitational constant  $G$  proposed in this theory during the emission of CMB (cosmic microwave background) accounts for the observed anisotropy and the angular power spectrum without the dark matter, as high  $G$  should produce the same gravitational effect as the equivalent high mass. Additionally, this theory explains the early formation of the galaxies/objects that were recently observed through the JWST (James Webb Space Telescope) by proving that they were formed much later than calculated.

## 2. Relativistic acceleration of space-time

Let us consider two points, A and B, on the space-time fabric. We can calculate the apparent outward acceleration of point B when observed from point A based on the redshift of the light coming from point B.

$D$  = Proper distance between points A and B

$\lambda = D$  (let us consider the wavelength of light to be equal to  $D$ )

Thus, by the time light travels from point B to point A, its wavelength would have expanded by  $D(1+z)$  based on the cosmological redshift ( $z$ ) phenomenon. Therefore, point B would have apparently moved from its original location by  $D(1+z)-D$ , which is equal to  $Dz$ .

The apparent velocity  $v$  of point B due to the redshift is given by  $v = \frac{\text{Distance}}{\text{time}} = \frac{Dz}{t}$ , where  $t$  is the time taken for the light to travel from point B to point A. As the proper distance between points A and B remains constant, the real velocity of point B is zero. Thus, apparent acceleration  $a$  is given by  $a = \frac{Dz}{t^2}$ . As space stretches like an elastic ruler, the proper distance between points A and B should always remain the same, irrespective of the apparent acceleration. Therefore, the light should take the same amount of time  $t$  to travel the apparent distance of  $D(1+z)$ , which is actually  $D$  owing to the constancy of the speed of light. Therefore, time  $t$  is given by  $t = D/c$

$$\text{So, } v = \frac{Dz}{\left(\frac{D}{c}\right)} = zc \text{ and } a = \frac{Dz}{\left(\frac{D}{c}\right)^2} = \frac{zc^2}{D}$$

$$v = zc \quad (1)$$

As  $v = zc$  has already been established in conjunction with the Hubble law, the above derivation proves that space is stretching like an elastic ruler, where the length and volume remain constant, as opposed to the raisin bread model, where length increases and the matter is diluted with the expansion of space-time. Therefore, this theory discards the raisin-bread model of the universe and provides a theoretical basis for  $v = zc$  (1) whereas the raisin-bread model does not provide any theoretical reasoning.

Based on the equations  $v = H_0 D$  and  $v = zc$  from Hubble's law,  $\frac{z}{D} = \frac{H_0}{c}$ , where  $v$  is the apparent receding velocity,  $H_0$  is the Hubble's constant,  $D$  is the distance,  $z$  is the cosmological redshift, and  $c$  is the speed of light. Therefore, the apparent acceleration  $a = \frac{H_0 c^2}{c} = cH_0 = 7.549 \times 10^{-10} \text{ m/s}^2$ ,

which is the current acceleration of the universe for  $H_0 = 77.7 \text{ (km/s)/Mpc}$  based on the  $H_0$  values within the range mentioned in the references (Chen et al., 2019; de Jaeger et al., 2020; Tully et al., 2016). As the universe is stretching with constant volume, there must be length-contraction and time dilation in the past, which are given below. As we are describing the past events from the perspective of the current cosmological reference frame and to maintain the constancy of the speed of light in all reference frames, both the length and time are divided by  $(1+z)$ .

$$\text{Cosmological length-contraction: } L = \frac{L_0}{(1+z)} \quad ; \quad \text{Cosmological time-dilation: } T = \frac{T_0}{(1+z)}$$

The cosmological time-dilation proposed in this theory is inherent to the fabric of space-time due to the cosmological potential similar to the gravitational time-dilation due to the gravitational potential and the associated redshift is the cosmological redshift due to the stretching of space-time. Therefore, the relativistic acceleration of the universe is below.

$$a = cH_0(1+z) \quad (2)$$

Therefore, point B will always move from point A with an apparent acceleration equal to the above, although the proper distance between the two points always remains the same. As Hubble's constant decreases over time, the rate of apparent acceleration of the universe or space-time is also considered to decrease over time. As the proper distance between the two points and the volume of the universe always remains constant, acceleration  $a = cH_0(1+z)$  is only considered apparent.

### 3. Relativistic Hubble law and the cosmological redshift

Based on the relativistic acceleration of space-time  $a = cH_0(1+z)$ , we can calculate the relativistic velocity  $v$  of point B from the big repulsion (point A) on the space-time fabric based on the cosmological redshift  $z$ .

$$\begin{aligned} \frac{dv}{dt} &= a; \quad dv = a dt; \quad dv = cH_0(1+z)dt; \\ \int_0^v dv &= \int_{t_p}^t cH_0(1+z)dt; \quad \int_0^v dv = c \int_{t_p}^{T_0} \frac{1}{T} \left(1 + \frac{v}{c}\right) dT \quad \text{from (1)} \end{aligned}$$

$$\int_0^v \frac{1}{(c+v)} dv = \int_{t_p}^{T_0} \frac{1}{T} dT$$

$$\ln(c+v) - \ln(c) = \ln(T_0) - \ln(t_p)$$

Planck time  $t_p$  is the minimum age of the universe at the beginning of the big repulsion owing to the quantization of space and time to Planck units.  $T$  is the age of the universe, which is considered to be Hubble time in this model of the universe. As  $\ln(0)$  is undefined, the singularity or the zero time of the big repulsion is also undefined.

$$\frac{c+v}{c} = \frac{T_0}{t_p} ; \quad 1 + \frac{v}{c} = \frac{1}{H_0 t_p} ; \quad 1+z = \frac{1}{H_0 t_p}$$

Therefore, the maximum possible cosmological redshift is below.

$$z = \frac{1}{H_0 t_p} - 1 \quad (3)$$

which is  $7.36614 \times 10^{60}$  for  $H_0 = 77.7$  (km/s)/Mpc, and the maximum apparent relativistic recession velocity that is possible is  $7.36614 \times 10^{60}$  times the speed of light.

$$\text{For } H_0 < H_D \quad 1+z = \frac{H_D}{H_0} = \frac{1}{H_0 \left( \frac{1}{H_0} - \frac{D}{c} \right)} = \frac{1}{1 - \frac{H_0 D}{c}} = \frac{c}{c - H_0 D} \quad (4)$$

where  $\frac{1}{H_D} \geq t_p$  and  $H_D$  is Hubble's constant when light is emitted at distance  $D$ .

Therefore, the relativistic Hubble law is below based on (1) and (4).

$$v = H_0 D (1+z) = \frac{H_0 D}{\left( 1 - \frac{H_0 D}{c} \right)} \quad (5)$$

Acceleration at the beginning of the universe (Big Repulsion) is below.

$$a = c H_0 (1+z) = c H_0 \frac{1}{H_0 t_p} = \frac{c}{t_p} = 5.56 \times 10^{51} \text{ m/s}^2 \text{ from (2) and (3)}$$

The acceleration of the universe at distance  $D$  from the present is below.

$$a = c H_0 (1+z) = c H_0 \left( \frac{c}{c - H_0 D} \right) \text{ from (2) and (4)}$$

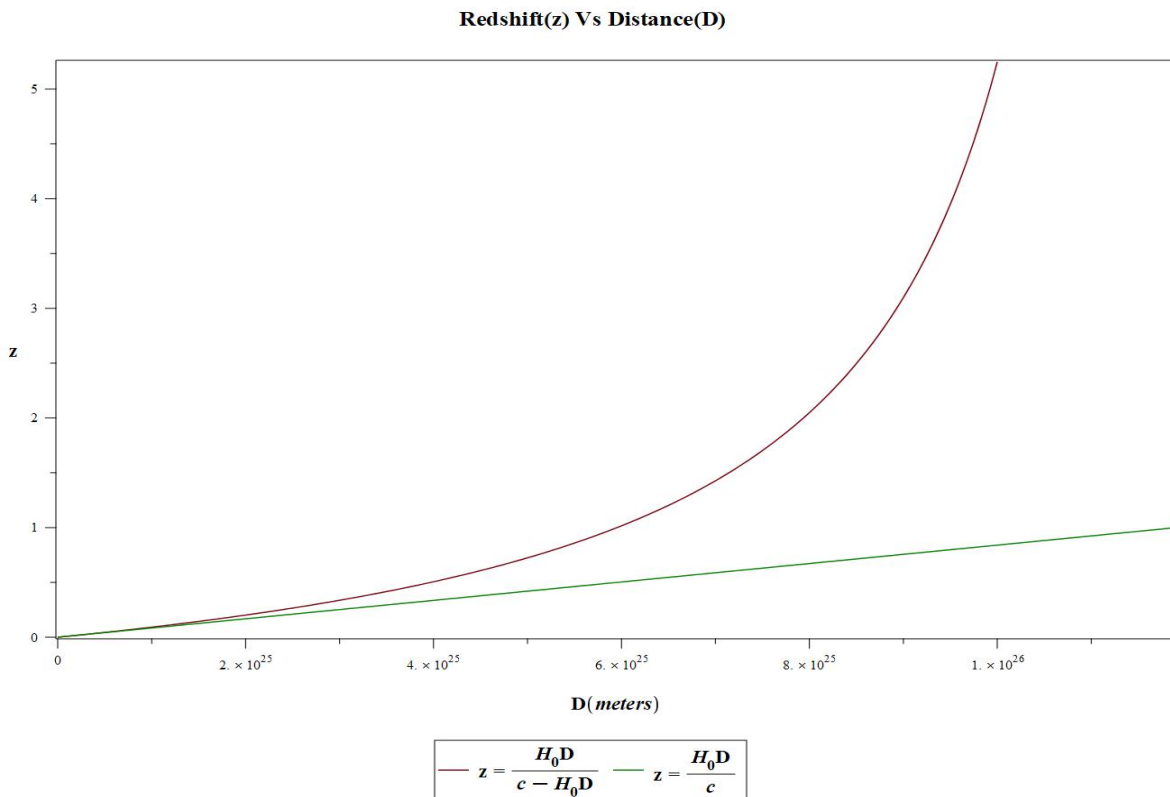
$$a = c H_0 \left( \frac{c}{c - H_0 D} \right) \quad (6)$$



We can see in the table below that the redshift ( $z$ ) values match the regular Hubble's formula and the new relativistic formula for low values of  $D$ . The new formula restricts the maximum cosmological redshift to  $7.36614 \times 10^{60}$  owing to the model of the universe that is considered. As shown in the table below, the  $z$  values differ from each other as  $D$  increases or as the time traveled by the light approaches the age of the universe. Therefore, Hubble's law  $v = H_0 D$  is only accurate up to moderate distances, as it is the limiting case of the relativistic Hubble law. We can also see that the new  $z$  values in Table I are in line with the accelerating model of the universe.

**Table I (Cosmological redshift values)**

$D$ (meters)	$z = \frac{H_0 D}{c}$	$z = \frac{H_0 D}{c - H_0 D}$
$1 \times 10^{10}$	$8.3994291420 \times 10^{-17}$	$8.3994291420 \times 10^{-17}$
$2 \times 10^{15}$	$1.6798858284 \times 10^{-11}$	$1.6798858284 \times 10^{-11}$
$3 \times 10^{20}$	$2.5198287426 \times 10^{-6}$	$2.5198350921 \times 10^{-6}$
$\sim 5.9527854 \times 10^{25}$	0.5	1
$1 \times 10^{26}$	0.8399429142	5.2477708815
$1.19 \times 10^{26}$	0.9995320679	2136.0622240085
$\sim 1.1905570 \times 10^{26}$	1	$7.3661442125 \times 10^{60}$ (Maximum/Big Repulsion)





The age of the universe is 12.58 billion years, which is the Hubble time. As the proper distance and the volume remain constant, the maximum observable universe is only  $12.58 \times 2$  billion light years across, which is 25.16 billion light-years for  $H_0=77.7$  (km/s)/Mpc.

Important formulas: based on (4)

Time since the big repulsion to the emission of light ( $T_b$ )	Time since the emission of light ( $T_z$ )	Distance vs Redshift
$T_b = \frac{1}{H_0(1+z)}$	$T_z = \frac{1}{H_0} \frac{z}{(1+z)}$	$D = \frac{zc}{H_0(1+z)}$

Therefore, the light from galaxy HD1 with cosmological redshift  $z=13.27$  (Zhe et al., 2022) should have been emitted after 0.8 billion years since the big repulsion, possibly giving enough time for it to form as a galaxy by also factoring in the high initial gravitational constant  $G$ , as mentioned in section 4 of this document, and hence resolving the issue of the early formation of galaxies/objects that were recently observed through the JWST.

#### 4. Variable constants $G$ , $h$ , $\epsilon_0$ , and $k_B$ and the mass increase

As the speed of light is constant, the Planck length and Planck time are considered constants. Therefore, the product of the gravitational constant  $G$  and Planck constant  $h$  is considered to be constant. As the Planck charge ( $q_p = \frac{e}{\sqrt{\alpha}}$ ) is conserved, alpha ( $\alpha$ ) or the fine-structure constant is considered to be constant. Therefore, the product of the Planck constant  $h$  and permittivity of free space  $\epsilon_0$  is considered to be constant. As the Planck temperature is considered to be constant, the product of the gravitational constant  $G$  and Boltzmann constant  $k_B$  is considered to be constant. Based on the above, we can calculate how constants  $G$ ,  $h$ ,  $\epsilon_0$ , and  $k_B$  change over time. Using the law of conservation of energy, we can calculate the values of  $G$ ,  $h$ ,  $\epsilon_0$ , and  $k_B$  in the past and the future.

$$\text{Cosmological redshift } z = \frac{\lambda_{obs} - \lambda_{emit}}{\lambda_{emit}} \text{ or } \lambda_{obs} = \lambda_{emit}(1+z) \quad (7)$$

$$\text{As the energy is conserved in cosmological redshift, } E = \frac{hc}{\lambda_{obs}} = \frac{h_{t_p} c}{\lambda_{emit}} \quad (8)$$

Here,  $h$  is the current Planck constant, and  $h_{t_p}$  is the old Planck constant when the age of the universe is Planck's time  $t_p$ .

$$z = \frac{1}{H_0 t_p} - 1 \text{ from (3), and } h = h_{t_p}(1+z) \text{ based on (7) and (8)} \quad (9)$$

$$\text{As } hG = h_{t_p} G_{t_p}, h\epsilon_0 = h_{t_p} \epsilon_{t_p} \text{ and } k_B G = k_{B_{t_p}} G_{t_p}$$

$$h = h_{t_p} (1+z) = h_{t_p} \left( \frac{1}{H_0 t_p} \right) \quad (10)$$

$$G = \frac{G_{t_p}}{(1+z)} = G_{t_p} (H_0 t_p) \quad (11)$$

$$\varepsilon_0 = \frac{\varepsilon_{t_p}}{(1+z)} = \varepsilon_{t_p} (H_0 t_p) \quad (12)$$

$$k_B = k_{B_{t_p}} (1+z) = k_{B_{t_p}} \left( \frac{1}{H_0 t_p} \right) \quad (13)$$

Below are the values of  $G_{t_p}$ ,  $h_{t_p}$ ,  $\varepsilon_{t_p}$ , and  $k_{B_{t_p}}$  in MKS units when the age of the universe was Planck time  $t_p$ , based on the current values of  $G$ ,  $h$ , and  $\varepsilon_0$  for  $H_0=77.7$  (km/s)/Mpc. Therefore, these constants vary proportionally to the relativistic Hubble flow.

$h = 6.62607015 \times 10^{-34}$	$G = 6.67430 \times 10^{-11}$	$\varepsilon_0 = 8.8541878128 \times 10^{-12}$	$k_B = 1.380649 \times 10^{-23}$
$h_{t_p} = 8.99530332 \times 10^{-95}$	$G_{t_p} = 4.91638563 \times 10^{50}$	$\varepsilon_{t_p} = 6.52212243 \times 10^{49}$	$k_{B_{t_p}} = 1.87431709 \times 10^{-84}$

As the factor  $(1+z)$  in all the above four constants changes directly in proportion to mass ( $m$ ) in the dimensional formulas, it follows that  $m_0 = m_z (1+z)$  as  $[E = h(1+z)v = m(1+z)c^2]$ , where  $m_0$  is the current mass and  $m_z$  is the original mass when the light was emitted in the past at cosmological redshift  $z$ . However, the mass increase due to the expansion of space-time should be only seen as an increase in cosmological potential energy  $U_c$  of the mass but not the change in rest mass. The rest mass remains constant. This is analogous to the relativistic mass of an object moving with some velocity with constant rest mass.

$$|U_c| = m_z z c^2$$

Changes in mass can be observed by converting  $m_0$  and  $m_z$  to energy, which is the observed energy difference in the cosmological redshift. Therefore, the energy of a photon does not increase with the space-time expansion, whereas the potential energy of a mass increases with the expansion, which perfectly explains the observed energy difference in the cosmological redshift. Similarly, for gravitational redshift, the change in mass is manifested as the gravitational potential energy ( $U$ ). The rest mass remains constant. Similar to cosmological redshift, the change in mass in gravitational redshift is also associated with the change in physical constants  $G$ ,  $h$ ,  $\varepsilon_0$ , and  $k_B$ .

Therefore, the energy of a photon does not change when moving against gravity, but the gravitational redshift  $z$ , which is the decrease in the frequency of the photon, is due to the increase in the Planck constant  $[E=h(1+z)\nu]$ . This proves that photons do not curve space-time by themselves but take the path of any curved space. As photons do not curve the space-time in this theory, the annihilation of matter would convert the cosmological potential energy, gravitational potential energy, and the rest mass to photons  $[E=m(1+z)c^2=h(1+z)\nu]$ .

Therefore, the gravitational potential energy  $U$  of mass  $m$  in the gravitational field of mass  $M$  is as follows:

$$U = -(m(1+z) - m)c^2 = -mzc^2, \text{ where } m(1+z) \text{ is the mass at infinity} \quad (14)$$

$$U = -mzc^2 \quad (15)$$

$$\frac{U}{m} = \phi = -zc^2 \quad (16)$$

where  $z$  is the gravitational redshift and  $\phi$  is the gravitational potential.

This theory predicts that the annihilation of matter should release some energy in the form of gravitational waves. However, it is not actively considered in this theory.

The change in mass with velocity, which is the relativistic mass, is also associated with changes in the physical constants  $G$ ,  $h$ ,  $\epsilon_0$ , and  $k_B$ . A person moving along with the mass will not observe the change in mass but observe the change in the above physical constants.

However, a stationary observer with respect to the moving mass will observe the increase in mass with no change in the physical constants. Therefore, both observers will see the same thing in two different ways, enabling the moving observer to know that the inertial frame is moving, although the laws of physics are the same in both frames.

For example, in this theory, a mass  $M$  moving with enough relativistic velocity to satisfy the Schwarzschild radius  $R$  can become a black hole. A person moving along with the mass will see it becoming a black hole too due to the increase in gravitational constant instead of relativistic mass as the mass remains constant in the moving inertial frame. Therefore, both stationary and moving observers will see the mass  $M$  becoming a black hole in two different ways.

From the stationary observer's perspective:  $R = \frac{2G(M\gamma)}{c^2}$

From the moving observer's perspective:  $R = \frac{2(G\gamma)M}{c^2}$ , where  $\gamma$  is the Lorentz factor.

As the gravitational constant  $G$  decreases with the expansion of the universe, this theory predicts that the radius of a black hole should proportionally decrease with constant mass  $M$ , ignoring accretion or evaporation of matter if any.

**Table II (Change in the values of the physical constants)****(17)**

Gravitational (From infinity)	$G(1+z)$	$\frac{h}{(1+z)}$	$\varepsilon_0(1+z)$	$\frac{k_B}{(1+z)}$
Transverse relativistic	$G(1+z)$ or $G\gamma$	$\frac{h}{(1+z)}$ or $\frac{h}{\gamma}$	$\varepsilon_0(1+z)$ or $\varepsilon_0\gamma$	$\frac{k_B}{(1+z)}$ or $\frac{k_B}{\gamma}$
Cosmological (From the big repulsion)	$\frac{G}{(1+z)}$	$h(1+z)$	$\frac{\varepsilon_0}{(1+z)}$	$k_B(1+z)$

where  $z$  is the gravitational, relativistic, and cosmological redshift, respectively.

For example, the Planck constant due to the cosmological redshift after 10 years from now would be  $h(1+\delta z)$  from (10). where  $\delta z$  is the change in the cosmological redshift value after 10 years

from now.  $(1+\delta z) = \frac{H_0}{H_1}$  using (4). where  $H_1$  is the Hubble constant after 10 years from now.

$\delta z$  values:

After 10 years:  $7.9463103517 \times 10^{-10}$

After 50 years:  $3.9731551758 \times 10^{-9}$

After 100 years:  $7.9463103517 \times 10^{-9}$

The accuracy of the values given below in Table III depends on the accuracy of the current values of  $h$ ,  $G$ ,  $\varepsilon_0$ ,  $k_B$ , and  $H_0$ .

**Table III (Change in physical constants due to cosmological/space-time expansion)****(18)**

Years	$\frac{G}{(1+\delta z)}$	$h(1+\delta z)$	$\frac{\varepsilon_0}{(1+\delta z)}$	$k_B(1+\delta z)$
0	$6.67430(15) \times 10^{-11}$	$6.62607015 \times 10^{-34}$	$8.8541878128(13) \times 10^{-12}$	$1.380649 \times 10^{-23}$
+10	$6.67429999 \times 10^{-11}$	$6.62607015 \times 10^{-34}$	$8.85418780 \times 10^{-12}$	$1.38064900 \times 10^{-23}$
+50	$6.67429997 \times 10^{-11}$	$6.62607017 \times 10^{-34}$	$8.85418777 \times 10^{-12}$	$1.38064900 \times 10^{-23}$
+100	$6.67429994 \times 10^{-11}$	$6.62607020 \times 10^{-34}$	$8.85418774 \times 10^{-12}$	$1.38064901 \times 10^{-23}$

Therefore, this theory provides falsifiable predictions by predicting the change in Planck's constant, the gravitational constant, the permittivity of free space, and the Boltzmann constant due to space-time expansion. Similar changes in the values of the above physical constants can be observed and calculated in the gravitational field and the relativistic frames using the factor  $(1+z)$ , where  $z$  is the gravitational redshift and the transverse relativistic redshift, respectively.

However, changes in the gravitational constant  $G$  in the distant past can be difficult to observe using the orbital frequencies of objects such as pulsars. Over a ten-year period, the value of the gravitational constant  $G$  would decrease to  $6.67429999 \times 10^{-11} \text{ N.m}^2.\text{kg}^{-2}$  (18) from the local  $G$  value ( $6.67430(15) \times 10^{-11} \text{ N.m}^2.\text{kg}$ ) at the orbiting object. As the decreased value of  $G$  is within the range of standard deviation ( $\pm 0.00015 \times 10^{-11}$ ) and/or combined with experimental errors and/or not having enough precision in the experiments to find the orbital periods could return a null result and hence difficult to observe the change in  $G$  in the distant past using the orbital frequencies. However, if the precision of the current value of  $G$  is sufficiently increased then this theory predicts that a change in  $G$  can be easily observed.

In the table below, we can see the similarity between the transverse relativistic mass increase, the gravitational mass increase, and the cosmological mass increase. In all three cases, mass increases by a factor of  $(1+z)$ . Here,  $m_0$  is the observed mass and  $m_z$  is the original mass at the point of the emission of the photon. The increase in cosmological and gravitational mass should be understood as an increase in their respective potential energies rather than an increase in their rest masses, which always remains constant.

**Table IV (Mass increase formulas)**

Transverse relativistic mass increase	Gravitational mass increase	Cosmological mass increase
$m_0 = m_z(1+z) = m_z \frac{1}{\sqrt{1-\frac{v^2}{c^2}}}$	$m_0 = m_z(1+z) = m_z \frac{1}{\sqrt{1-\frac{2GM}{Rc^2}}}$	$m_0 = m_z(1+z) = m_z \frac{1}{1-\frac{H_0 D}{c}}$

## 5. Acceleration of the universe due to electrostatic repulsion

In this theory of quantum dark energy in a seven-dimensional universe, the expansion of space-time is due to electrostatic repulsion between the Planck charges in the Planck volumes of the seven-dimensional space.

Acceleration at the beginning of the big repulsion is  $a = cH_0(1+z) = \frac{c}{t_p}$  from (2) and (3), which

can be reformulated as  $\sqrt{\rho_{t_p} G_{t_p}}$ , where  $\rho_{t_p}$  is the Planck energy density and  $G_{t_p}$  is the gravitational constant when the age of the universe is  $t_p$ . As acceleration  $a$  is directly related to the Planck energy density  $\rho$ , it proves the proposed model of the universe having Planck charges in Planck volumes, which exactly gives the Planck energy density  $\rho_{t_p}$  in the hidden three-dimensional space of the

seven-dimensional universe. Therefore,  $\frac{c}{t_p} = \sqrt{\rho_{t_p} G_{t_p}}$ , which can be generalized as follows.

$$a^2 = \rho G \quad (19)$$

Therefore, the product of the net electrostatic energy density of the universe in the hidden three dimensions and the gravitational constant is equal to the square of the acceleration of the universe.

$\rho$  = Electrostatic energy density of the universe responsible for the acceleration

$G$  = Gravitational constant

$a$  = Current acceleration of the universe, which is  $cH_0$

Planck energy  $E_p = \frac{1}{4\pi\epsilon_0} \frac{e^2}{\alpha(l_p)} = \sqrt{\frac{\hbar c^5}{G}}$ , where  $l_p$  is the Planck length. When the age of the universe is Planck time  $t_p$ , the energy density of the hidden three-dimensional space  $\rho_{t_p}$  of the universe is  $\frac{E_{p(t_p)}}{(l_p)^3}$ .

$$\rho_{t_p} = \frac{1}{4\pi\epsilon_{t_p}} \frac{e^2}{\alpha(l_p)^4} = \frac{\sqrt{\frac{\hbar_{t_p} c^5}{G_{t_p}}}}{(l_p)^3} \quad (20)$$

$$\rho_{t_p} G_{t_p} = \frac{1}{4\pi\epsilon_{t_p}} \frac{e^2}{\alpha(l_p)^4} G_{t_p} = \left(\frac{c}{t_p}\right)^2 = a^2$$

$$G_{t_p} = \frac{G}{H_0 t_p} \text{ from (11)}$$

$$\rho_{t_p} G_{t_p} = \frac{\rho_{t_p} G}{H_0 t_p} = \left(\frac{c}{t_p}\right)^2$$

$$\rho_{t_p} = \frac{H_0 c^2}{G t_p} \quad (21)$$

Therefore,  $\rho_{t_p} = 6.29 \times 10^{52} \text{ J/m}^3$ , which can be generalized as follows.

$$\rho G(1+z) = (cH_0(1+z))^2$$

$$\rho = \frac{(cH_0)^2(1+z)}{G} = \frac{(cH_0)^2 \left(\frac{c}{c-H_0 D}\right)}{G} \text{ from (4)}$$

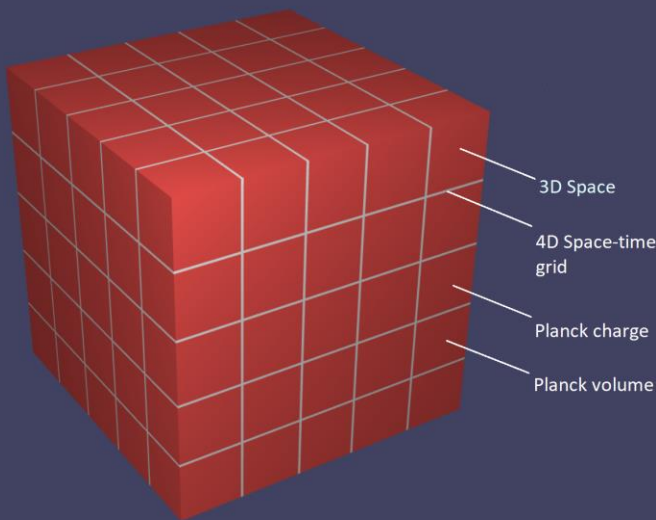
The current electrostatic energy density  $\rho$  responsible for the acceleration of the universe can be found by setting  $D = 0$  or  $z = 0$ .

$$\rho = \rho_{t_p} H_0 t_p = \frac{\rho_{t_p}}{1+z} = \frac{(cH_0)^2}{G} = 8.54 \times 10^{-9} \text{ J/m}^3, \text{ which resolves to (19).}$$

As the universe accelerates due to electrostatic repulsion, the electrostatic potential energy in the hidden three-dimensional space is gradually transferred to the four-dimensional space-time grid and stored as potential energy. As the elastic space-time grid resists expansion, the rate of acceleration gradually decreases to follow  $a = cH_0$ , which will asymptote to zero. As the total energy is conserved, the potential energy density of the four-dimensional space-time grid  $\rho_g$  is below.

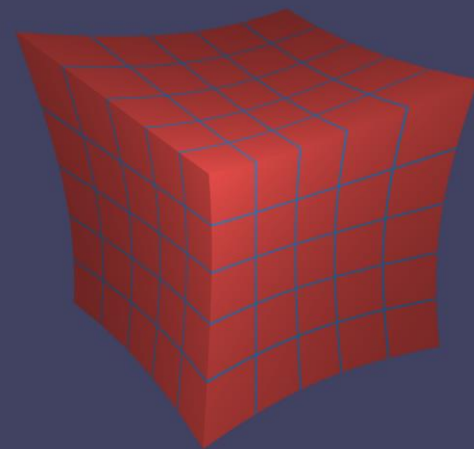
$$|\rho_g| = \rho_{t_p} - \rho \quad (22)$$

As  $\rho$  decreases with time and  $\rho_{t_p}$  is constant, the potential energy density of the four-dimensional space-time grid  $\rho_g$  will continue to increase until infinity owing to the declining permittivity of free space  $\epsilon_0$ .



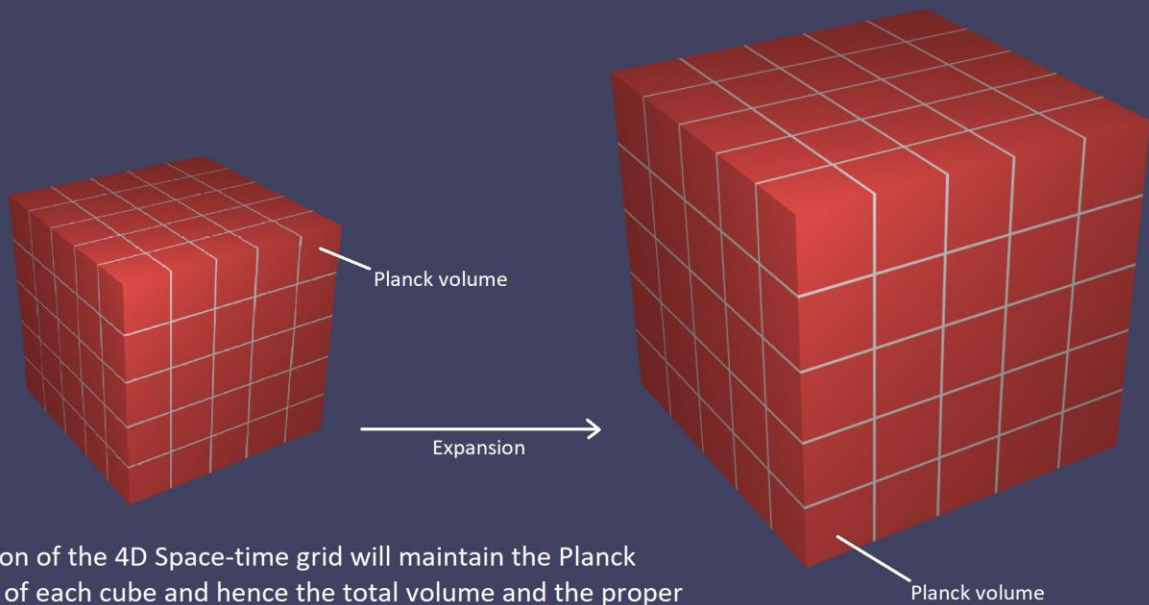
Total number of dimensions = 3D + 4D = 7D

Fig 1



Matter in the 4D Space-time grid curves the space-time around it due to the electrostatic repulsion outside the matter being greater than inside as mass also increases/expands along with space-time.

Fig 2



Expansion of the 4D Space-time grid will maintain the Planck volume of each cube and hence the total volume and the proper distance remains constant. Expansion is similar to the stretching of an elastic ruler where the length remains constant.

Fig 3



## 6. The temperature of the universe

The internal energy  $U$  of black-body photon gas is given by  $U = \left( \frac{8\pi^5 k^4}{15h^3 c^3} \right) VT^4$ .

where  $k$  = Boltzmann constant,  $h$  = Planck constant,  $c$  = Speed of light,  $V$  = Volume, and  $T$  = Temperature. (Leff, 2002)

As the number of photons  $N = \left( \frac{16\pi k^3 \zeta(3)}{h^3 c^3} \right) VT^3$ , we can substitute  $V = \left( \frac{Nh^3 c^3}{16\pi k^3 \zeta(3)} \right) \frac{1}{T^3}$  in  $U$ .

As  $k = k_{t_p} (1+z)$ , the maximum possible temperature of the universe when the age of the universe was  $t_p$  is given below by upholding the law of conservation of energy and considering CMB a black-body radiation.

As the energy  $U$  and the number of photons  $N$  remain constant with the expansion of the universe,

$$U = \frac{\pi^4 N k T}{30 \zeta(3)} = \frac{\pi^4 N \frac{k}{(1+z)} T_{t_p}}{30 \zeta(3)} ; T_{t_p} = T(1+z), \text{ which can be generalized as follows.}$$

$$T = T_0(1+z) \quad (23)$$

As  $T_0 = 2.725$  K and the maximum possible cosmological redshift  $z$  is  $7.4 \times 10^{60}$  from (3),  $T_{t_p} = 2 \times 10^{61}$  K. As the  $T_{t_p}$  value is greater than the Planck temperature, the maximum possible temperature of the universe is the Planck temperature, which is  $1.416784 \times 10^{32}$  K. However, this does not mean that this was the temperature at the time of the big repulsion but only the maximum possible temperature. As the CMB was emitted after the initial  $t_p$  of the big repulsion, the original temperature of the CMB when it was first emitted should be less than the Planck temperature. The future temperature of CMB radiation can also be calculated using the above formula. For example, the CMB temperature after  $10^7$  years would be 2.722 K.

As the entropy of a photon gas  $S = \frac{4U}{3T}$ , it follows that  $S = \frac{S_0}{(1+z)}$  (23), where  $S_0$  is the current entropy of CMB and  $S$  is the entropy at cosmological redshift  $z$ .

As the interaction of the electrostatic expansion energy with the quantum vacuum fluctuations to create matter at the time of big repulsion will be the same throughout space-time, the created primordial elementary particles, atoms, and their attributes, such as temperature and density, should be uniform throughout space-time without having the particles interact with one another or without being causally connected, which eliminates the need for cosmic inflation and solves the horizon problem, the flatness problem, and explains the uniformity of CMB radiation. However, minor fluctuations in the density of the created particles due to the randomness of the quantum vacuum fluctuations and the subsequent concentration of the matter due to gravity could explain the temperature anisotropy. However, the angular power spectrum of the CMB still needs to be explained, possibly through BAO (Baryon acoustic oscillations) and a high gravitational constant without dark matter.

## 7. The absolute frame of reference

The four-dimensional space-time grid proposed in this theory acts as an absolute inertial frame. As the physical constants ( $G$ ,  $h$ ,  $\epsilon_0$ , and  $k_B$ ) change with the change in gravitational potential and velocity, any frame of reference that is not absolute can be easily identified by knowing the values of the physical constants. Physical constants change in the gravitational field and due to velocity according to the formulas given in Table II (17). Therefore, in the gravitational field, the absolute reference frame is the one at an infinite distance from the mass generating the gravitational field, and the values of the physical constants change as we move from infinity toward the mass. For example, the gravitational constant would be higher near a mass than away from the mass, and the Planck constant would be lower near a mass than away from the mass.

As the Hubble constant decreases due to the space-time expansion, the values of the physical constants of the absolute inertial frame need to be updated to keep up with the expansion. Changes in the physical constants of the absolute inertial frame due to cosmological expansion are given in Table III (18). Additionally, the universal time in this theory is the time in the absolute inertial frame, which is the Hubble time.

Similarly, a frame moving with a velocity will have a higher gravitational constant and a lower Planck constant compared to the absolute inertial frame. However, a source of light emitting green photons will still emit green photons in both frames with different Planck constants due to the time-dilation in the moving frame.

For example, in this theory, acceleration due to Newtonian gravity is covariant in all inertial frames. Consider two masses each of mass  $m$  separated by a distance  $R$  moving with velocity  $v$  in an inertial frame. The time-dilation of the moving masses in the stationary reference frame compensates for the time-dilation of the observer moving along with the masses. Therefore, the observers in both the stationary and moving reference frames will observe the same formula for acceleration ( $a$ ) due to gravity in two different ways, one with relativistic mass  $\gamma m$  and the other with relativistic gravitational constant  $\gamma G$ , as shown in Table V, except for the time-dilation effect observed by the stationary observer, which makes the acceleration due to gravity covariant similar to electromagnetism. Gravitomagnetic effects are covered in section 8.5.

**Table V (Covariant acceleration due to Newtonian gravity)**

	WRT stationary reference frame	WRT moving reference frame
Masses aligned perpendicular to the direction of motion	$(\gamma m)a = F = \frac{G(\gamma m)(\gamma m)}{R^2} ; a = \frac{G(\gamma m)}{R^2}$	$a = \frac{(G\gamma)m}{R^2}$
Masses aligned parallel to the direction of motion	$\gamma^3 ma = F = \frac{G(\gamma m)(\gamma m)}{\left(\frac{R}{\gamma}\right)^2} ; a = \frac{G(\gamma m)}{R^2}$	$a = \frac{(G\gamma)m}{R^2}$

## 8. The relativistic Newtonian theory of gravity

In this theory, the presence of mass increases the permittivity of the four-dimensional space-time grid and hence reduces the electrostatic repulsion within the grid enveloped by the mass compared to the electrostatic repulsion outside the mass. As the contraction of the space-time grid within mass will be greater than the contraction outside due to the permittivity difference and the net compressing force on the matter from the surrounding Planck charges, space-time becomes naturally curved around any mass and hence the gravitational force. Therefore, in this theory, the gravitational force is considered to be a pushing type force rather than a pulling type of the regular non-relativistic Newtonian gravity and hence considered a real force rather than a fictitious force of the general theory of relativity.

The change in gravitational constant  $G$  in the gravitational field is similar to the change in the gravitational constant due to cosmological space-time expansion. Therefore, we can take Newton's law of gravitation and introduce variable  $G$  and the gravitational length-contraction to derive a relativistic law. Here, the change in gravitational potential energy is not included as part of the mass, as it does not cause additional gravitational force; only the rest and relativistic masses cause gravitational force in this theory.

$$F = \frac{G(1+z)Mm}{\left(\frac{R}{1+z}\right)^2} = (1+z)^3 \frac{GMm}{R^2} \text{ using (17), where } G \text{ is the gravitational constant at infinity or the}$$

absolute reference frame. As the space-time grid shrinks as we move from infinity toward the mass, the length-contraction as observed from infinity or the absolute reference frame is included for  $R$  to factor in the local curvature of space-time.

$$\frac{dU}{dR} = F = (1+z)^3 \frac{GMm}{R^2} = \left(1 - \frac{U}{mc^2}\right)^3 \frac{GMm}{R^2} \text{ using (15)}$$

$$\int_0^U \left(\frac{1}{mc^2 - U}\right)^3 dU = \frac{GMm}{(mc^2)^3} \int_{\infty}^R \frac{1}{R^2} dR$$

$$\frac{1}{2(mc^2 - U)^2} - \frac{1}{2(mc^2)^2} = -\frac{GMm}{(mc^2)^3} \frac{1}{R}, \text{ divide the denominator by } (mc^2)^2 \text{ on both sides.}$$

$$\frac{1}{2\left(1 - \frac{U}{mc^2}\right)^2} - \frac{1}{2} = -\frac{GMm}{mc^2} \frac{1}{R}$$

$$\frac{1}{2(1+z)^2} - \frac{1}{2} = -\frac{GMm}{mc^2} \frac{1}{R} \text{ using (15); after solving, we obtain the following equation.}$$

$$1+z = \sqrt{\frac{1}{1 - \frac{2GM}{Rc^2}}} \quad (24)$$

As the above gravitational redshift matches the redshift from the Schwarzschild solution of the Einstein field equations without using the weak field approximation, it validates this theory, the concept of variable physical constants, and proves the existence of an absolute reference frame. Additionally, the formulas below for the gravitational force and the gravitational potential energy are not approximations but complete solutions that work in both strong and weak gravitational fields for nonrotating stationary spherical masses without using the weak field approximation as in the general theory of relativity. If there is kinetic energy of the particles within the masses (e.g., kinetic energy due to temperature and/or angular momentum), the respective relativistic masses should be used instead of the rest masses  $M$  and  $m$  to obtain accurate results.

Gravitational redshift  $z$  becomes infinity at  $R = \frac{2GM}{c^2}$  in (24), which gives the Schwarzschild radius without using a weak field approximation.

Relativistic gravitational force  $F$  for  $M \geq m$ :

$$F = (1+z)^3 \frac{GMm}{R^2} = \left( \sqrt{\frac{1}{1 - \frac{2GM}{Rc^2}}} \right)^3 \frac{GMm}{R^2} \quad (25)$$

In Table V, we can see that the acceleration due to gravity  $a$  is the same in both the stationary and moving inertial frames except for the time-dilation effect in the stationary frame. Relativistic acceleration due to gravity in both frames can be obtained by including the  $(1+z)^3$  factor (25).

$\frac{d\phi}{dR} = a = (1+z)^3 \frac{G\gamma m}{R^2} = \left(1 - \frac{\phi}{c^2}\right)^3 \frac{G\gamma m}{R^2}$  using (16);  $\int_0^U \left(\frac{1}{c^2 - \phi}\right)^3 d\phi = \frac{G\gamma m}{(c^2)^3} \int_{\infty}^R \frac{1}{R^2} dR$ ; after solving, we obtain the following equation.

$$1+z = \sqrt{\frac{1}{1 - \frac{2G\gamma m}{Rc^2}}}$$

Relativistic gravitational potential energy  $U$  for  $M \geq m$ :

$$U = -mzc^2 = -m \left( \sqrt{\frac{1}{1 - \frac{2GM}{Rc^2}}} - 1 \right) c^2 \quad (26)$$

Equating the above relativistic gravitational potential energy  $U$  with the relativistic kinetic energy  $(m(\gamma-1)c^2)$  with escape velocity  $c$  directly gives the Schwarzschild radius  $R_s = \frac{2GM}{c^2}$ .

For weak gravitational fields, (26) reduces to the regular Newtonian gravitational potential energy and hence validates this theory.

$$U = -m \left( \left( 1 - \frac{2GM}{Rc^2} \right)^{\frac{1}{2}} - 1 \right) c^2 \approx -m \left( 1 + \frac{GM}{Rc^2} - 1 \right) c^2 = -\frac{GMm}{R}$$

Relativistic gravitational potential  $\phi$  for  $M \geq m$ :

$$\phi = \frac{U}{m} = -zc^2 = - \left( \sqrt{\frac{1}{1 - \frac{2GM}{Rc^2}}} - 1 \right) c^2 \quad (27)$$

The maximum energy (rest + potential)  $E_{\max}$  that can be gained by a mass  $m$  at radius  $R$  from mass  $M$  is below.

$$E_{\max} = m(1+z)c^2 = m \left( \sqrt{\frac{1}{1 - \frac{2GM}{Rc^2}}} \right) c^2 \text{ from (14)}$$

As the energy is conserved in the freefall motion of mass  $m$  in the gravitational field, the velocity of mass  $m$  can be derived from the maximum energy  $E_{\max}$ . Through dimensional analysis, the

energy at any point during the freefall must be of the form  $m \left( \sqrt{\frac{1}{1 - \frac{v^2}{c^2}}} \right) c^2$ . Therefore,  $v^2 = \frac{2GM}{R}$ .

$v = \sqrt{\frac{2GM}{R}}$ , which is the velocity of mass  $m$  in a freefall motion from the point of maximum potential energy or infinity. This is also equal to the escape velocity  $v_e$ . As the freefall motion is inertial, the Lorentz factor  $\gamma$  is directly realized from energy conservation, which is  $\sqrt{\frac{1}{1 - \frac{v^2}{c^2}}}$ , and

hence, this theory upholds the special theory of relativity in all inertial frames.

The length-contraction, time-dilation, and mass change in the gravitational field are given by

$L = \frac{L_{\infty}}{(1+z)}$ ,  $T = \frac{T_{\infty}}{(1+z)}$ , and  $m = \frac{m_{\infty}}{(1+z)}$ , where  $L_{\infty}$ ,  $T_{\infty}$ , and  $m_{\infty}$  are the absolute length, absolute time, and maximum mass, respectively, at infinity or the absolute reference frame (28)

As the radial length  $L$ ,  $T$ , and  $m$  becomes zero at Schwarzschild's radius, space-time terminates at the event horizon of a black hole. Black hole singularities do not exist in this model of the universe as the mass becomes zero and is converted to the equivalent pure informational entropy at the event horizon. Space-time terminates at the event horizon and cannot be extended beyond the event horizon, as in the general theory of relativity, as the radial length, time and tangible mass become zero at the event horizon and hence discards the Riemannian geometry of space-time.

Additionally, charged black holes do not exist in this model of the universe as the permittivity of free space becomes infinity (from (17)) at the event horizon. Therefore, stationary black holes can only have two properties, namely, informational (entropic) mass and angular momentum, and hence only partially satisfy the no-hair theorem.

### 8.1. Photon sphere

As the photon takes the curved path around the black hole due to the space-time curvature, matching the fictitious acceleration due to gravity for the photon with the centripetal acceleration gives the radius of the photon sphere for a nonspinning black hole, as predicted by the general theory of relativity. Here, the forces are considered to be fictitious, as the photon does not curve space-time by itself but only takes the path of curved space-time due to optical refraction. The path taken by a photon or a particle due to refraction in curved space-time is equivalent to the geodesic path in the general theory of relativity. Here, the refraction of the photon due to the curvature is considered equivalent to the centripetal acceleration. Therefore, the fictitious centripetal acceleration of the photon due to refraction in the curved space-time around mass  $M$  is  $\frac{c^2}{\left(\frac{R}{1+z}\right)}$ .

Here, the factor  $(1+z)$  accounts for the local curvature of space-time.

$$(1+z)^3 \frac{GM}{R^2} = \frac{c^2}{\left(\frac{R}{1+z}\right)} \text{ using (25), and solving using (24), we obtain the radius of the photon}$$

sphere.

$$\left( \sqrt{1 - \frac{2GM}{Rc^2}} \right)^2 \frac{GM}{Rc^2} = 1 \quad (29)$$

$$R = \frac{3GM}{c^2} \quad (30)$$

The probability wave function ( $\psi$ ) of any matter particle ( $m \ll M$ ) also takes the path of the curved space-time and is refracted like a photon but curves the space-time unlike a photon and causes a wobble in mass  $M$ ; hence, the forces are considered real. A photon does not cause a wobble in mass  $M$ , as it is massless. As the matter particles move with velocities less than  $c$ , the centripetal acceleration experienced by a matter particle due to the refraction of its wave function is  $\frac{v^2}{\left(\frac{R}{1+z}\right)}$ .

The innermost stable circular orbit (ISCO) does not apply to this theory, as the relativistic gravitational potential ( $\phi$ ) always changes for any radius greater than the Schwarzschild radius.

## 8.2. Gravitational lensing

As the angle of deflection  $\theta$  for the photon sphere is  $\pi$  radians owing to the symmetry of the sphere, as shown in Figure 4, we can calculate the angle of deflection for a nonspinning spherical mass for any radius  $R$  within the relativistic Newtonian regime by replacing 1 in (29) with  $\text{Tan}\left(\frac{\pi}{4}\right)$ , which is the ratio of the gravitational acceleration to the centripetal acceleration of the photon at its closest approach to mass  $M$ . As this ratio should be proportional to the angle of deflection  $\theta$ , we can generalize this equation by replacing  $\text{Tan}\left(\frac{\pi}{4}\right)$  with  $\text{Tan}\left(\frac{\theta}{4}\right)$  to find the angle of deflection for any radius  $R$ .

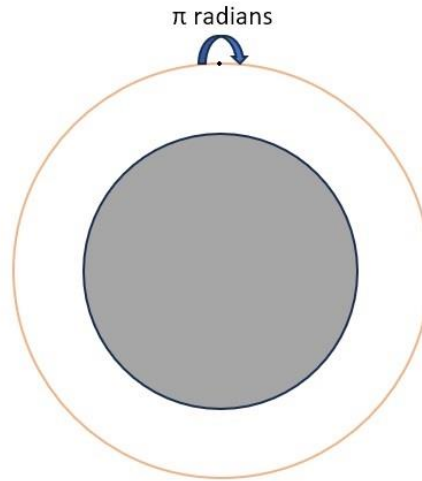


Fig 4

$$\left( \sqrt{1 - \frac{2GM}{Rc^2}} \right)^2 \frac{GM}{Rc^2} = \text{Tan}\left(\frac{\pi}{4}\right)$$

$$\theta = 4 \text{Tan}^{-1} \left( (1+z)^2 \frac{GM}{Rc^2} \right) \quad (31)$$

For  $\theta = 2\pi$ , we obtain  $R = \frac{2GM}{c^2}$ , which is the Schwarzschild radius and hence validates the above equation (31). Therefore, light bends onto itself at the event horizon of a black hole and does not travel beyond it as in the general theory of relativity and hence provides another proof that space-time ends at the event horizon and the nonexistence of singularity at the center of a black hole. For weak gravitational fields, the angle of deflection in (31) reduces to  $\theta = \frac{4GM}{Rc^2}$ , which is the same as the deflection predicted by the general theory of relativity. Therefore, equation (31) predicts the angle of deflection in a closed form that works in both strong and weak gravitational fields better than the general theory of relativity without using any approximations or higher-order terms.



### 8.3. Perihelion precession of Mercury

From Kepler's second law,  $dA = \frac{1}{2} r^2 d\theta$ , where  $dA$  is the change in the area swept out by the orbiting mass,  $d\theta$  is the change in angle and  $r$  is the radius.

Let  $dA'$  and  $d\theta'$  be the local change in area and angle, respectively. Divide the above equation on both sides by the local time  $dt'$  and apply length-contraction (28) to factor in the local curvature of space-time.

$$\frac{dA'}{dt'} = \frac{1}{2} \left( \frac{r}{1+z} \right)^2 \frac{d\theta'}{dt'}$$

Apply time-dilation using (28), where  $dt'$  is the local time and  $dt$  is the local time without the space-time curvature.

$$\frac{dA'}{dt'} = \frac{1}{2} \left( \frac{r}{1+z} \right)^2 \frac{d\theta'}{dt(1+z)} = \frac{1}{2} r^2 \frac{d\theta'}{dt} \frac{1}{(1+z)^3}$$

As the angular momentum is conserved, the above equation can be compared with the nonrelativistic equation without the space-time curvature to obtain the following equation.

$$d\theta' = d\theta(1+z)^3 = d\theta \left( 1 - \frac{2GM}{Rc^2} \right)^{\frac{3}{2}} \approx d\theta \left( 1 + \frac{3GM}{Rc^2} \right) \text{ by ignoring the higher-order terms.}$$

The polar equation of the ellipse is  $R = \frac{a(1-\varepsilon^2)}{1-\varepsilon \cos \theta}$ .

$$\theta' = \int_0^{2\pi} \left( 1 + \left( \frac{3GM}{c^2} \frac{1-\varepsilon \cos \theta}{a(1-\varepsilon^2)} \right) \right) d\theta$$

$\theta' = 2\pi + \frac{6\pi GM}{c^2 a(1-\varepsilon^2)}$ , the second term gives the precession angle  $\Delta\phi$  of the perihelion per revolution. where  $M$  is the mass of the sun,  $a$  is the semimajor axis of Mercury and  $\varepsilon$  is the orbital eccentricity.

$$\Delta\phi = \frac{6\pi GM}{c^2 a(1-\varepsilon^2)}$$

Solving the above, we obtain a perihelion precession of 43"/century, which is the same as that predicted by the general theory of relativity (Park et al., 2017).

## 8.4. Shapiro time delay

We can calculate the gravitational time delay of light passing by a mass such as the sun. As the gravitational length-contraction and gravitational time-dilation go together in this theory, we can calculate the effective speed of light as observed by an observer on Earth as the light from a distant object passes near the sun and reaches the Earth.

The speed of light is constant at every point in space in the gravitational field, as the length-contraction is compensated by the time-dilation. Therefore,  $\frac{\frac{\text{Distance}}{(1+z)}}{\frac{\text{Time}}{(1+z)}} = \frac{\text{Distance}}{\text{Time}} = c$  (the speed

of light) remains constant. However, a stationary observer would see a change in the speed of light w.r.t. his reference frame as the light goes through the gravitational field. Here, the time expansion is applied instead of contraction, as the time delay is observed from the observer's frame of reference whose local time is faster than the time near the sun. Therefore, the length-contraction and time expansion do not cancel out in the observer's reference frame and hence produce a time delay according to this theory, which is called Shapiro time delay.

The effective speed of light as experienced by the stationary observer is  $\frac{\frac{\text{Distance}}{(1+z)}}{\text{Time}(1+z)} = \frac{c}{(1+z)^2}$ .

$\frac{dx}{dt} = \frac{c}{(1+z)^2} = c \left( 1 - \frac{2GM}{Rc^2} \right)$  using (24), where  $R$  is the distance between the sun and the traveling photon. Neglecting the deflection of the light near the sun, the path of the light from  $A$  to  $B$  is a straight line.

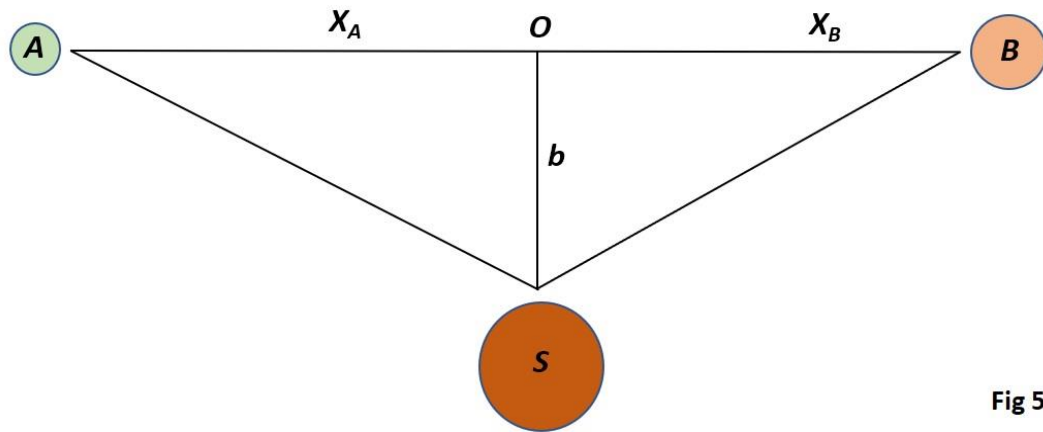


Fig 5

$dt = \frac{1}{c} \left( \frac{1}{1 - \frac{2GM}{Rc^2}} \right) dx \approx \frac{1}{c} \left( 1 + \frac{2GM}{\sqrt{x^2 + b^2}c^2} \right) dx$  by ignoring the higher-order terms, where  $b$  is the impact parameter. Integrating on the left side from  $T_A$  to  $T_B$  and on the right side from  $X_A$  to  $X_B$ .

$$\int_{T_A}^{T_B} dt = \int_{X_A}^{X_B} \frac{1}{c} \left( 1 + \frac{2GM}{\sqrt{x^2 + b^2} c^2} \right) dx$$

$$T_B - T_A = \frac{X_B - X_A}{c} + \frac{2GM}{c^3} \ln \left( \frac{X_B + \sqrt{X_B^2 + b^2}}{X_A + \sqrt{X_A^2 + b^2}} \right)$$

The second term gives the additional one-way time delay  $\Delta t$  of the light coming from point A as observed by an observer at point B, which better fits the experimental data than the Schwarzschild metric using Schwarzschild coordinates (Pössel, 2019, Page 8).

$$\Delta t = \frac{2GM}{c^3} \ln \left( \frac{X_B + \sqrt{X_B^2 + b^2}}{X_A + \sqrt{X_A^2 + b^2}} \right)$$

## 8.5. Relativistic Poisson and gravitoelectromagnetism (GEM) equations

Relativistic Poisson and Maxwell-like GEM equations can be generated to make them work in strong gravitational fields and relativistic speeds as well by applying the Laplace operator ( $\nabla^2$ ) on the relativistic gravitational potential. Relativistic GEM equations can explain the frame-dragging effect, orbital precession, geodetic effect, and gravitational waves. Both the relativistic GEM and Poisson equations together can form the equivalent of Einstein field equations for this theory with the scope for further generalization.

$$\text{As } z = \frac{-\phi}{c^2} \text{ (16), } (1+z) = \left( 1 - \frac{\phi}{c^2} \right).$$

### 8.5.1. Relativistic Poisson equations

The relativistic gravitational potential in Cartesian coordinates based on (27) is given below.

$$\phi = - \left( \frac{1}{\sqrt{1 - \frac{2GM}{(\sqrt{x^2 + y^2 + z^2}) c^2}}} - 1 \right) c^2$$

The relativistic Poisson equation outside the point mass can be generated by applying the Laplace operator to the above gravitational potential.

$$\begin{aligned} \nabla^2 \phi = & - \frac{3G^2 M^2 x^2 (1+z)^5}{(x^2 + y^2 + z^2)^3 c^2} - \frac{3G^2 M^2 y^2 (1+z)^5}{(x^2 + y^2 + z^2)^3 c^2} - \frac{3G^2 M^2 z^2 (1+z)^5}{(x^2 + y^2 + z^2)^3 c^2} \\ & - \frac{3GMx^2 (1+z)^3}{(x^2 + y^2 + z^2)^{\frac{5}{2}}} - \frac{3GM y^2 (1+z)^3}{(x^2 + y^2 + z^2)^{\frac{5}{2}}} - \frac{3GM z^2 (1+z)^3}{(x^2 + y^2 + z^2)^{\frac{5}{2}}} + \frac{3GM (1+z)^3}{(x^2 + y^2 + z^2)^{\frac{3}{2}}} \end{aligned}$$

$\nabla^2\phi = -2\pi G\rho\left[(1+z)^5 - (1+z)^3\right]$ , where  $\rho = \frac{M}{\frac{4}{3}\pi r^3}$  is the relativistic mass density of the point mass.

$$\text{Outside the point mass: } \nabla^2\phi = -2\pi G\rho\left[\left(1-\frac{\phi}{c^2}\right)^5 - \left(1-\frac{\phi}{c^2}\right)^3\right] \approx 4\pi G\rho\frac{\phi}{c^2} \quad (32)$$

The relativistic Poisson equation at the point mass can be generated by considering only the positive terms and ignoring the negative terms in the Laplacian ( $\nabla^2\phi$ ) as the positive terms represent the positive divergence or the source of the gravitational flux.

$$\nabla^2\phi = \frac{3GM(1+z)^3}{(x^2 + y^2 + z^2)^{\frac{3}{2}}} = 4\pi G\rho(1+z)^3, \text{ where } \rho \text{ is the relativistic mass density of the point mass.}$$

$$\text{At the point mass: } \nabla^2\phi = 4\pi G\rho\left(1-\frac{\phi}{c^2}\right)^3 \quad (33)$$

For weak gravitational fields, where the gravitational redshift  $z$  is negligible, the above equations (32) and (33) reduce to Laplace and nonrelativistic Poisson equations, respectively.

$$\text{Outside the point mass: } \nabla^2\phi = -2\pi G\rho\left[(1+z)^5 - (1+z)^3\right] \approx -2\pi G\rho[1-1] = 0$$

$$\text{At the point mass: } \nabla^2\phi = 4\pi G\rho(1+z)^3 \approx 4\pi G\rho[1] = 4\pi G\rho$$

### 8.5.2. Relativistic GEM equations

Relativistic Maxwell-like GEM equations can be generated using (33) and (32) at the point mass/current and outside the point mass/current, respectively. where  $E_G$  is the gravitoelectric field,  $B_G$  is the gravitomagnetic field,  $\phi$  is the relativistic gravitational potential due to the gravitoelectric effect, and  $J$  is the relativistic mass current density.

At the point mass/current:

$$\begin{aligned} \nabla \cdot E_G &= -4\pi G\rho\left(1-\frac{\phi}{c^2}\right)^3 \\ \nabla \times B_G &= -\frac{4\pi G\left(1-\frac{\phi}{c^2}\right)^3}{c^2}J + \frac{1}{c^2}\frac{\partial E_G}{\partial t} \\ \nabla \cdot B_G &= 0 \\ \nabla \times E_G &= -\frac{\partial B_G}{\partial t} \end{aligned}$$

Outside the point mass/current:

$$\begin{aligned} \nabla \cdot E_G &= 2\pi G\rho\left[\left(1-\frac{\phi}{c^2}\right)^5 - \left(1-\frac{\phi}{c^2}\right)^3\right] \\ \nabla \times B_G &= \frac{2\pi G\left[\left(1-\frac{\phi}{c^2}\right)^5 - \left(1-\frac{\phi}{c^2}\right)^3\right]}{c^2}J + \frac{1}{c^2}\frac{\partial E_G}{\partial t} \\ \nabla \cdot B_G &= 0 \\ \nabla \times E_G &= -\frac{\partial B_G}{\partial t} \end{aligned}$$

## 8.6. Space-time interval ( $ds^2$ )

The invariant line element ( $ds^2$ ) in the Minkowski space-time in spherical coordinates is given by  $ds^2 = -c^2 dt^2 + dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$ . The presence of a gravitational field would dilate the relativistic proper time  $d\tau$  to  $\frac{d\tau}{(1+z)}$ . Therefore,  $d\tau_g = \frac{d\tau}{(1+z)}$ , where  $d\tau_g$  is the relativistic proper time with gravity included and  $z$  is the gravitational redshift. We can substitute  $d\tau = d\tau_g(1+z)$  in the line element ( $ds^2$ ) of the Minkowski space-time to find the line element with gravity ( $ds^2 = c^2 d\tau_g^2$ ) included.

$c^2 (d\tau_g(1+z))^2 = -c^2 dt^2 + dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$ , where  $dt$ ,  $dr$ , and  $r$  are the coordinate time and lengths respectively as measured from the absolute reference frame or the reference frame at infinity.

$$ds^2 = -\frac{c^2 dt^2}{(1+z)^2} + \frac{dr^2}{(1+z)^2} + \frac{r^2}{(1+z)^2} (d\theta^2 + \sin^2 \theta d\phi^2) \quad (34)$$

Applying the gravitational length-contraction and time-dilation (24) (16), we obtain the equivalent of the Schwarzschild metric for this theory.

Schwarzschild metric equivalent:

$$ds^2 = -c^2 dt^2 \left(1 - \frac{2GM}{rc^2}\right) + dr^2 \left(1 - \frac{2GM}{rc^2}\right) + r^2 \left(1 - \frac{2GM}{rc^2}\right) (d\theta^2 + \sin^2 \theta d\phi^2)$$

OR

$$ds^2 = -c^2 dt^2 \left(1 - \frac{\phi}{c^2}\right)^{-2} + dr^2 \left(1 - \frac{\phi}{c^2}\right)^{-2} + r^2 \left(1 - \frac{\phi}{c^2}\right)^{-2} (d\theta^2 + \sin^2 \theta d\phi^2)$$

Once the  $(1+z)$  factor  $\left(1 - \frac{\phi}{c^2}\right)$  is identified by solving the relativistic Poisson equation (33), it can be substituted in (34) to obtain the space-time interval equation for this theory, which is equivalent to the equation  $ds^2 = g_{\mu\nu} dx^\mu dx^\nu$  in the general theory of relativity (Einstein, 1922). For any other inertial/noninertial reference frame, the respective coordinate length, time, and gravitational redshift  $z$  up to the origin of the reference frame must be used in (34) to make the line element ( $ds^2$ ) invariant in all reference frames.

In this theory, space-time ends at the event horizon of a black hole as the time and length become zero and hence eliminates singularity, unlike in the general theory of relativity, where the length  $dr$  becomes infinity and custom coordinate systems are used to avoid infinities and to extend the space-time up to the singularity at the center of a black hole. Additionally, the slope of the light ray of the light cone is always equal to 1 (45 degrees) in this theory, even near the event horizon, unlike in the general theory of relativity, where the light cone gradually becomes narrow and the slope becomes infinity at the event horizon. Setting the line element  $ds^2=0$ , which is lightlike, and  $c=1$  in (34), ignoring  $d\theta$  and  $d\phi$ , we obtain the slope  $\frac{dt}{dr} = 1$ .

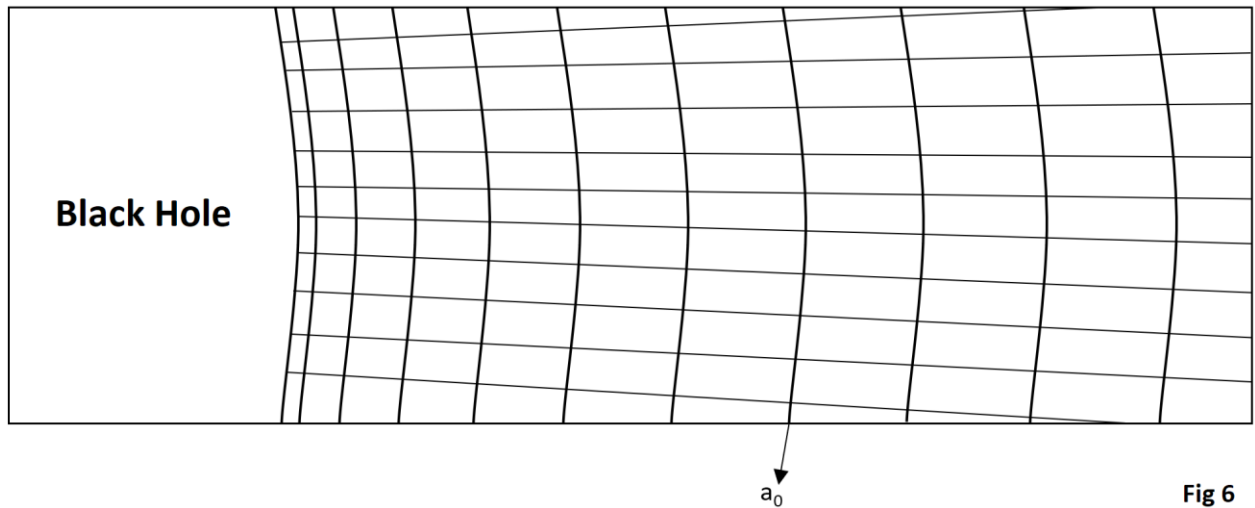
Applying the cosmological length-contraction and time-dilation (4), we obtain the equivalent of the FLRW metric of the flat space-time for this theory.

FLRW metric equivalent:

$$ds^2 = -c^2 dt^2 \left(1 - \frac{H_0 r}{c}\right)^2 + dr^2 \left(1 - \frac{H_0 r}{c}\right)^2 + r^2 \left(1 - \frac{H_0 r}{c}\right)^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

## 9. MOND (Modified Newtonian dynamics)

We can generate the relativistic MOND formula and the respective Poisson equation using the  $(1+z)$  factor similar to the relativistic Newtonian theory of gravity. However, the relativistic effect in the deep-MOND regime is not significant due to lower accelerations. This theory proposes that the MOND regime can only be present around a black hole but not around ordinary masses such as gas and stars. In the absence of black holes, the gravitational field around ordinary masses will always be Newtonian at any acceleration. Therefore, the cutoff acceleration for the MOND regime, which is  $a_0$  ( $\sim 1.2 \times 10^{-10} \text{ m/s}^2$ ) (Bekenstein, 2009), is only applicable for the area around a central black hole but not for ordinary masses. Each cell in Figure 6 of the curved space around the black hole represents the Planck volume with Planck charge. The radial length becomes zero at the event horizon but will continue to expand to follow the relativistic Newton's law of gravitation as we move away from the black hole. Additionally, the number of Planck volumes on the circumference remains constant at any radius from the event horizon.



To understand the reason for the transition of the Newtonian regime to the MOND regime at  $a_0$ , we can flatten the above-curved space, which is depicted in Figure 7. As the relative Planck volume cannot increase more than the maximum allowed by the expansion of the universe, the relative size of the Planck volume should remain constant after  $a_0$ , which is the case in Figure 7 below.

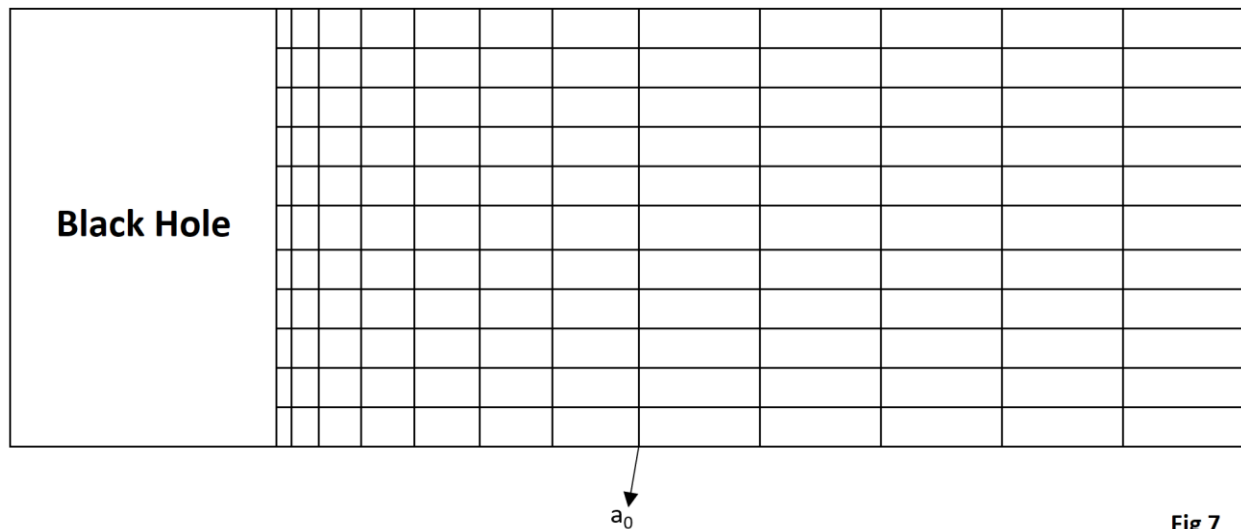


Fig 7

Once the space becomes curved, as is the case around the black hole, the Planck volumes beyond  $a_0$  also become curved, and hence, the relative Planck volumes also gradually increase as the stretching of space-time continues to increase beyond  $a_0$  due to the curvature, as the circumference farther from  $a_0$  should stretch more than the one closer to  $a_0$ . However, the rate of change in the Planck volume for  $a < a_0$  is less than the rate of change for  $a > a_0$ . Therefore, the Newtonian law of gravitation switches to the MOND law of gravitation at  $a_0$ .

Let us consider an elastic rubber band stretched around a cylinder; the radial force exerted by the rubber band on the cylinder would be  $\frac{1}{2\pi}$  times the tension in the rubber band owing to the circumference of the cylinder given by  $2\pi r$ . As the current acceleration of the universe is  $cH_0(2)$  per this theory, the radial acceleration at  $a_0$  should be  $\frac{cH_0}{2\pi}$ , which is  $\sim 1.2 \times 10^{-10} \text{ m/s}^2$  per the above analogy. Here, the whole universe acts like a rubber band wrapped around the black hole at  $a_0$ . Additionally, we can see that the space-time grid ends at the event horizon of the black hole and does not extend into it as in the general theory of relativity.

Below, Figure 8 represents the distortion of the space-time grid around the ordinary matter. As the matter is present only in the four-dimensional space-time grid, it is depicted in red in the middle of Figure 8. Each cube in this figure is of one Planck volume with a Planck charge. As the number of circumferential Planck volumes and the respective area increase radially as we move away from the matter, cutoff acceleration such as  $a_0$  is not applicable, and Newton's law of gravitation can be applied at any acceleration without using any modification such as MOND in the absence of black holes. However, the gravity of ordinary matter will still switch to MOND at  $a_0$  if it is present within the  $a_0$  radius of a black hole owing to the geometry of the space-time as shown in Figure 6 and Figure 7. This explains why the expected gravitational lensing is not observed around the gaseous part of the Bullet cluster (1E 0657-56) (Paraficz et al., 2016) but around the galactic matter as it is subjected to MONDian gravitational lensing due to the presence of black holes at the centers of the galaxies.



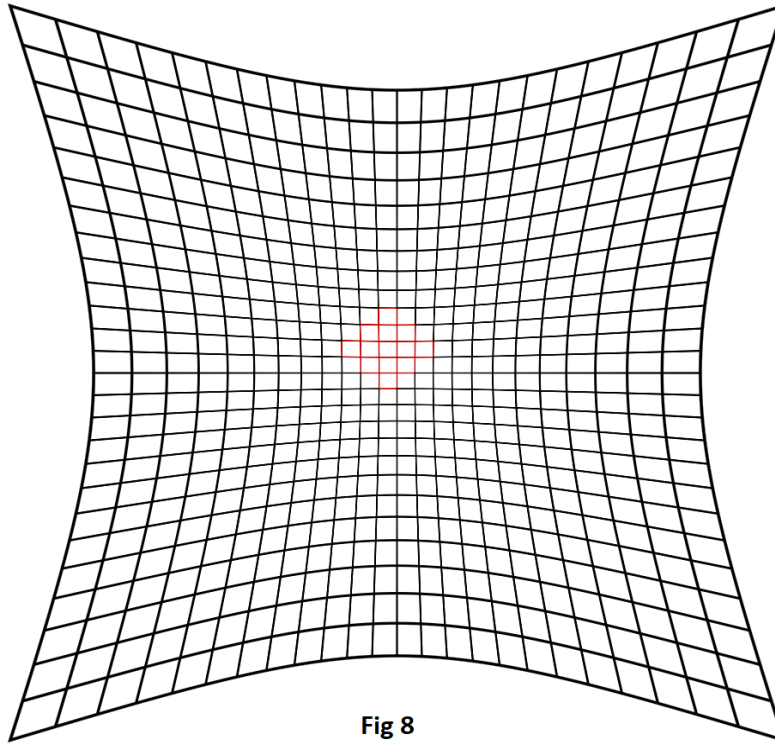


Fig 8

As dark matter has not been detected in galaxy NGC 1052-DF2 (Van Dokkum et al., 2018), this theory predicts that galaxies that do not need dark matter to explain the rotation curves beyond the  $a_0$  orbit should not have black holes at their centers or anywhere in the galaxies.

### 9.1. The relativistic MONDian gravity

The MONDian gravitational law is given below.

$$F = \frac{GMm}{r^2} f\left(\frac{r}{r_o}\right) \text{ (Sanders \& Mcgaugh, 2002)}$$

$f(x) \rightarrow 1$  for  $x \ll 1$ ,  $f(x) \rightarrow x$  for  $x \gg 1$  and  $r_o$  is the radius at which  $a = a_0$ .

We can take the above MOND formula and introduce variable  $G$  and gravitational length-contraction to develop a relativistic law around a stationary black hole similar to the relativistic Newtonian theory of gravity.

$$F = \frac{G(1+z)Mm}{\frac{R_o}{(1+z_o)} \frac{R}{(1+z)}} = (1+z_o)(1+z)^2 \frac{GMm}{R_o R}, \text{ where } z_o \text{ is the gravitational redshift (24) at } R_o, z \text{ is the}$$

gravitational redshift (35) at  $R$ , and  $M$  is the total mass present within the  $a_0$  radius which is  $R_0$ .

$$\frac{dU}{dR} = F = (1+z_o)(1+z)^2 \frac{GMm}{R_o R} = (1+z_o) \left(1 - \frac{U}{mc^2}\right)^2 \frac{GMm}{R_o R} \text{ using (15)}$$

$$\begin{aligned}
\int_U^U \left( \frac{1}{mc^2 - U} \right)^2 dU &= (1 + z_o) \frac{GMm}{R_0(mc^2)^2} \int_R^{R_o} \frac{1}{R} dR \\
\frac{1}{(mc^2 - U_o)} - \frac{1}{(mc^2 - U)} &= (1 + z_o) \frac{GMm}{R_0(mc^2)^2} \ln \left( \frac{R_o}{R} \right), \text{ divide the denominators by } (mc^2). \\
\frac{1}{\left(1 - \frac{U_o}{mc^2}\right)} - \frac{1}{\left(1 - \frac{U}{mc^2}\right)} &= (1 + z_o) \frac{GM}{R_0 c^2} \ln \left( \frac{R_o}{R} \right) \\
\frac{1}{(1 + z_o)} - \frac{1}{(1 + z)} &= (1 + z_o) \frac{GM}{R_0 c^2} \ln \left( \frac{R_o}{R} \right) \text{ using (15)} \\
1 + z &= \frac{(1 + z_o)}{1 - (1 + z_o)^2 \frac{GM}{R_0 c^2} \ln \left( \frac{R_o}{R} \right)} \quad (35)
\end{aligned}$$

Relativistic MONDian gravitational force  $F$  for  $M \geq m$ :

$$F = \left[ 1 + z_o \mu \left( \frac{z}{z_o} \right) \right] (1 + z)^2 \frac{GMm}{R^2} \mu \left( \frac{R}{R_o} \right) \quad (36)$$

where  $\mu(x) \rightarrow 1$  for  $x \ll 1$ ,  $\mu(x) \rightarrow x$  for  $x \gg 1$

The above relativistic MONDian gravitational formula reduces to the relativistic Newtonian formula for gravity (25) for  $R \leq R_o$  and  $z \geq z_o$ .

Relativistic MONDian gravitational potential energy  $U$  for  $M \geq m$ :

$$U = -mzc^2$$

Relativistic MONDian gravitational potential  $\phi$  for  $M \geq m$ :

$$\phi = \frac{U}{m} = -zc^2 \quad (37)$$

As per this theory, any mass present beyond the  $a_0$  radius of a black hole should follow the regular relativistic Newtonian theory of gravity. So, any spherically and uniformly distributed mass beyond the  $a_0$  radius ( $R_0$ ) can be considered as a point mass located at the center of the galaxy exerting relativistic Newtonian gravity at  $R > R_0$  and at any acceleration without needing MOND. So, the total gravitational force  $F_T$  experienced at radius  $R$  beyond the  $a_0$  radius of a galaxy with a central black hole is below.

$$F_T = F_M + F_N \quad (38)$$

where  $F_M$  is the MONDian force due to the mass present within  $R_0$  and  $F_N$  is the Newtonian force due to the mass present between  $R_0$  and  $R$ , where  $R > R_0$ , ignoring the EFE (external field effect) if any. Therefore, as per this theory, even for accelerations greater than  $a_0$  due to the gravitational force  $F_T$ , there can still be a contribution from the MONDian gravity as opposed to the current understanding of MOND where it is applicable only for accelerations less than  $a_0$ .

## 9.2. MONDian gravitational lensing

In the deep-MOND regime, the fictitious gravitational force experienced by a photon can be equated to the fictitious centripetal acceleration of the photon to find the lensing formula around a stationary black hole similar to the formula (31) in the relativistic Newtonian regime.

$$a = (1 + z_o)(1 + z)^2 \frac{GM}{R_o R} = \frac{c^2}{\left(\frac{R}{1 + z}\right)}$$

$(1 + z_o)(1 + z) \frac{GM_0}{R_o c^2} = 1$ , which is the ratio of the gravitational acceleration to the centripetal acceleration of the photon at its closest approach to mass  $M_0$  present within the  $a_0$  radius which is  $R_0$ . Replacing ratio 1 with  $\tan\left(\frac{\pi}{4}\right)$ , we obtain the gravitational lensing equation in the deep-MOND regime. As this ratio should be proportional to the angle of deflection  $\theta$ , we can generalize this equation by replacing  $\tan\left(\frac{\pi}{4}\right)$  with  $\tan\left(\frac{\theta}{4}\right)$  to find the angle of deflection for any radius  $R$ .

$$\theta = 4 \tan^{-1} \left( (1 + z_o)(1 + z) \frac{GM_0}{R_o c^2} \right) \approx \frac{4GM_0}{R_o c^2} = \frac{4\sqrt{GM_0 a_0}}{c^2} \quad (39)$$

Therefore, in the absence of EFE, the gravitational deflection angle  $\theta$  remains almost constant in the deep-MOND regime.

If there is a spherically distributed mass  $M_N$  present between  $R_0$  and  $R$  of a black hole then the total deflection angle at  $R > R_0$  is given by summing the MONDian (39) and Newtonian (31) deflection angles which can be applied for both individual galaxies and galaxy clusters.

$$\theta = 4 \tan^{-1} \left( (1 + z_o)(1 + z_m) \frac{GM_0}{R_o c^2} \right) + 4 \tan^{-1} \left( (1 + z_n)^2 \frac{GM_N}{R c^2} \right) \approx \frac{4\sqrt{GM_0 a_0}}{c^2} + \frac{4GM_N}{R c^2} \quad (40)$$

where  $z_m$  is the MONDian gravitational redshift (35) due to mass  $M_0$  at radius  $R$ ,  $z_n$  is the Newtonian gravitational redshift (24) due to mass  $M_N$  at radius  $R$ , and  $M_0 + M_N$  is the total mass.

### 9.3. Relativistic Poisson equation for MOND

The relativistic MONDian gravitational potential  $\phi$  in Cartesian coordinates based on (37) and (35) is given below. The relativistic Poisson equation outside the point mass can be generated by applying the Laplace operator to the gravitational potential.

$$\frac{\phi}{c^2} = -\frac{(1+z_o)}{1-(1+z_o)^2 \frac{GM}{R_o c^2} \ln \left( \frac{R_o}{\sqrt{x^2 + y^2 + z^2}} \right)} - 1$$

$$\nabla^2 \phi = -\frac{2G^2 M^2 x^2 (1+z_o)^2 (1+z)^3}{(x^2 + y^2 + z^2)^2 R_o^2 c^2} - \frac{2G^2 M^2 y^2 (1+z_o)^2 (1+z)^3}{(x^2 + y^2 + z^2)^2 R_o^2 c^2} - \frac{2G^2 M^2 z^2 (1+z_o)^2 (1+z)^3}{(x^2 + y^2 + z^2)^2 R_o^2 c^2}$$

$$- \frac{2GMx^2 (1+z_o)(1+z)^2}{(x^2 + y^2 + z^2)^2 R_o} - \frac{2GM y^2 (1+z_o)(1+z)^2}{(x^2 + y^2 + z^2)^2 R_o} - \frac{2GM z^2 (1+z_o)(1+z)^2}{(x^2 + y^2 + z^2)^2 R_o} + \frac{3GM (1+z_o)(1+z)^2}{(x^2 + y^2 + z^2) R_o}$$

The relativistic Poisson equation at the point mass can be generated by considering only the positive terms and ignoring the negative terms in the Laplacian ( $\nabla^2 \phi$ ).

$$\nabla^2 \phi = \frac{3GM(1+z_o)(1+z)^2}{(x^2 + y^2 + z^2) R_o} = 4\pi G \rho (1+z_o)(1+z)^2 \frac{R}{R_o}, \text{ where } \rho = \frac{M}{\frac{4}{3}\pi R^3} \text{ is the relativistic mass}$$

density of the point mass present within the  $a_0$  radius.

$$\frac{a_0}{|\nabla \phi|} = \frac{(1+z_o)^3 \frac{GM}{R_o^2}}{(1+z_o)(1+z)^2 \frac{GM}{R_o R}} = \frac{(1+z_o)^2 R}{(1+z)^2 R_o} \text{ based on (25) and (36).}$$

$$\nabla^2 \phi = 4\pi G \rho (1+z)^4 (1+z_o)^{-1} \frac{a_0}{|\nabla \phi|} \quad ; \quad \nabla \cdot \left[ \mu \left( \frac{|\nabla \phi|}{a_0} \right) \nabla \phi \right] = 4\pi G \rho (1+z)^4 \left[ 1 + z_o \xi \left( \frac{z}{z_o} \right) \right]^{-1}$$

Relativistic Poisson equation for MOND:

$$\nabla \cdot \left[ \mu \left( \frac{|\nabla \phi|}{a_0} \right) \nabla \phi \right] = 4\pi G \rho \left( 1 - \frac{\phi}{c^2} \right)^4 \left[ 1 - \frac{\phi_o}{c^2} \xi \left( \frac{\phi}{\phi_o} \right) \right]^{-1} \quad (41)$$

where  $\mu(x) \rightarrow 1$  for  $x \gg 1$ ,  $\mu(x) \rightarrow x$  for  $x \ll 1$  and  $\xi(x) \rightarrow 1$  for  $x \ll 1$ ,  $\xi(x) \rightarrow x$  for  $x \gg 1$

The above relativistic Poisson equation for MOND (41) reduces to the relativistic Poisson equation in the Newtonian regime (33) for  $\phi \geq \phi_o$  and  $|\nabla \phi| \geq a_0$ , and to the nonrelativistic equation when the gravitational redshift  $z$  is considered negligible. The nonrelativistic Poisson equation based on AQUAL (“A QUAdratic Lagrangian”) is given below.

$$\nabla \cdot \left[ \mu \left( \frac{|\nabla \phi|}{a_0} \right) \nabla \phi \right] = 4\pi G \rho \quad (\text{Mamon et al., 2005}), \text{ where } \mu(x) \rightarrow 1 \text{ for } x \gg 1, \mu(x) \rightarrow x \text{ for } x \ll 1.$$

#### 9.4. Relativistic AQUAL for MOND

Relativistic AQUAL  $\mathcal{L}$  and the respective action  $\mathcal{S}$  based on (41) are below.

$$\mathcal{L} = -\frac{1}{8\pi G} a_0^2 F \left( \frac{|\nabla \phi|^2}{a_0^2} \right) - \int \rho \left[ \left( 1 - \frac{\phi}{c^2} \right)^4 \left( 1 - \frac{\phi_o}{c^2} \xi \left( \frac{\phi}{\phi_o} \right) \right)^{-1} \right] d\phi \quad (42)$$

Where  $F(x) \rightarrow \frac{2}{3}x^{\frac{3}{2}}$  for  $x \ll 1$ ,  $F(x) \rightarrow x$  for  $x \gg 1$  and  $\xi(x) \rightarrow 1$  for  $x \ll 1$ ,  $\xi(x) \rightarrow x$  for  $x \gg 1$

$$\mathcal{S} = \iint \mathcal{L} d^3r dt$$

#### 9.5. MOND in galaxy clusters

The need for dark matter is not completely eliminated when MOND is applied to the galaxy clusters (Sanders, 2018) as it is applied only for accelerations less than  $a_0$ . As there can be a contribution from MONDian gravity even for accelerations greater than  $a_0$  as per this theory (38), it can be applied to the galaxy clusters to eliminate the need for dark matter.

For a simple gravitationally bound system such as a point mass  $m$  rotating around mass  $M$  with  $M \gg m$ , the total kinetic energy  $T$  is given by  $T = \frac{1}{2}mv^2$ , the Newtonian potential energy  $U$  is

given by  $U = -\frac{GMm}{R}$ , and the rotational velocity  $v$  is given by  $v = \sqrt{\frac{GM}{R}}$ . Substituting  $v$  in  $T$ , we

get the regular Newtonian-based virial theorem which is  $\langle T \rangle = -\frac{1}{2}\langle U \rangle$ . However, when both the

Newtonian and MONDian dynamics are involved, the regular virial theorem cannot be applied to a system that is gravitationally bound or in hydrostatic equilibrium with the binding potential such as a galaxy cluster and hence the need for dark matter to fit the cluster dynamics into the regular virial theorem. We can eliminate the need for dark matter in the galaxy clusters using the updated virial theorem with both Newtonian and MONDian gravity included.

$$F_T = F_M + F_N = \frac{GM_0}{RR_0} + \frac{GM_N}{R^2} \quad (38), \text{ where } M_0 \text{ is the mass present within } R_0, M_N \text{ is the spherically}$$

distributed mass present between  $R_0$  and  $R$ , where  $R > R_0$ , and  $M_0 + M_N$  is the total baryonic mass.

$$\frac{GM_0}{RR_0} + \frac{GM_N}{R^2} = \frac{v^2}{R} \quad ; \quad v = \sqrt{\frac{GM_0}{R_0} + \frac{GM_N}{R}} \quad ; \quad T = \frac{1}{2}m \left( \frac{GM_0}{R_0} + \frac{GM_N}{R} \right) \text{ which is the updated virial}$$

theorem based on this theory for a simple gravitationally bound system with contributions from both Newtonian and MONDian gravity.

We can apply the above updated virial theorem to a galaxy cluster by using two point masses located at the center of the cluster, one point mass  $M_N$  exerting Newtonian gravity and the other point mass  $M_0$  exerting MONDian gravity. The sum of all the baryonic masses within the  $a_0$  radius of the individual galaxies in the cluster would become the mass of the point mass  $M_0$  exerting MONDian gravity, and the sum of the rest of the baryonic masses in the cluster would become the mass of the point mass  $M_N$  exerting Newtonian gravity in the absence of intergalactic black holes.

In the galaxy clusters, let us consider that about one tenth of the total baryonic mass is present within the  $a_0$  radius of the individual galaxies with central black holes, i.e.,  $M_N=9M_0$  and consider a gravitationally bound point mass  $m$  rotating around the cluster at its very edge, we can identify the additional dynamical mass (dark matter) predicted by the regular Newtonian virial theorem by comparing it with the updated (Newtonian + MONDian) virial theorem.

$$T = \frac{1}{2} m \left( \frac{GM_0}{R_0} + \frac{GM_N}{R} \right) = x \frac{1}{2} m \frac{G(M_0 + M_N)}{R}, \text{ where } x \text{ is the factor for dark matter.}$$

$$\text{For } M_N=9M_0, T = \frac{1}{2} m \left( \frac{GM_0}{R_0} + \frac{9GM_0}{R} \right) = x \left( \frac{1}{2} m \frac{10GM_0}{R} \right); a_0 = \frac{GM_0}{R_0^2}; R_0 = \sqrt{\frac{GM_0}{a_0}}$$

For the Coma cluster, solving the above quadratic equation for the dark matter factor  $x$  with the radius  $R \approx 6.17 \times 10^{23}$  ( $\sim 20$  Mpc), the Newtonian gravitational mass  $x(10M_0) \approx 6.2 \times 10^{15}$  solar masses (Chernin et al., 2013), and  $a_0 = 1.2 \times 10^{-10}$  in MKS units, we get  $x \approx 7.24$ . Therefore, the regular Newtonian virial theorem predicts that the total gravitational mass of the Coma cluster is  $\sim 7.24$  times the total baryonic mass which means that  $\sim 86$  percent of the total gravitational mass in the Coma cluster is identified as dark matter.

For the Virgo cluster, with the Newtonian gravitational mass  $x(10M_0) \approx 7.4 \times 10^{14}$  and the radius  $R \approx 2.4 \times 10^{23}$  ( $\sim 7.4$  Mpc) (Kashibadze et al., 2020) in MKS units, we get  $x \approx 8.74$ . Therefore,  $\sim 88$  percent of the total gravitational mass in the Virgo cluster is identified as dark matter. As the percentages of the dark matter predicted in the galaxy clusters are in line with the standard model of cosmology and also completely accounted for with the updated virial theorem using only the baryonic mass, this theory completely eliminates the need for dark matter in galaxy clusters. Therefore, the generalized virial theorem for any gravitationally bound system having both MONDian and Newtonian dynamics based on (35) and (37) for  $z_0 \approx 0$  is below, where  $R$  is the virial radius.

$$\langle T \rangle = -\frac{1}{2} \left( \frac{\langle U_0 \rangle}{\ln \left( \frac{R_0}{R} \right)} + \langle U_N \rangle \right)$$

## 10. Conclusions:

This theory proves that both space-time and matter expand at the Hubble rate like an elastic ruler with constant proper distance and volume, establishes the relativistic acceleration of space-time by proposing cosmological time-dilation similar to gravitational time-dilation and proves the relativistic Hubble law. Fundamental constants are shown to vary proportionally to the Hubble rate

of the universe and propose a relativistic Newtonian theory of gravity and relativistic MOND in the context of an absolute reference frame. The derived Schwarzschild radius, photon sphere, gravitational lensing, perihelion precession of mercury, and Shapiro time delay are shown to agree with the general theory of relativity or even exceed it in accuracy without using any weak field approximations or higher order terms. The relativistic GEM equations are derived to prove the existence of gravitational waves and the gravitomagnetic effects similar to the general theory of relativity. This theory predicts the gravitational/cosmological time-dilation, the length-contraction and eliminates black hole singularities whereas the general theory of relativity predicts the gravitational time-dilation, the gravitational length-expansion, singularities, and the redshift-based apparent cosmological time-dilation due to the creation of new space and does not predict the inherent cosmological time-dilation in the space-time fabric, the cosmological length-contraction, the cosmological potential energy, the expansion of matter with constant proper distance and volume. This theory clearly defines the relativistic scalar gravitational potential whereas the general theory of relativity does not have a clear definition. The general theory of relativity does not explain the mechanism of dark energy that accelerates the expansion of the universe and requires dark matter that is not yet discovered whereas this theory successfully explains the dark energy and dark matter in galaxy clusters. Overall, this theory proposes a hidden extradimensional electrostatic background to gravity to explain the curvature of space-time, relativistic gravity, dark energy, dark matter in galaxy clusters, cosmological redshift, cosmological time-dilation, and eliminates cosmic inflation, the cosmic event horizon, the cosmic scale factor ( $a$ ), the critical density ( $\Omega$ ), the cosmological constant ( $\Lambda$ ), and black hole singularities.

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## Methods:

The computations in this paper were performed by using Maple™. Maple 2020.2. Maplesoft, a division of Waterloo Maple Inc., Waterloo, Ontario. Maple is a trademark of Waterloo Maple Inc.

## Data availability statement:

Data sharing is not applicable to this article as no new data were created or analyzed in this study.